1) **Pantograph Drag**: A high-speed train picks up its electrical power via a pantograph from an overhead line. The locomotive travels at speed $U$ and the pantograph exerts a constant vertical force $F$ on the power line.

We make the usual small amplitude approximations and assume (not unrealistically) that the line is supported in such a way that its vertical displacement obeys an inhomogeneous Klein-Gordon equation

$$\rho \ddot{y} - T y'' + \rho \Omega^2 y = F \delta(x - Ut),$$

with $c = \sqrt{T/\rho}$ being the velocity of propagation of short-wavelength transverse waves on the overhead cable.

a) Assume that $U < c$ and solve for the steady state displacement of the cable about the pickup point. (Hint: the disturbance is time-independent when viewed from the train.)

b) Now assume that $U > c$. Again find an expression for the displacement of the cable. (The same hint applies, but the physically appropriate boundary conditions are very different!)

c) By equating the rate at which wave-energy

$$E = \int \left\{ \frac{1}{2} \rho \dot{y}^2 + \frac{1}{2} Ty'^2 + \frac{1}{2} \rho \Omega^2 y^2 \right\} \, dx$$

is being created to rate at which the locomotive is doing work, calculate the wave-drag on the train. In particular, show that there is no drag at all until $U$ exceeds $c$. (Hint: While the front end of the wake is moving at speed $U$, the trailing end of the wake is moving forward at the *group velocity* of the wave-train.)

d) By carefully considering the force the pantograph exerts on the overhead cable, again calculate the induced drag. You should get the same answer as in part c) (Hint: The tension in the cable is the same before and after the train has passed, but the direction in which the tension acts is different. The force $F$ is therefore not *exactly* vertical, but has a small forward component. Don’t forget that the resultant of the forces is accelerating the cable.)

This problem of wake formation and drag is related both to Čerenkov radiation and to the Landau criterion for superfluidity.
2) Non-linear Waves:

a) Suppose a fluid has equation of state $P = \lambda^2 \rho^3 / 3$. From the continuity equation

$$\partial_t \rho + \partial_x \rho v = 0,$$

and Euler’s equation of motion

$$\rho (\partial_t v + v \partial_x v) = -\partial_x P,$$

deduce that

$$\left( \frac{\partial}{\partial t} + (\lambda \rho + v) \frac{\partial}{\partial x} \right) (\lambda \rho + v) = 0,$$

$$\left( \frac{\partial}{\partial t} + (-\lambda \rho + v) \frac{\partial}{\partial x} \right) (-\lambda \rho + v) = 0.$$

In what limit do these equations become equivalent to the wave equation for one-dimensional sound? What is the sound speed in this case?

b) Show that the Riemann invariants $v \pm \lambda \rho$ are constant on suitably defined characteristic curves. What is the local speed of propagation of the waves moving to the right or left?

c) The fluid starts from rest, $v = 0$, but with a region where the density is higher than elsewhere. Show that the Riemann equations will inevitably break down at some later time due to the formation of shock waves.

The notion of Riemann invariants can be extended to other equations of state, but the expressions are more complicated, and the left and right-going waves interact with each other.

3) Burgers Shocks: As simple mathematical model for the formation and decay of a shock wave consider Burgers’ Equation:

$$\partial_t u + u \partial_x u = \nu \partial_x^2 u.$$

Note its similarity to the Riemann equations of the previous question. The additional term on the right-hand side introduces dissipation and prevents the solution becoming multivalued.

a) Show that if $\nu = 0$ any solution of Burgers’ equation having a region where $u$ decreases to the right will always eventually become multivalued.

b) Show that the Hopf-Cole transformation, $u = -2\nu \partial_x \ln \psi$, leads to $\psi$ obeying a heat diffusion equation

$$\partial_t \psi = \nu \partial_x^2 \psi.$$
c) Show that

\[ \psi(x, t) = Ae^{\nu a^2 t - ax} + Be^{\nu b^2 t - bx} \]

is a solution of the heat equation, and so deduce that Burgers’ equation has a shock-wave-like solution which travels to the right at speed \( C = \nu (a + b) = \frac{1}{2}(u_L + u_R) \), the mean of the wave speeds to the left and right of the shock. Show that the width of the shock is \( \approx 4\nu / |u_L - u_R| \).