

AC Measurement of Magnetic Susceptibility

Episode 2 – Loss and Thermal Effects

Prof. Jeff Filippini

Physics 401

Spring 2020



Outline

Combine the tools and techniques we've learned to characterize **magnetic properties of materials**

- **Magnetic Materials:** Dia- and paramagnetism (Bonus: superconductors!)
- **Week 1 Refresher:** AC Measurement of Susceptibility
- **Magnetic Losses:** Complex permeability
- **Temperature Dependence:** Curie-Weiss Law
- **Finishing the Final Lab:** What you'll get, what you need to do
- *Closing out the Semester*

This is the **second week** of the **final lab**

Report counts as your **final exam**



Reminder: Magnetic Response of Materials

Two things are often called the “magnetic field”: **B** and **H**

B

Magnetic induction
Magnetic flux density

Determines forces **on**
moving *free* charges
via **Lorentz force law**:

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

$$\vec{B} = \mu_0 (\vec{H} + \vec{M})$$

M

Magnetization
Magnetic polarization

Field created only **by**
moving *bound* charges,
i.e. magnetic response
of the medium

H

Magnetic field intensity
Magnetizing field

Field created only **by**
moving *free* charges.

In vacuum, $B = \mu_0 H$.

Reminder: Magnetic Response of Materials

$$\vec{B} = \mu_0(\vec{H} + \vec{M}) = \mu_0(1 + \chi)\vec{H} = \mu_0\mu_r\vec{H}$$

We classify materials into three major categories:

Diamagnetic	$\chi < 0$	$\mu_r < 1$	Weakly repelled
Paramagnetic	$\chi > 0$	$\mu_r > 1$	Weakly attracted
Ferromagnetic	$\chi \gg 0$	$\mu_r \gg 1$	Strongly attracted

Magnetic Materials: Diamagnetism

Many materials are **diamagnetic** ($\chi < 0$), i.e. they magnetize in a way that reduces the applied field and are repelled from magnets.

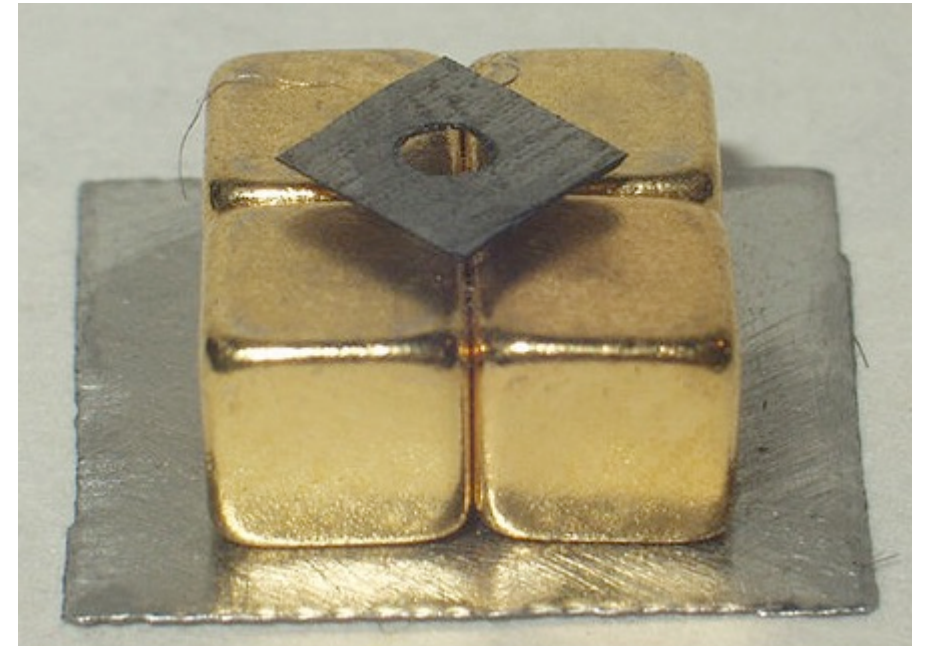
Effect of **paired electrons** in all materials, generally very weak

Material	χ_v (10^{-5})
Bismuth	-16.6
Carbon (diamond)	-2.1
Carbon (graphite)	-1.6
Copper	-1.0
Lead	-1.8
Mercury	-2.9
Pyrolytic carbon	-40.0
Silver	-2.6
Superconductor	-10^5
Water	-0.91

[Wikipedia](#)

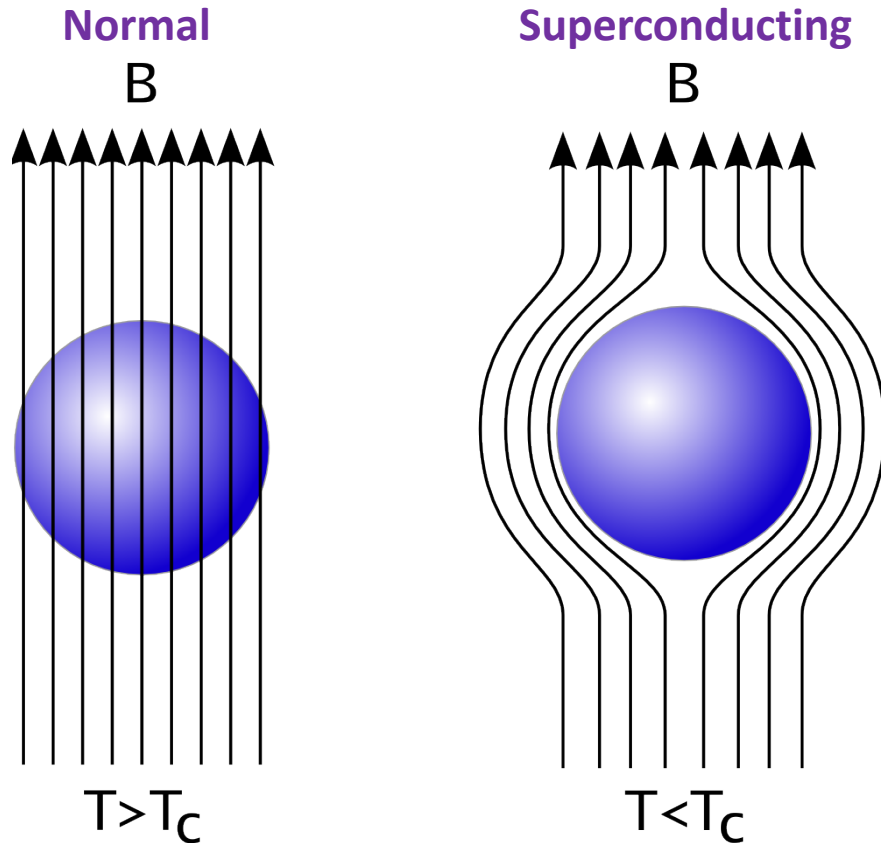


Frog levitated in 16T field
(Wikipedia)



Levitating sample of pyrolytic carbon (Wikipedia)

Superconductors: the Meissner Effect



[Wikipedia](#)

Superconductors are **perfect diamagnets** ($\chi = -1$)

At $T < T_c$, electrons **condense** into a macroscopic quantum state that can flow without loss.

Persistent image currents expel incident fields to maintain $B=0$. Above some **critical field H_c** , superconductivity is destroyed.

Important for designing superconducting solenoid magnets

Deep analogy with the Higgs mechanism in particle physics!

Magnetic Materials: Paramagnetism

In **paramagnetic** ($\chi > 0$) materials, atomic/molecular dipoles from unpaired electron spins align with **external fields**.

Material	Magnetic susceptibility, χ_v [10^{-5}]
Tungsten	6.8
Caesium	5.1
Aluminium	2.2
Lithium	1.4
Magnesium	1.2
Sodium	0.72

[Wikipedia](#)



Liquid oxygen trapped between poles of a strong magnet (Wikipedia)

Magnetic Materials: Ferromagnetism

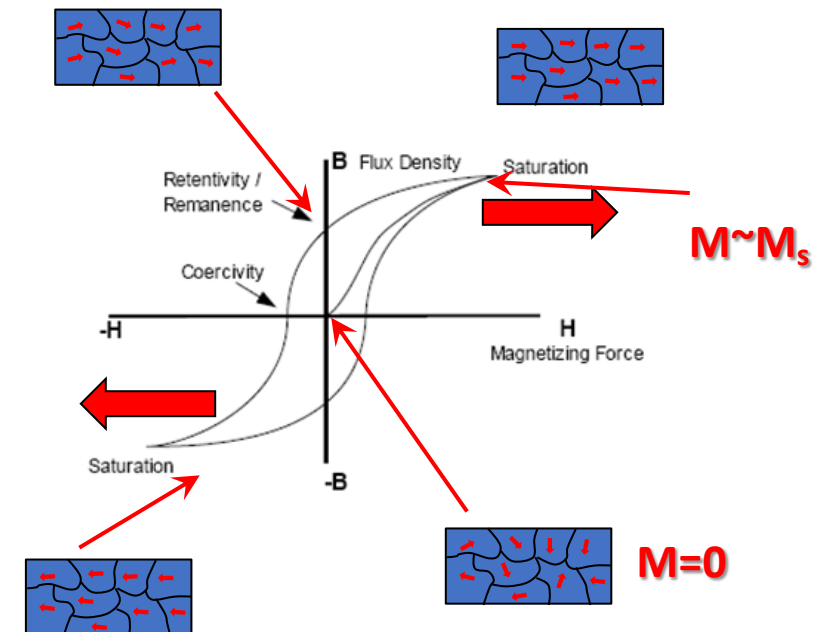
Material	μ_r	B_{rem} (T)
Fe, 99.8% pure	5000	1.3
Permalloy	100,000	0.7
Superpermalloy	1,000,000	0.7
Co, 99% pure	250	0.5
Ni, 99% pure	600	0.4

Wikipedia

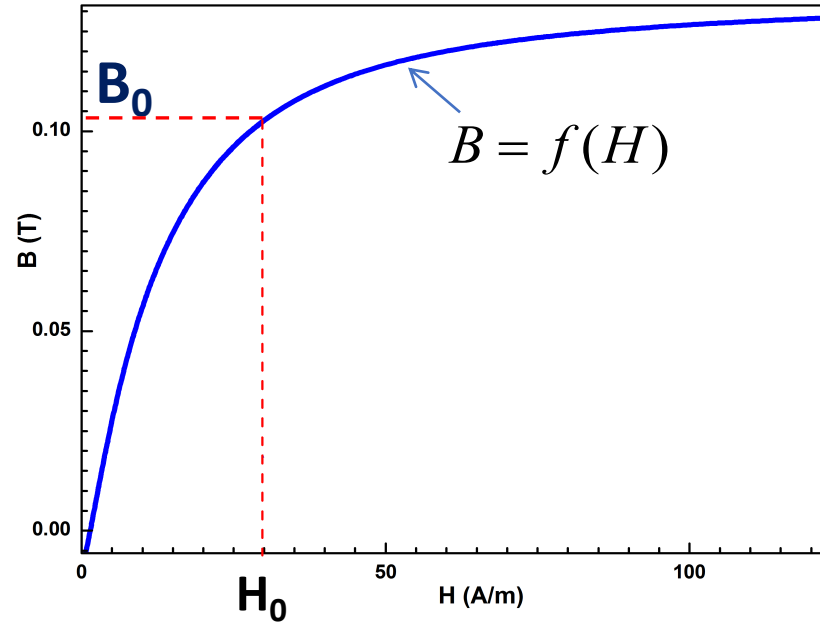
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In **ferromagnetic** ($\chi \gg 1$) materials, neighboring atomic/molecular dipoles from unpaired electron spins **spontaneously align** with one another, forming **domains**.

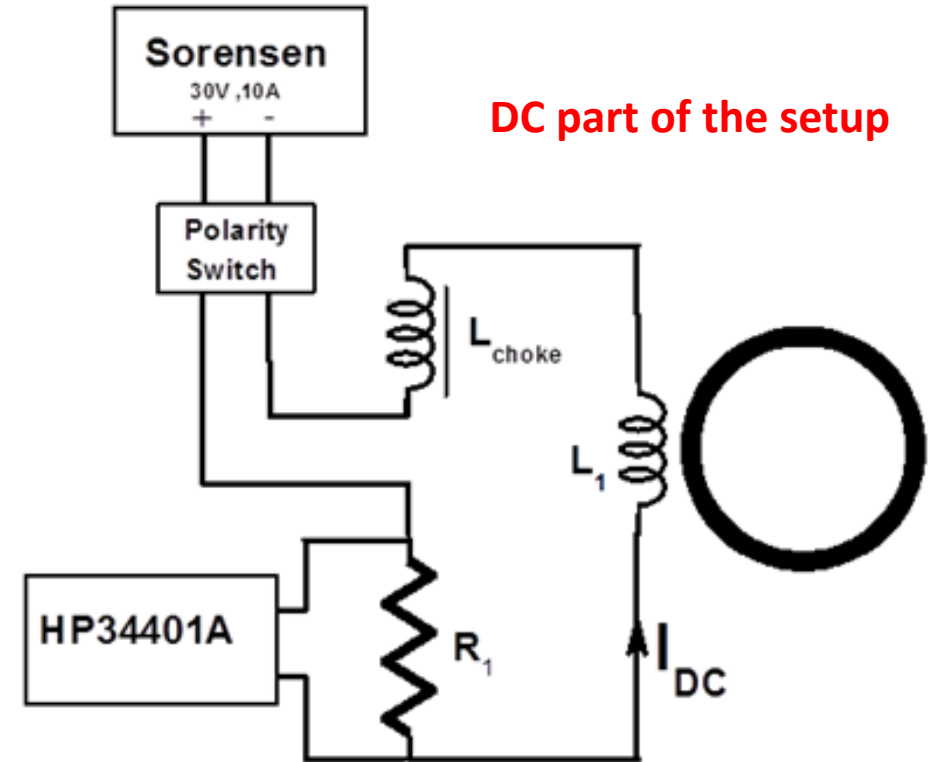


Measuring Permeability: DC Field



$$H_0 = \frac{N_p I_{DC}}{2\pi r}$$

$$\mathbf{H} = \mathbf{H}_0 + \mathbf{H}_1 \sin \omega t$$



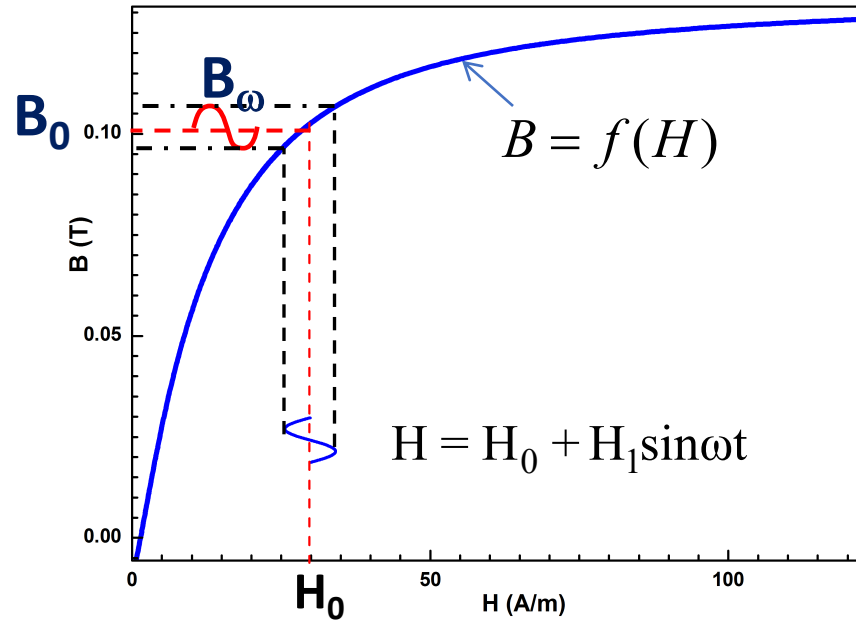
DC part of the setup

Here N_p is the number of turns in the DC primary coil

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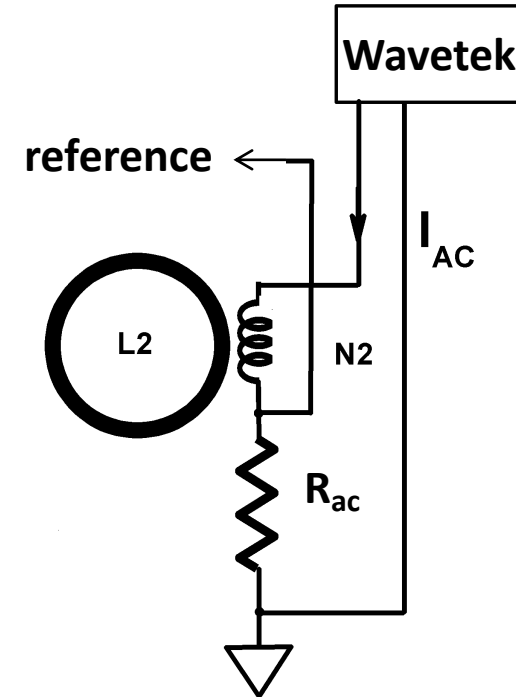


Measuring Permeability: AC Modulation



$$H = H_0 + H_1 \sin \omega t$$

$$H_1 = \frac{N_p I_{AC}}{2\pi r}$$



$$B_\omega \sim \frac{df}{dH} = \frac{dB}{dH} = \mu = \mu_0 \mu_r$$

Here N_p is the number of turns in the AC primary coil

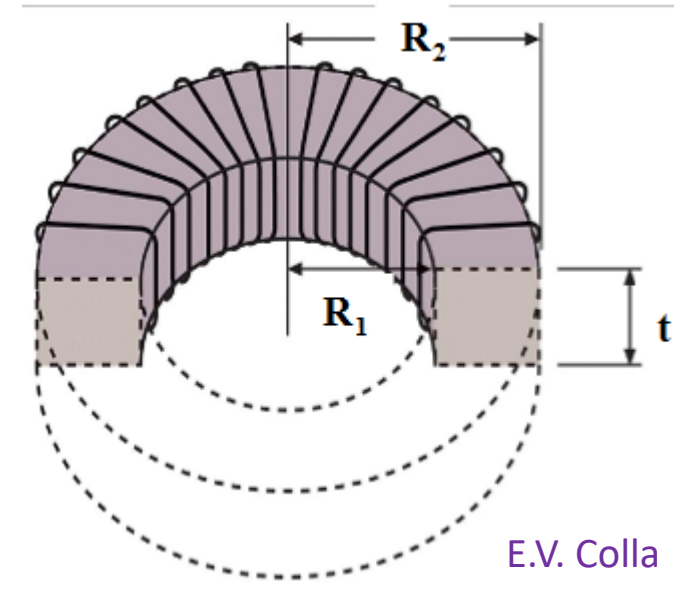
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Measuring Permeability: AC Toroid Flux

Primary coil is a **toroid** of N_p turns carrying a current I_p creates a magnetic field H :

$$H = \frac{N_p I_p}{2\pi r}$$

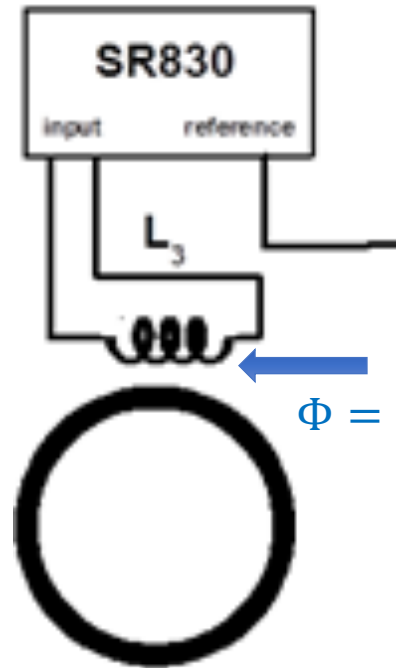


... and this magnetizing field adds a flux $d\Phi$ to *each turn* of the pickup coil

$$d\Phi = \mu \int \vec{H} \cdot d\vec{a} = \frac{\mu I N t}{2\pi} \int_{R_1}^{R_2} \frac{dr}{r} = \frac{\mu I N t}{2\pi} \ln \frac{R_2}{R_1}$$

Measuring Permeability: Pickup Coil

Faraday's Law



$$V_{lock-in} = -N_{pickup} \frac{d\Phi}{dt} \quad \Phi = \int \vec{B} \cdot d\vec{a} = \mu \int \vec{H} \cdot d\vec{a}$$

Only the AC flux contributes to this time derivative:

$$\Phi_1 = \mu \int \vec{H}_{ac} \cdot d\vec{a} = \frac{\mu I_{ac} N t}{2\pi} \ln \frac{R_2}{R_1}$$

This AC flux is sourced by the current in coil L2: $I_{ac} = \frac{V_0 \sin \omega t}{R_{ac}}$

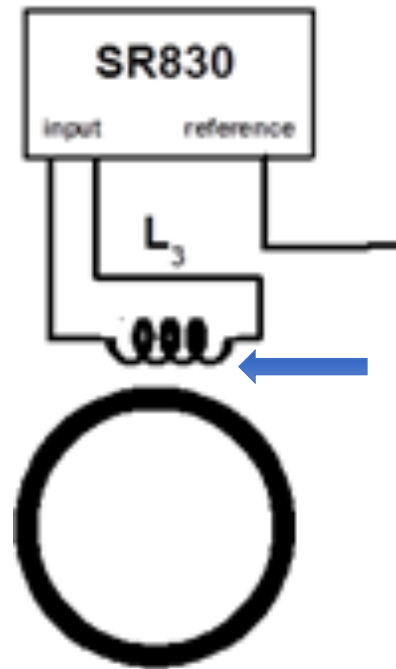
$$V_{lock-in} = -N_{pickup} \frac{\mu N t}{2\pi} \ln \frac{R_2}{R_1} \frac{dI_{ac}}{dt} = -N_{pickup} \frac{\mu N t}{2\pi} \ln \frac{R_2}{R_1} \omega \cos \omega t$$

$$V_{lock-in} = -\mu_r L_0 \frac{V_{ac}}{R_{ac}} \omega \cos \omega t$$

where $L_0 \equiv N_{pickup} \frac{\mu_0 N t}{2\pi} \ln \frac{R_2}{R_1}$

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Measuring Permeability: Pickup Coil



Relative permeability of core

Time derivative of source current

Geometry of the toroid

$$V_{lock-in} = -\mu_r L_0 \frac{V_{ac}}{R_{ac}} \omega \cos \omega t$$

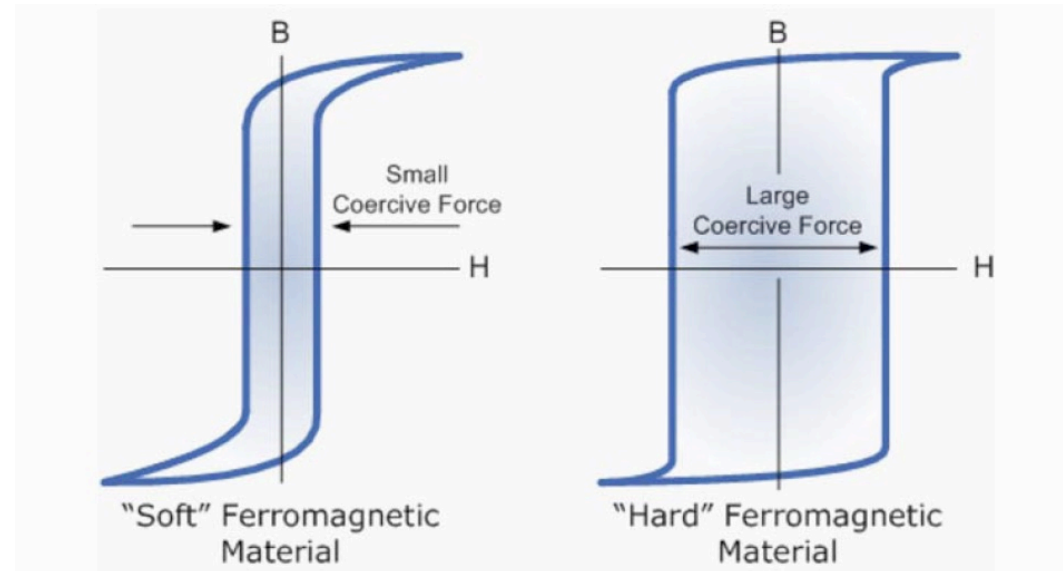
Geometric mutual inductance in the absence of the core

$$L_0 \equiv N_{pickup} \frac{\mu_0 N t}{2\pi} \ln \frac{R_2}{R_1}$$

In general: $\mu_r = \mu' - i\mu''$

- The **real** component μ' appears at the **Y** (quadrature phase) channel of the lock-in, due to its $\pi/2$ **phase shift** with respect to the driving current (cos vs. sin!)
- The **imaginary** component μ'' characterizes losses, and appears **in-phase**

Hysteresis, Coercivity, and Work



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Hysteresis necessarily involves energy losses from re-magnetization.
If the domains are “sticky”, we need to do work to overcome that.

$$W = V \int \vec{H} \cdot d\vec{B}$$

For uniform fields over volume V
(analogous to $dW = F dx$)

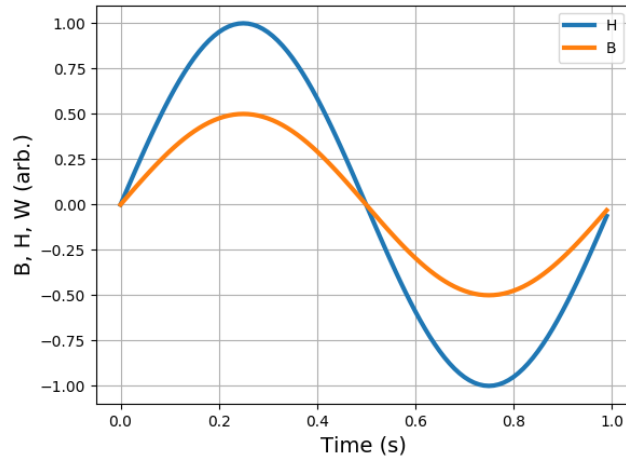
$$W_{loop} = V \oint \vec{H} \cdot d\vec{B} = V * A_{loop}$$

Aside: Loss from Complex Permeability

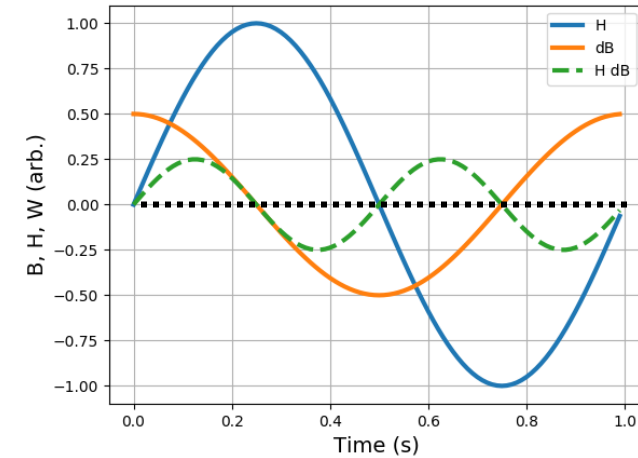
$$\vec{B} = \mu_0(\vec{H} + \vec{M}) = (\mu' - i\mu'')\mu_0\vec{H}$$

Why is a material with complex permeability ($\mu'' \neq 0$) lossy?

Real μ_r :

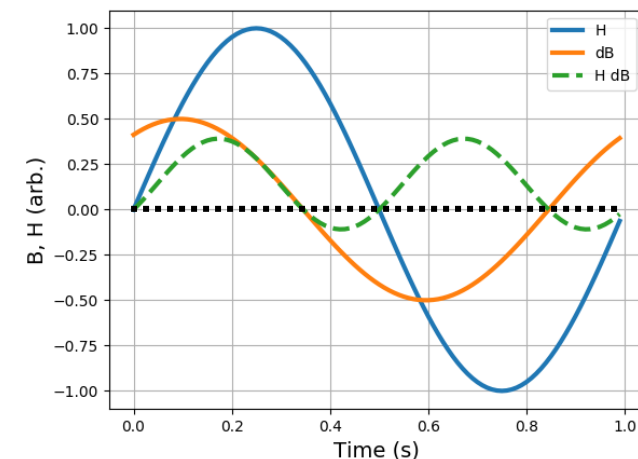
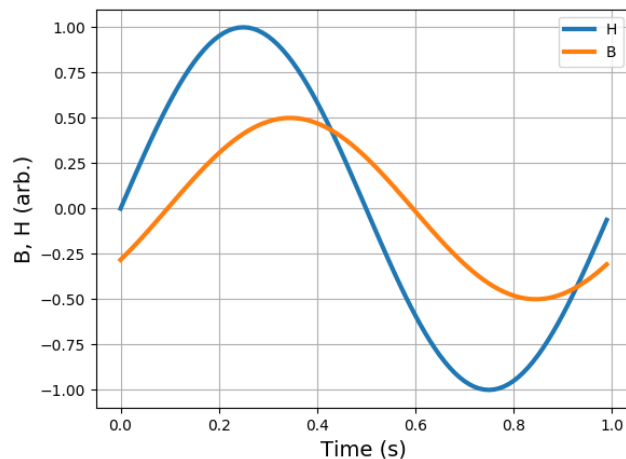


In analogy with
 $dW = F dx$,
 we have
 $dW = H dB$



Zero
 integral

Complex μ_r :



Nonzero
 integral!



Calculating the Magnetic Induction, B

$$B = \mu_0(1 + \chi)H = \mu_0\mu_r H = \mu H$$
$$\mu = \mu_0\mu_r = \frac{dB}{dH} \Rightarrow B = \mu_0 \int \mu_r(H) dH$$

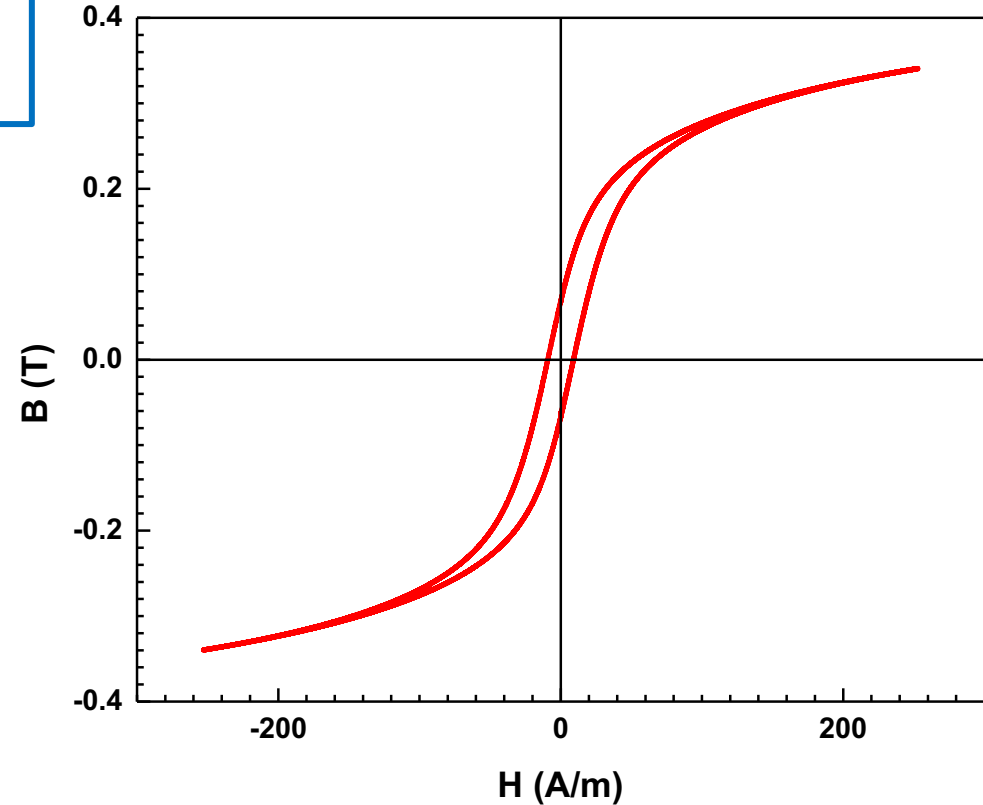
Integrate using OriginPro to find B(H) contour

Work for any motion through this plane:

$$W = V \int \vec{H} \cdot d\vec{B}$$

For uniform fields over volume V
(analogous to $dW = F dx$)

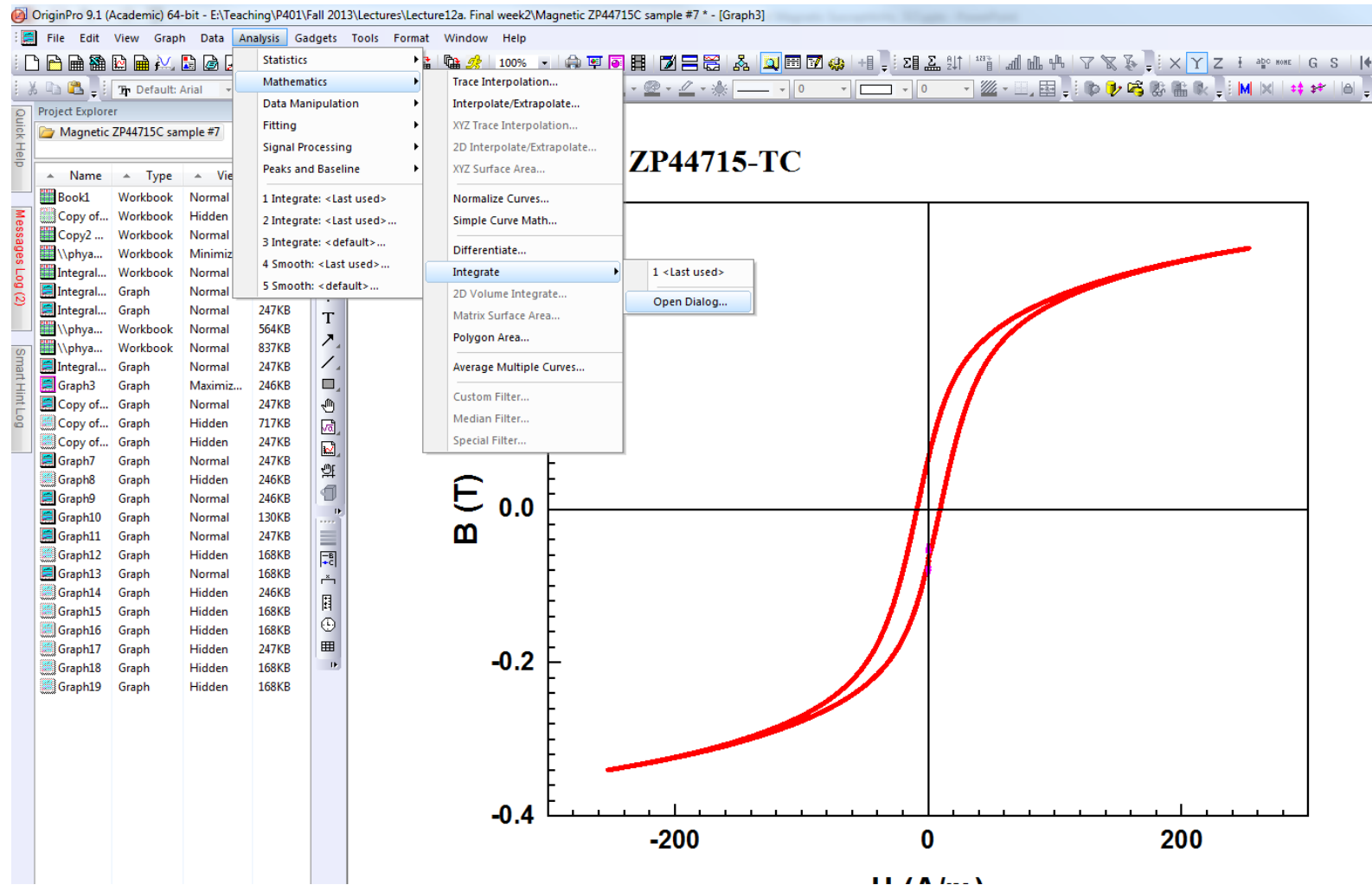
Magnetics ZP44715-TC



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Hysteresis and Remagnetization Losses

$$W_{loop} = V \oint \vec{H} \cdot d\vec{B} = V * A_{loop}$$

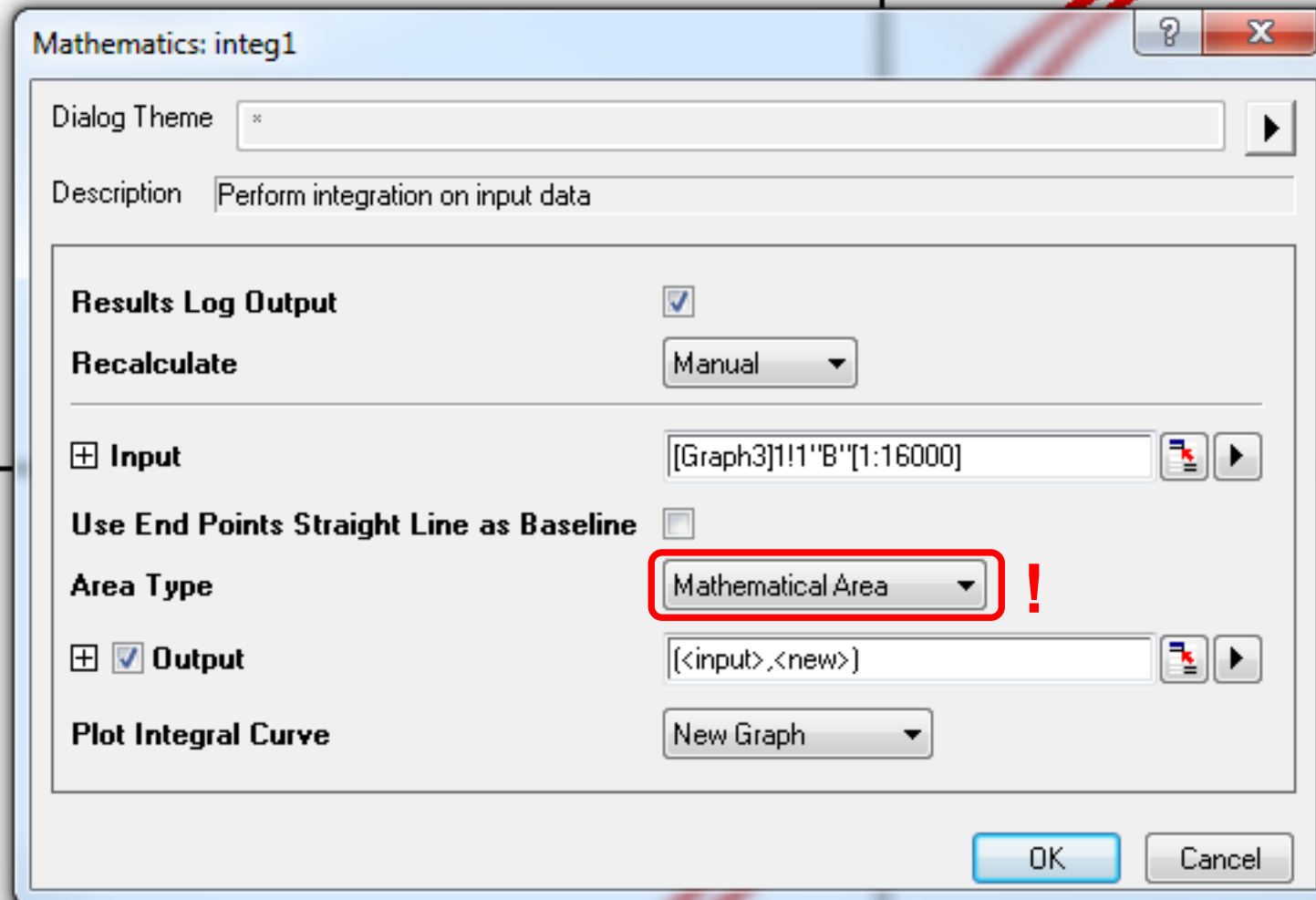


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Hysteresis and Remagnetization Losses

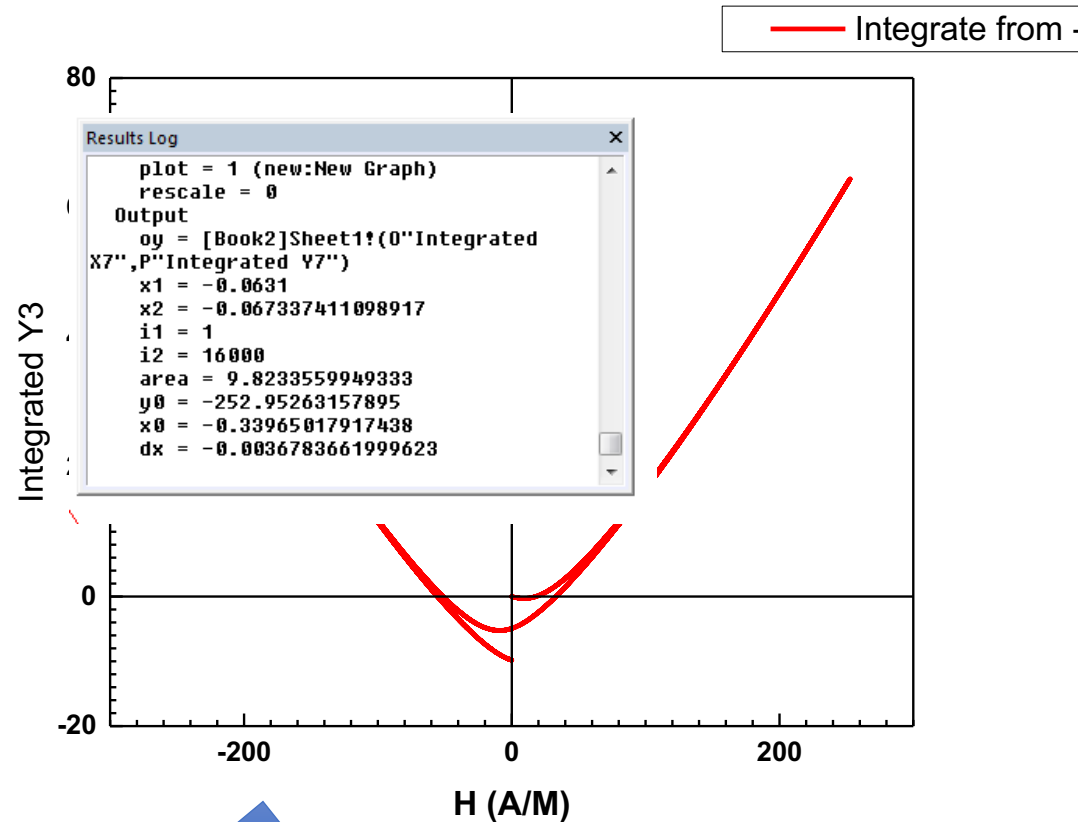
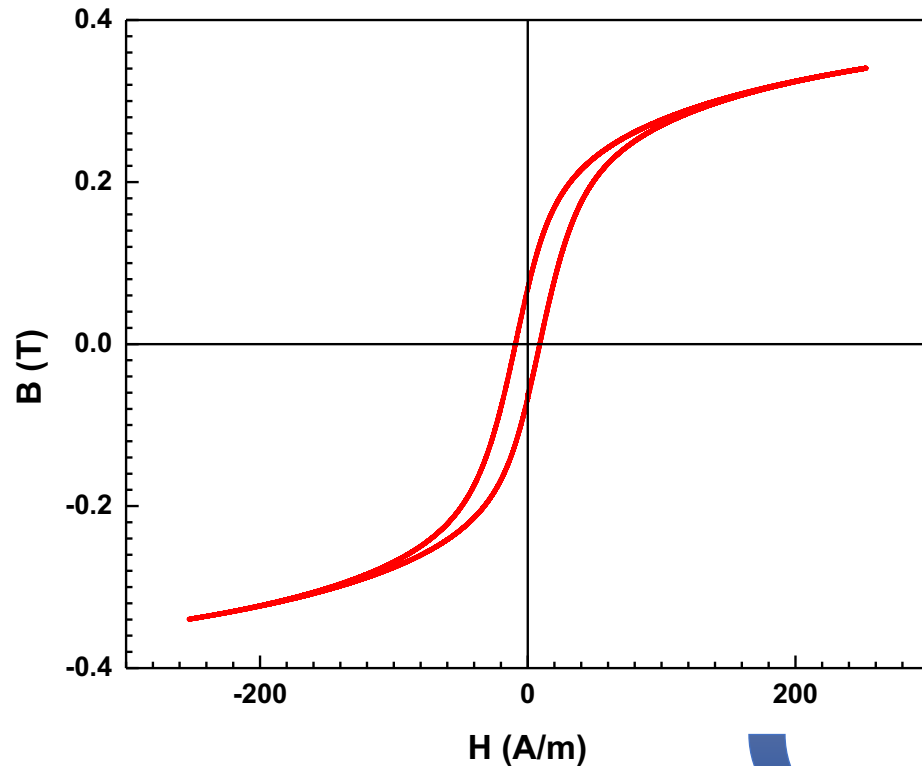
$$W_{loop} = V \oint \vec{H} \cdot d\vec{B} = V * A_{loop}$$



Hysteresis and Remagnetization Losses

$$W_{loop} = V \oint \vec{H} \cdot d\vec{B} = V * A_{loop}$$

Magnetics ZP44715-TC



Hysteresis and Remagnetization Losses

$$W_{loop} = V \oint \vec{H} \cdot d\vec{B} = V * A_{loop}$$

```
Results Log
plot = 1 (new:New Graph)
rescale = 0
Output
oy = [Book2]Sheet1!(0"Integrated
x7",P"Integrated Y7")
x1 = -0.0631
x2 = -0.067337411098917
i1 = 1
i2 = 16000
area = 9.8233559949333
y0 = -252.95263157895
x0 = -0.33965017917438
dx = -0.0036783661999623
```

Units :

$$V(\text{volume}) \rightarrow m^3$$

$$H(\text{field}) \rightarrow A \cdot m^{-1}$$

$$B(\text{magn.induction}) \rightarrow kg \cdot s^{-2} \cdot A^{-1}$$

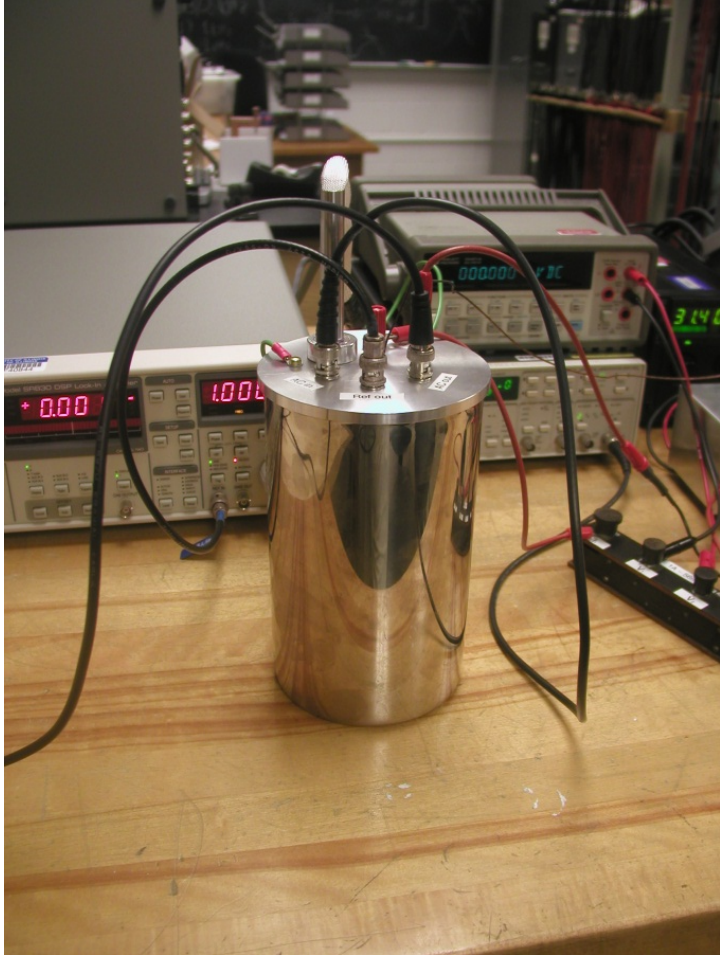
$$[V \cdot B \cdot H] \rightarrow m^2 \cdot kg \cdot s^{-2} \equiv J(\text{joule})$$

This is the **energy** lost around a circuit of the loop.

To make this a **loss power**, we need to account for the frequency:

$$P_{loss} = \frac{W}{T} = Wf, \text{ where } T \text{ is the period and } f \text{ is the frequency}$$

Temperature Dependence of Permeability



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In this experiment we measure $\chi(T)$: magnetic permeability as a function of temperature

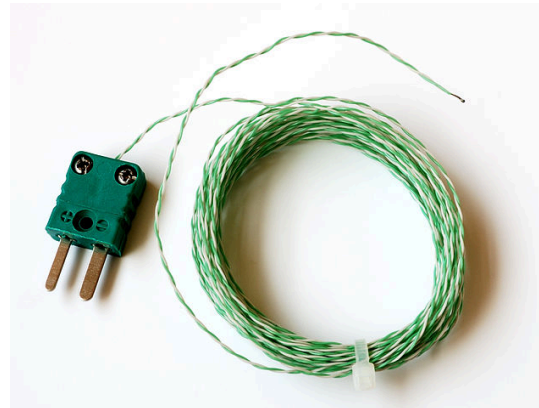
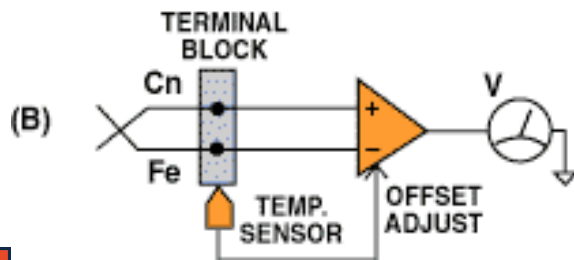
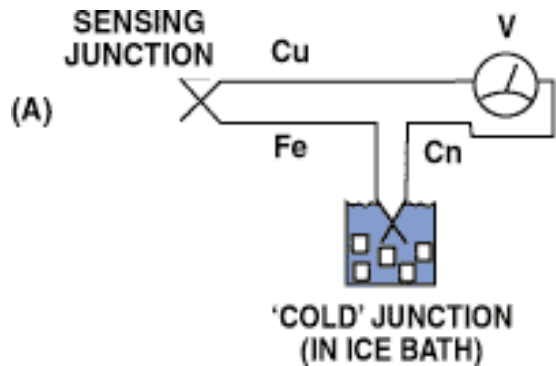
We measure at **fixed H_0** , and so fixed I_{DC} .
The default is to measure at $I_{DC}=0$ (zero DC field).

Our thermometer will be a **T-type thermocouple**.
We'll use the DMM to measure the EMF (voltage) across it, and from this infer the temperature.

Thermocouples

In a thermocouple, a pair of **dissimilar metal** wires generates a temperature-dependent EMF

Seebeck effect: Charge carriers diffuse away from hot areas, setting up a (material-dependent) EMF that is a function of the temperature gradient.



[Wikipedia](#)

Type	Names of Materials	T Range
B	Platinum 30% Rhodium (+) Platinum 6% Rhodium (-)	2500 -3100F 1370-1700C
C	W5Re Tungsten 5% Rhenium (+) W26Re Tungsten 26% Rhenium (-)	3000-4200F 1650-2315C
E	Chromel (+) Constantan (-)	200-1650F 95-900C
J	Iron (+) Constantan (-)	200-1400F 95-760C
K	Chromel (+) Alumel (-)	200-2300F 95-1260C
N	Nicrosil (+) Nisil (-)	1200-2300F 650-1260C
R	Platinum 13% Rhodium (+) Platinum (-)	1600-2640F 870-1450C
S	Platinum 10% Rhodium (+) Platinum (-)	1800-2640F 980-1450C
T	Copper (+) Constantan (-)	-330-660F -200-350C

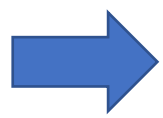
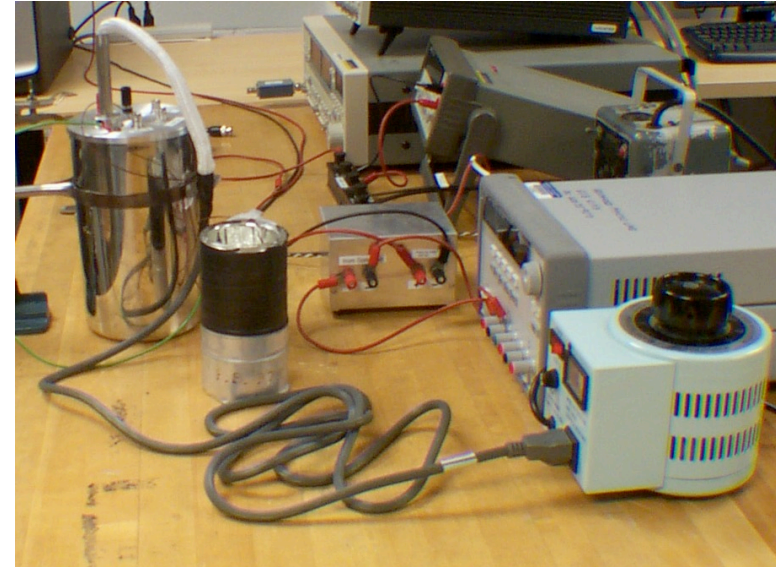
Type T at 0°C: $dV/dT=41.5 \mu\text{V}/^\circ\text{C}$

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Temperature Ramp

Option 1: Manually change the voltage applied to the heater



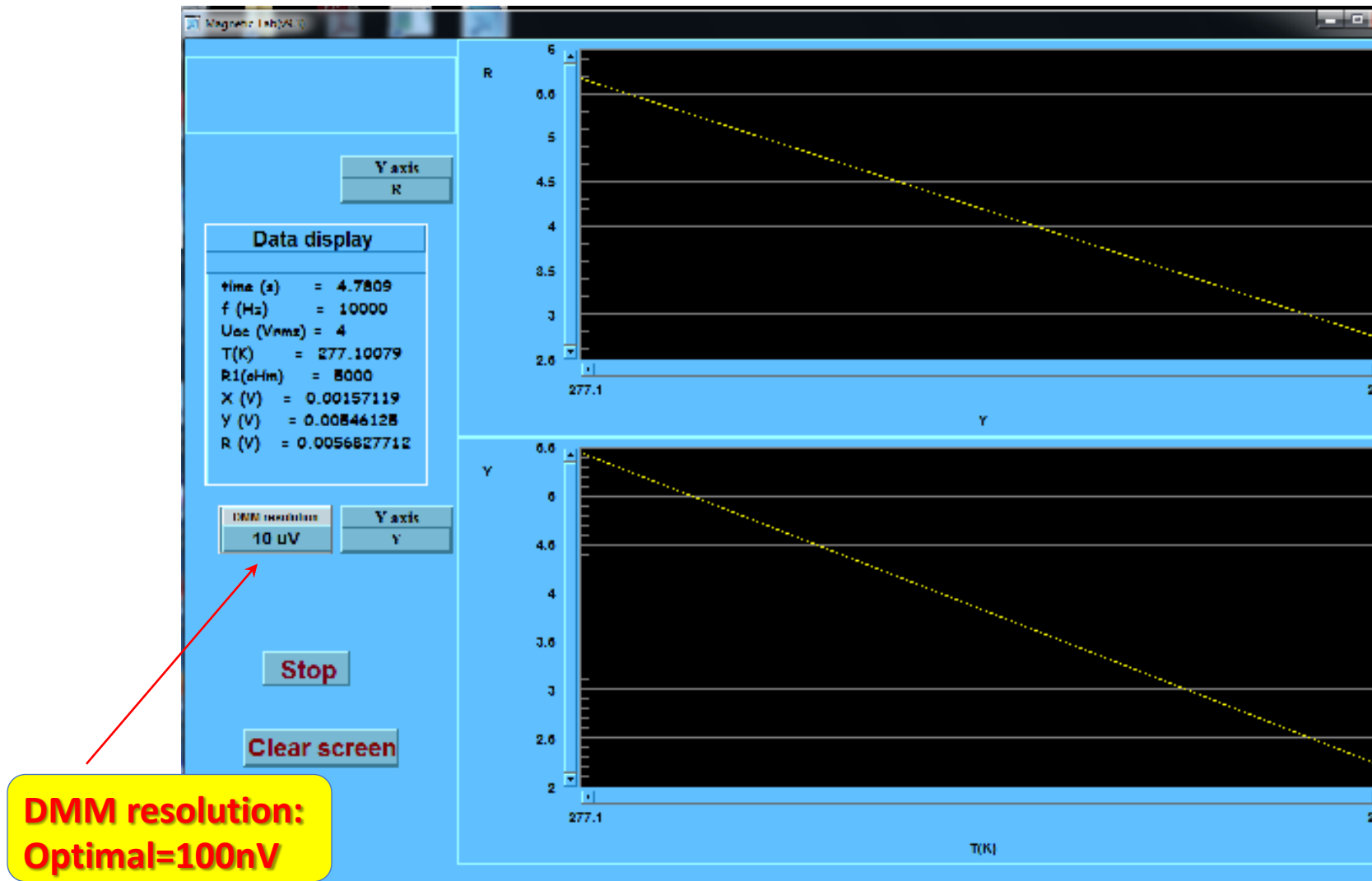
Option 2: Automate using the Omega **PID*** temperature controller



** **P**roportional, **I**ntegral, **D**ifferential: Three standard linear operations to perform on error signal to determine control (feedback) signal to stabilize input at target*

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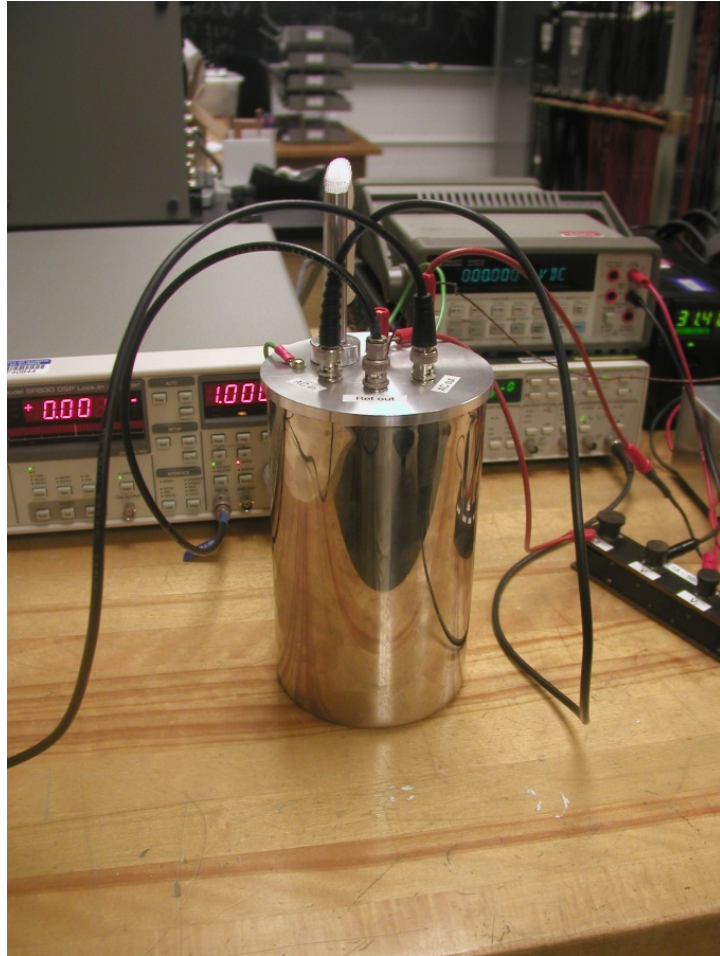
Temperature Measurement



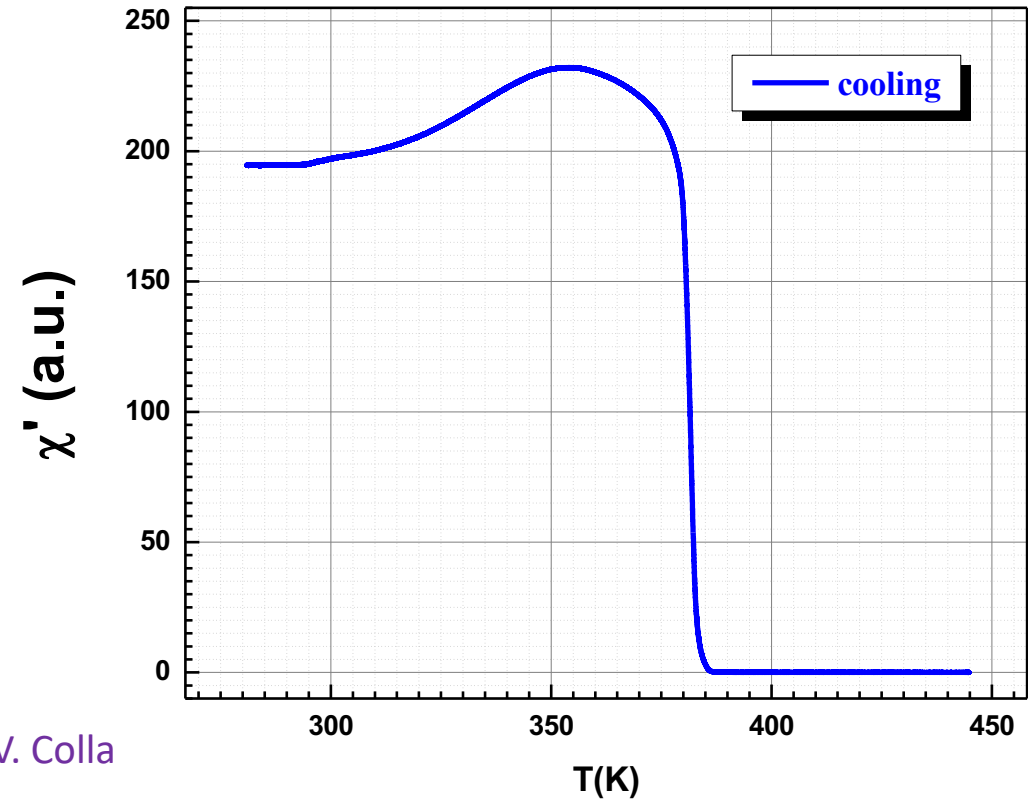
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Temperature Dependence of Permeability



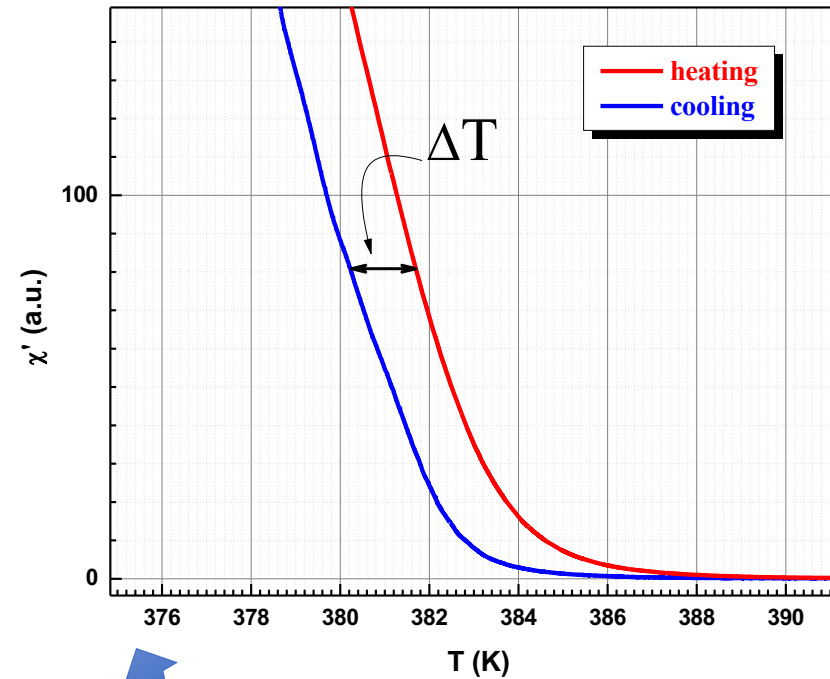
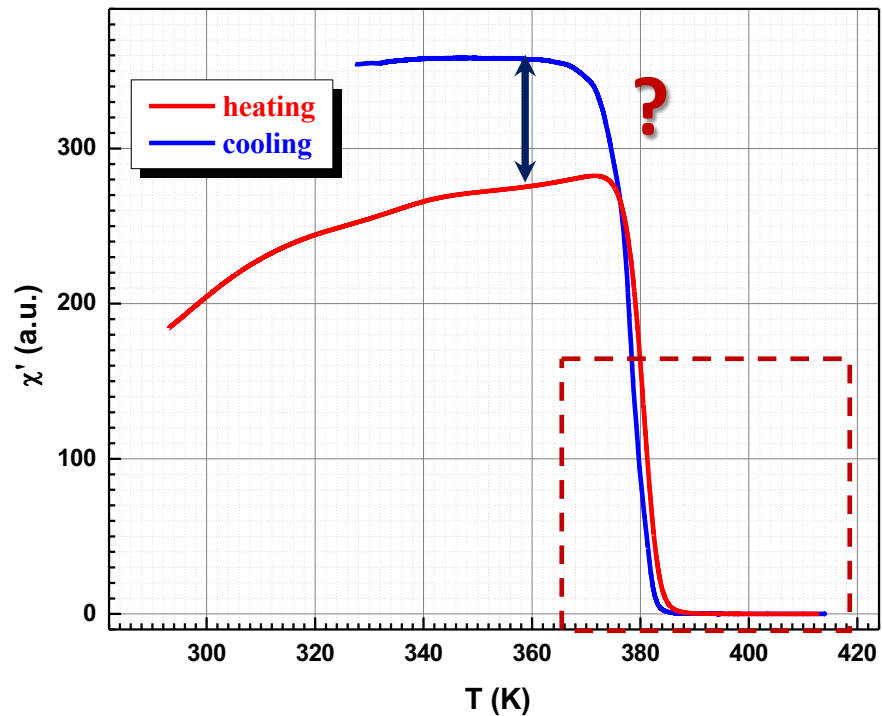
Ferroxcube 3e8



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Temperature Dependence of Permeability

Ferroxcube 4A20

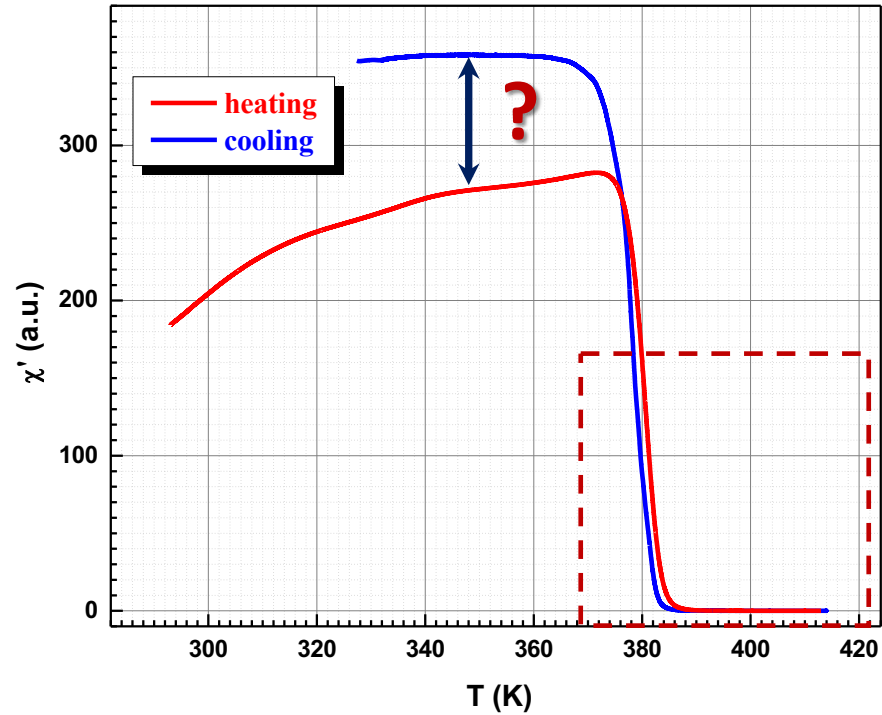


Hysteresis in T_c and χ
What is the source?

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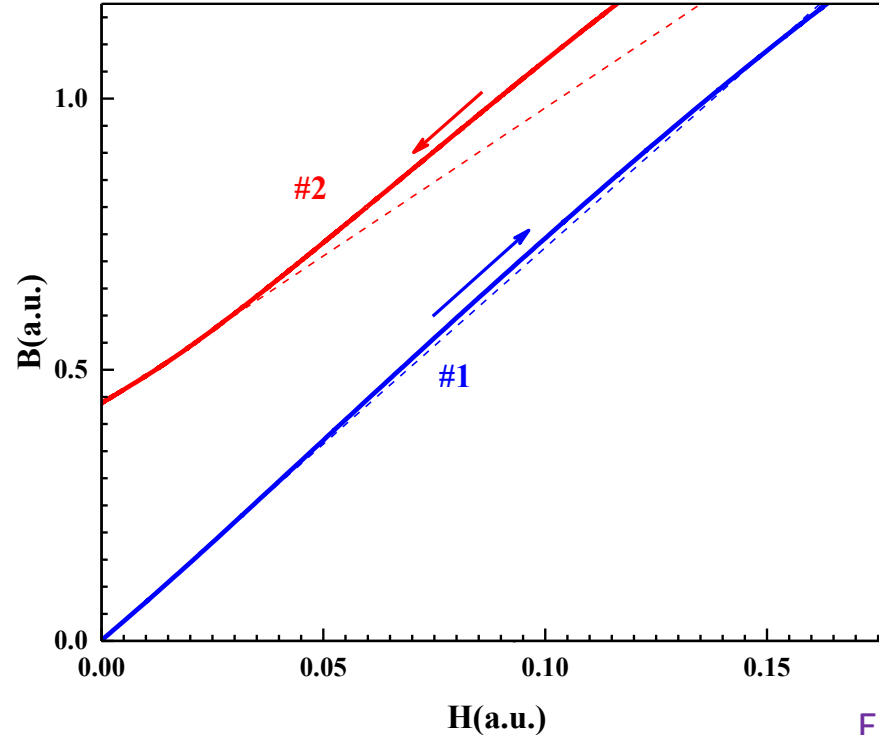


Temperature Dependence of Permeability



Slope#1 ~0.728
Slope#2 ~0.546

$$\mu(T) = \mu_0 \mu_r(T) = \frac{dB}{dH}(T)$$



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Curie-Weiss Law



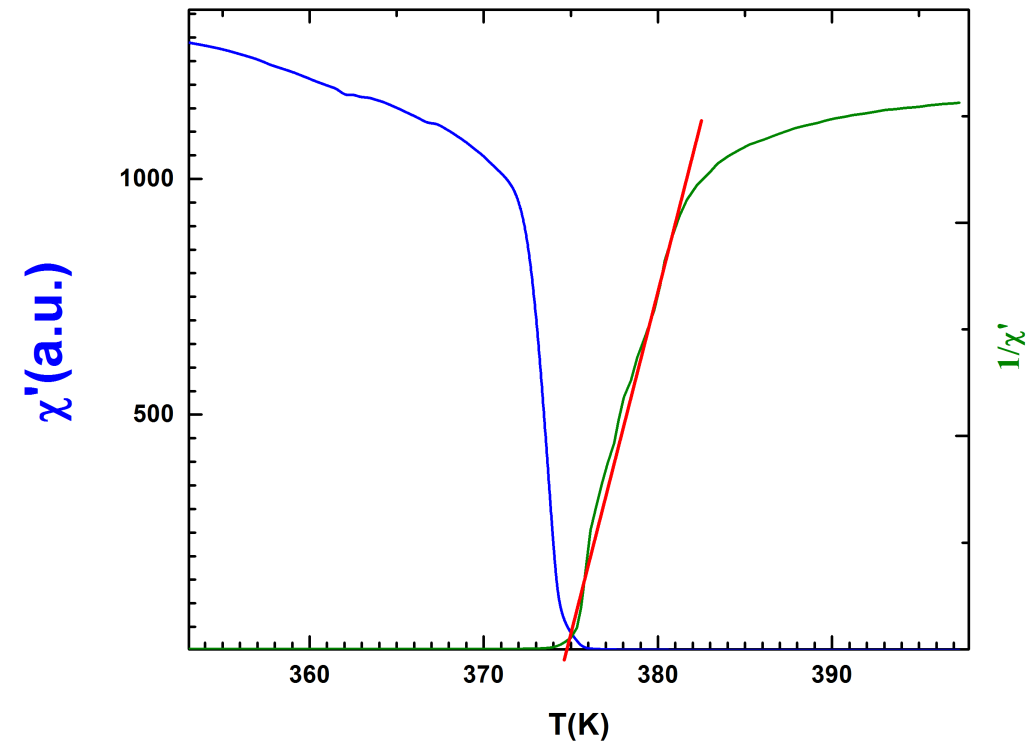
Pierre Curie
1859 – 1906



Pierre Weiss
1865 – 1940

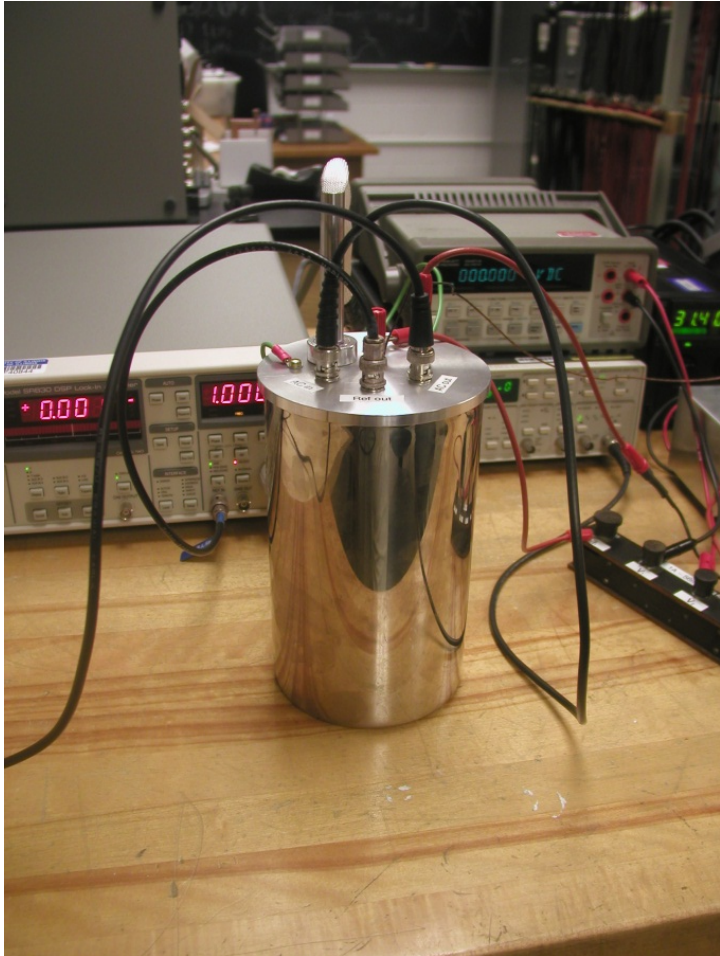
$$\chi' = \frac{C}{T - T_c}$$

3E8 from Ferroxcube

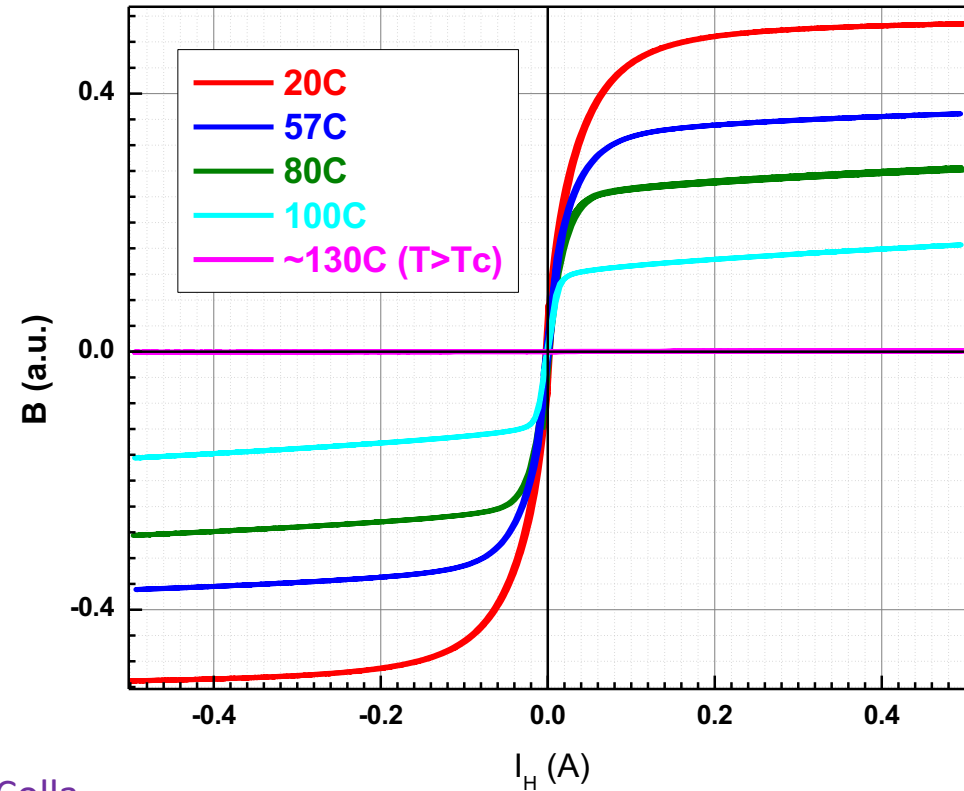


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Temperature Dependence of Permeability



Ferroxcube 3e8



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"Week 3": Full curves at varying H_0

Our (Socially Distanced!) Plan for the Final Lab

- Apparatus is largely prepared, data collection is largely automated
- We usually divide lab tasks up into three weeks:
 1. **B(H)** for **four** samples at **room** temperature (χ , μ'' , *hysteresis*)
 2. **$\chi(H=0)$** for **one** sample at **several** temperatures (*Curie-Weiss law*, T_{Curie})
 3. **B(H,T)** for **one** sample at **several** temperatures (*saturation, remanence, ...*)
- We will post videos and data on the following schedule:
 - **Week 1**: Wednesday 4/22
 - **Week 2/3**: Sunday 4/26 (*combined, same sample*)
- You will perform the data analysis and write a report as usual, following the lecture and write-up
- *Contact me and your TAs with questions! Don't wait!*

References

- Information about magnetic materials can be found in:
[\\engr-file-03\phyinst\APL Courses\PHYCS401\Experiments\AC_Magnetization\Magnetic Materials](#)
- SR830 (Lock-in Amplifier) manual
[\\engr-file-03\phyinst\APL Courses\PHYCS401\Common\EquipmentManuals\SR830m.pdf](#)

End of Semester

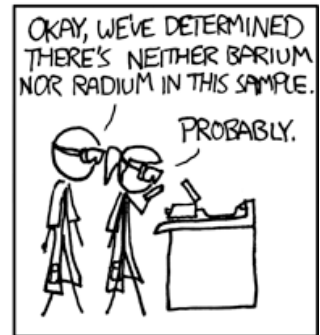
- **No lecture** 4/27 or 5/4 – *this is our last meeting!*
- Please turn in any late **lab reports!** You can revise and resubmit one report for a **replacement grade**. *Coordinate with your TA!*
- **Sunday, May 10th, 11:59 pm** is the deadline for the final lab, *and* for all other reports/resubmissions. *No extensions and no vouchers!*
- **Grade thresholds** may be adjusted if needed on a section-by-section basis if substantial disparities arise



MOVIE SCIENCE MONTAGE



ACTUAL SCIENCE MONTAGE



C

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Thank You

for your hard work

and especially for bearing with us during this most unusual and challenging semester

Best of luck in your future endeavors, in and out of the laboratory!

I ILLINOIS