

## Standing Wave – Simple Theory

In this example, we show plots of acoustic quantities associated with the superposition of two counter-propagating 1-D monochromatic plane traveling waves – *i.e.* a simple standing wave – in “free air”. Please refer to Physics 406 Lecture Notes 12 p. 6-18 for details.

The individual right and left-traveling complex **time-domain** over-pressure amplitudes are:  
 $\tilde{p}_A(x,t) = \tilde{A}e^{i(\omega t-kx)}$  and  $\tilde{p}_B(x,t) = \tilde{B}e^{i(\omega t+kx)}$  with  $\tilde{A} \neq \tilde{B}$  {necessarily}, where  $\tilde{A} = |\tilde{A}|e^{i\varphi_A^o} \equiv Ae^{i\varphi_A^o}$   
and  $\tilde{B} = |\tilde{B}|e^{i\varphi_B^o} \equiv Be^{i\varphi_B^o}$ . Thus:  $\tilde{p}_{tot}(x,t) = \tilde{p}_A(x,t) + \tilde{p}_B(x,t) = \tilde{A}e^{i(\omega t-kx)} + \tilde{B}e^{i(\omega t+kx)}$  (*Pascals*).

The individual right and left-traveling complex **time-domain** 1-D particle velocity amplitudes are:  $\tilde{u}_A^{\parallel}(x,t) = \frac{\tilde{A}}{\rho_o c} e^{i(\omega t-kx)} \equiv \tilde{u}_{A_o}^{\parallel} e^{i(\omega t-kx)}$  and:  $\tilde{u}_B^{\parallel}(x,t) = -\frac{\tilde{B}}{\rho_o c} e^{i(\omega t+kx)} \equiv -\tilde{u}_{B_o}^{\parallel} e^{i(\omega t+kx)}$  (using  $c = \omega/k$  ).

$$\text{Thus: } \tilde{u}_{tot}^{\parallel}(x,t) = \tilde{u}_A^{\parallel}(x,t) + \tilde{u}_B^{\parallel}(x,t) = \tilde{u}_{A_o}^{\parallel} e^{i(\omega t-kx)} + \tilde{u}_{B_o}^{\parallel} e^{i(\omega t+kx)} = \frac{\tilde{A}}{\rho_o c} e^{i(\omega t-kx)} - \frac{\tilde{B}}{\rho_o c} e^{i(\omega t+kx)} (\text{m/s}).$$

Defining:  $\tilde{R} \equiv \frac{\tilde{B}}{\tilde{A}} = \frac{|\tilde{B}|e^{i\varphi_B^o}}{|\tilde{A}|e^{i\varphi_A^o}} = \frac{|\tilde{B}|}{|\tilde{A}|} e^{i(\varphi_B^o - \varphi_A^o)} = |\tilde{R}| e^{i(\varphi_B^o - \varphi_A^o)} = |\tilde{R}| e^{i\Delta\varphi_{BA}^o}$  where:  $\Delta\varphi_{BA}^o \equiv \varphi_B^o - \varphi_A^o$ , then:

$$\tilde{p}_{tot}(x,t) = \tilde{A} \left[ e^{i(\omega t-kx)} + |\tilde{R}| e^{i(\omega t+kx)} \cdot e^{i\Delta\varphi_{BA}^o} \right] = \tilde{A} e^{i(\omega t-kx)} \left[ 1 + |\tilde{R}| e^{i(2kx+\Delta\varphi_{BA}^o)} \right]$$

and:

$$\tilde{u}_{tot}^{\parallel}(x,t) = \frac{\tilde{A}}{\rho_o c} \left[ e^{i(\omega t-kx)} - |\tilde{R}| e^{i(\omega t+kx)} \cdot e^{i\Delta\varphi_{BA}^o} \right] = \frac{\tilde{A}}{\rho_o c} e^{i(\omega t-kx)} \left[ 1 - |\tilde{R}| e^{i(2kx+\Delta\varphi_{BA}^o)} \right]$$

The **magnitudes** of the complex total/resultant over-pressure  $|\tilde{p}_{tot}(x)|$  and longitudinal particle velocity  $|\tilde{u}_{tot}^{\parallel}(x)|$  are:

$$|\tilde{p}_{tot}(x)| = |\tilde{A}| \sqrt{1 + 2|\tilde{R}|\cos(2kx+\Delta\varphi_{BA}^o) + |\tilde{R}|^2} \quad \text{and: } |\tilde{u}_{tot}^{\parallel}(x)| = \frac{|\tilde{A}|}{\rho_o c} \sqrt{1 - 2|\tilde{R}|\cos(2kx+\Delta\varphi_{BA}^o) + |\tilde{R}|^2}$$

The **phases** associated with the complex total/resultant over-pressure  $\varphi_{p_{tot}}(x)$  and longitudinal particle velocity  $\varphi_{u_{tot}}$  are:

$$\varphi_{p_{tot}}(x) = \tan^{-1} \left[ \frac{\sin kx(1 + |\tilde{R}|\cos(2kx+\Delta\varphi_{BA}^o)) + |\tilde{R}|\cos kx \cdot \sin(2kx+\Delta\varphi_{BA}^o)}{\cos kx(1 + |\tilde{R}|\cos(2kx+\Delta\varphi_{BA}^o)) - |\tilde{R}|\sin kx \cdot \sin(2kx+\Delta\varphi_{BA}^o)} \right]$$

$$\varphi_{u_{tot}}(x) = \tan^{-1} \left[ \frac{\sin kx(1 - |\tilde{R}|\cos(2kx+\Delta\varphi_{BA}^o)) - |\tilde{R}|\cos kx \cdot \sin(2kx+\Delta\varphi_{BA}^o)}{\cos kx(1 - |\tilde{R}|\cos(2kx+\Delta\varphi_{BA}^o)) + |\tilde{R}|\sin kx \cdot \sin(2kx+\Delta\varphi_{BA}^o)} \right]$$

The complex **frequency-domain** longitudinal specific acoustic impedance, it's magnitude and phase are, noting {[here](#)} that  $\omega = ck$ :

$$\tilde{z}_{a_{tot}}^{\parallel}(x, \omega) = z_o \frac{\left[ \begin{array}{l} \left\{ 1 - |\tilde{R}|^2 \right\} + 2i|\tilde{R}| \sin(2kx + \Delta\phi_{BA}^o) \\ \left\{ 1 + |\tilde{R}|^2 \right\} - 2|\tilde{R}| \cos(2kx + \Delta\phi_{BA}^o) \end{array} \right]}{\left[ \begin{array}{l} \left\{ 1 + |\tilde{R}|^2 \right\} - 2|\tilde{R}| \cos(2kx + \Delta\phi_{BA}^o) \end{array} \right]}$$

$$|\tilde{z}_{a_{tot}}^{\parallel}(x, \omega)| = z_o \frac{\sqrt{\left\{ 1 - |\tilde{R}|^2 \right\}^2 + 4|\tilde{R}|^2 \sin^2(2kx + \Delta\phi_{BA}^o)}}{\left[ \begin{array}{l} \left\{ 1 + |\tilde{R}|^2 \right\} - 2|\tilde{R}| \cos(2kx + \Delta\phi_{BA}^o) \end{array} \right]}$$

$$\varphi_{z_{a_{tot}}}(x, \omega) = \tan^{-1} \left( \frac{2|\tilde{R}| \sin(2kx + \Delta\phi_{BA}^o)}{\left\{ 1 - |\tilde{R}|^2 \right\}} \right) = \Delta\varphi_{p_{tot} - u_{tot}^{\parallel}}(x, \omega) = \varphi_{p_{tot}}(x, \omega) - \varphi_{u_{tot}^{\parallel}}(x, \omega)$$

The complex **frequency-domain** longitudinal acoustic intensity, it's magnitude and phase are:

$$\tilde{I}_{a_{tot}}^{\parallel}(x, \omega) = \frac{1}{2} \frac{|\tilde{A}|^2}{z_o} \left[ \begin{array}{l} \left\{ 1 - |\tilde{R}|^2 \right\} + 2i|\tilde{R}| \sin(2kx + \Delta\phi_{BA}^o) \end{array} \right]$$

$$|\tilde{I}_{a_{tot}}^{\parallel}(x, \omega)| = \frac{1}{2} \frac{|\tilde{A}|^2}{z_o} \sqrt{\left\{ 1 - |\tilde{R}|^2 \right\}^2 + 4|\tilde{R}|^2 \sin^2(2kx + \Delta\phi_{BA}^o)}$$

$$\varphi_{I_{a_{tot}}}(x, \omega) = \tan^{-1} \left( \frac{2|\tilde{R}| \sin(2kx + \Delta\phi_{BA}^o)}{\left\{ 1 - |\tilde{R}|^2 \right\}} \right) = \varphi_{z_{a_{tot}}}(x, \omega) = \Delta\varphi_{p_{tot} - u_{tot}^{\parallel}}(x, \omega) = \varphi_{p_{tot}}(x, \omega) - \varphi_{u_{tot}^{\parallel}}(x, \omega)$$

The complex **frequency-domain** potential, kinetic and total energy densities associated with counter-propagating 1-D monochromatic traveling plane waves (*n.b.* all **purely real** quantities):

$$w_{potl}(x, \omega) \equiv \frac{1}{4} \frac{1}{\rho_o c^2} |\tilde{p}_{tot}(x, \omega)|^2 = \frac{1}{4} \frac{|\tilde{A}|^2}{\rho_o c^2} \left[ 1 + |\tilde{R}|^2 + 2|\tilde{R}| \cos(2kx + \Delta\phi_{BA}^o) \right]$$

$$w_{kin}(x, \omega) \equiv \frac{1}{4} \rho_o |\tilde{u}_{tot}^{\parallel}(x, \omega)|^2 = \frac{1}{4} \frac{|\tilde{A}|^2}{\rho_o c^2} \left[ 1 + |\tilde{R}|^2 - 2|\tilde{R}| \cos(2kx + \Delta\phi_{BA}^o) \right]$$

$$w_{tot}(x, \omega) \equiv w_{potl}(x, \omega) + w_{kin}(x, \omega) = \frac{1}{2} \frac{|\tilde{A}|^2}{\rho_o c^2} \left[ 1 + |\tilde{R}|^2 \right] = \frac{1}{2} \frac{|\tilde{A}|^2}{z_o c} \left[ 1 + |\tilde{R}|^2 \right]$$

We coded up the above formulas using MATLAB and show plots in the figures below of the magnitudes and phases of complex total over-pressure, particle velocity, longitudinal specific acoustic impedance and longitudinal acoustic intensity  $|\tilde{p}_{tot}(\theta)|$ ,  $|\tilde{u}_{tot}^{\parallel}(\theta)|$ ,  $|\tilde{z}_{a_{tot}}^{\parallel}(\theta)|$  and  $|\tilde{I}_{a_{tot}}^{\parallel}(\theta)|$  vs.  $\theta \equiv kx$  for  $0 \leq |\tilde{R}| \leq 1$  in steps of 0.1, with  $\Delta\phi_{BA}^o = 0.0$ . Please note the suppressed zeroes in some of the following plots!

Figure 1.  $|\tilde{R}| = 0.0$

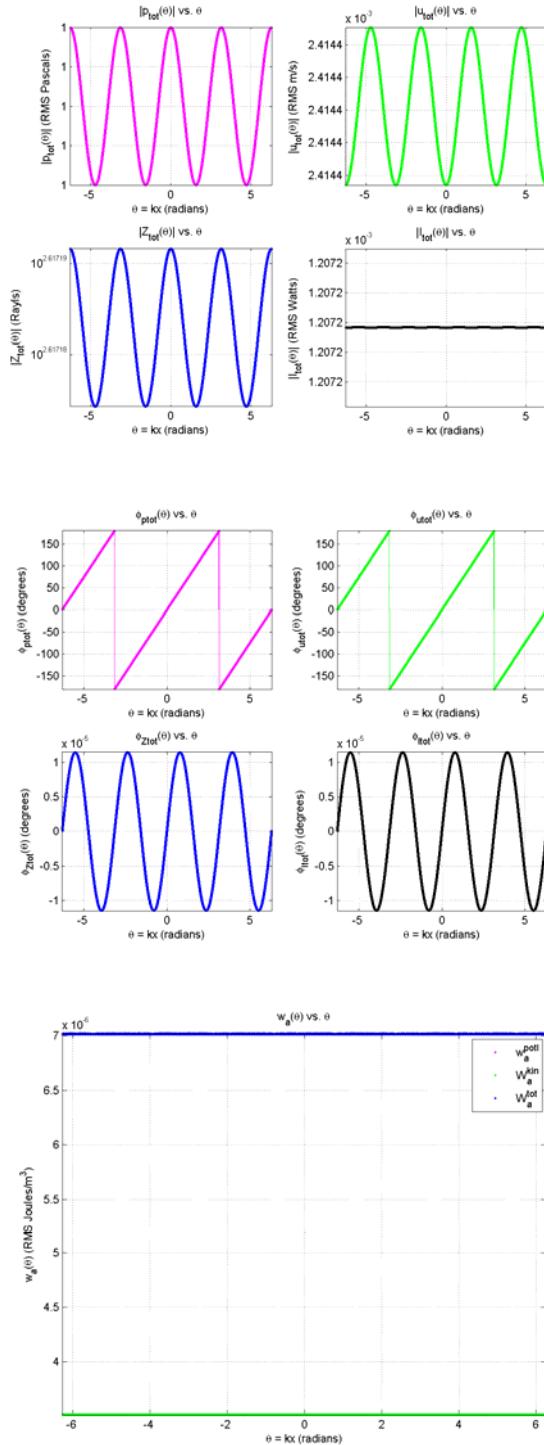


Figure 2.  $|\tilde{R}| = 0.1$

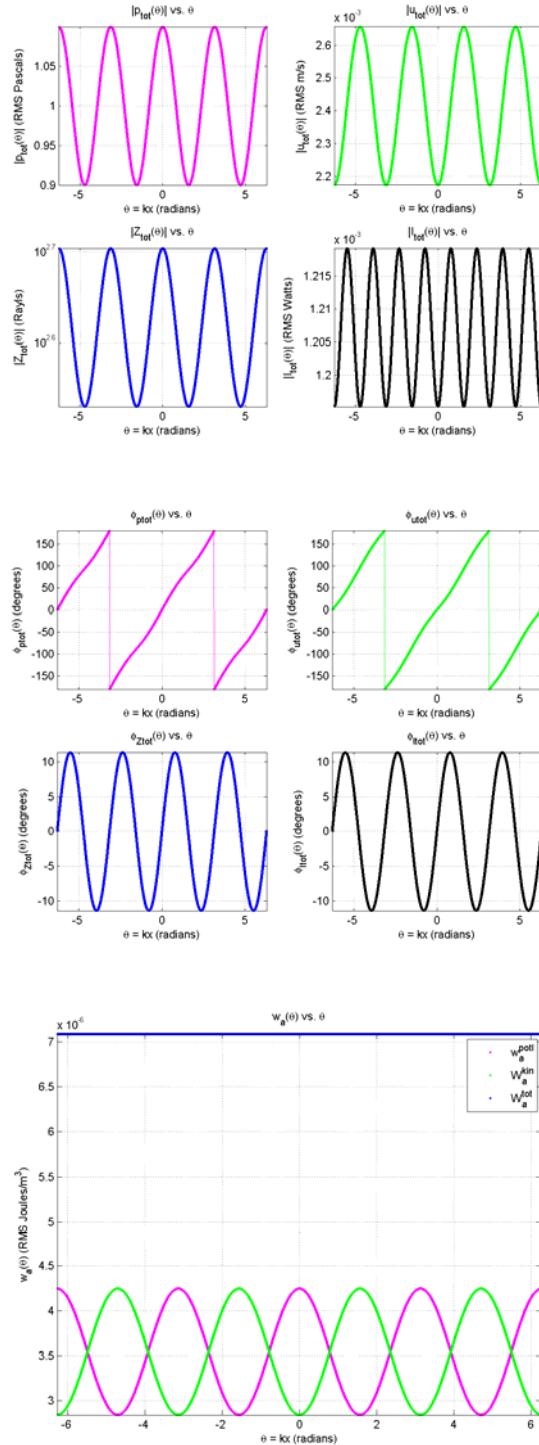


Figure 3.  $|\tilde{R}| = 0.2$

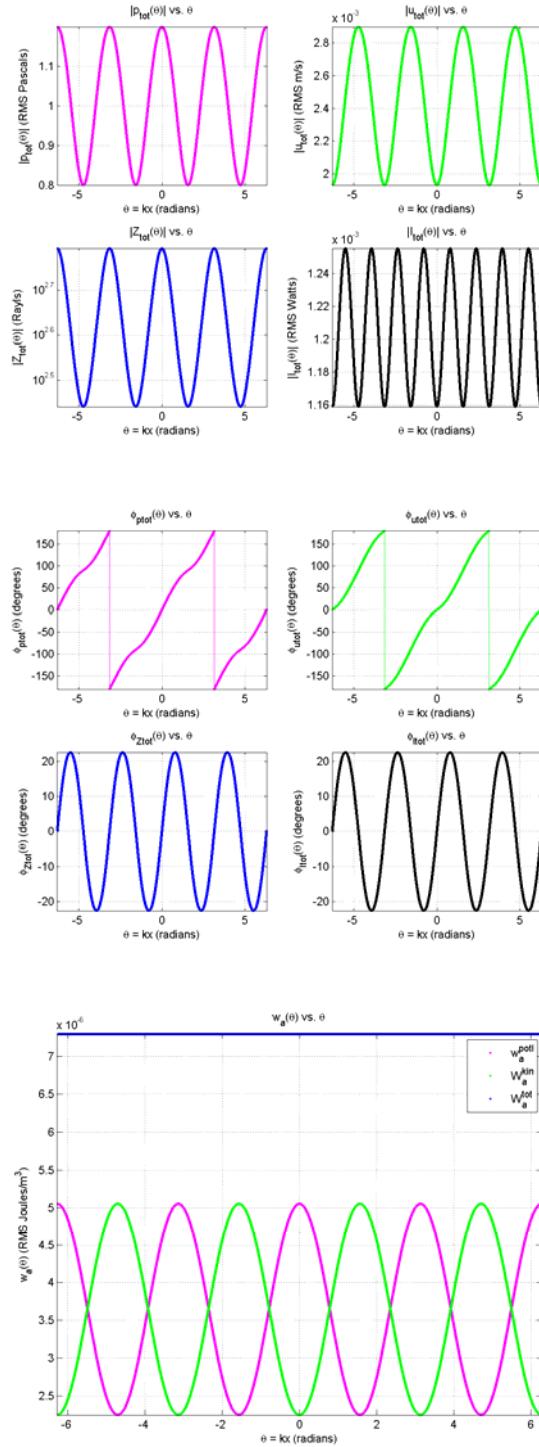


Figure 4.  $|\tilde{R}| = 0.3$

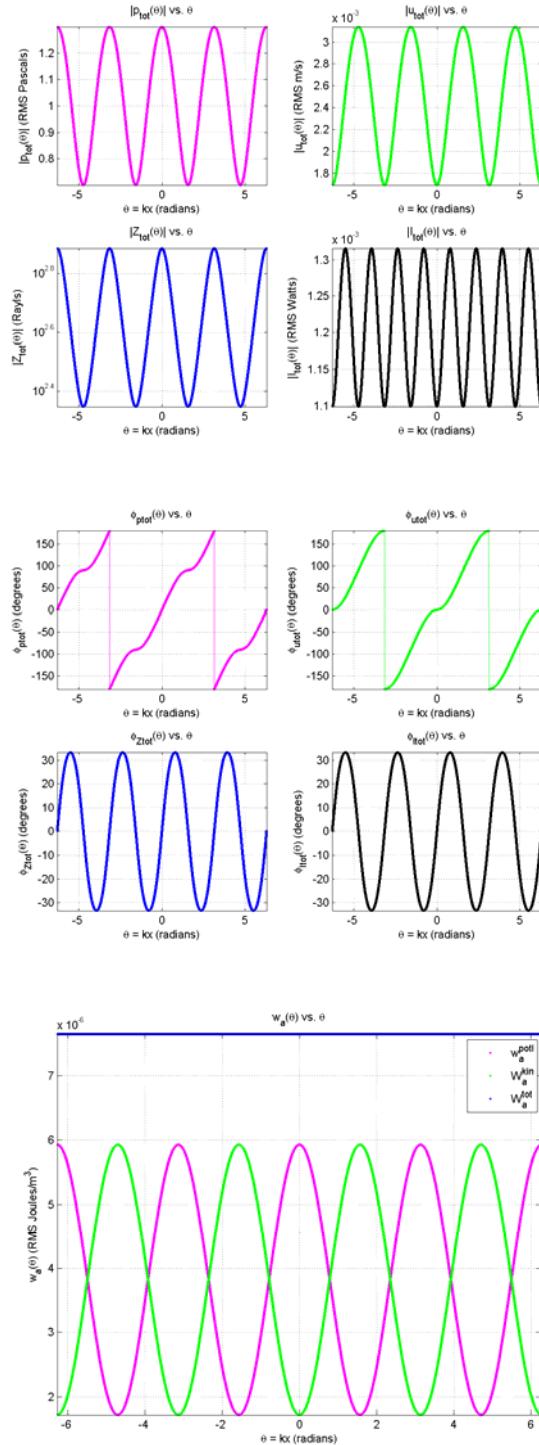


Figure 5.  $|\tilde{R}| = 0.4$

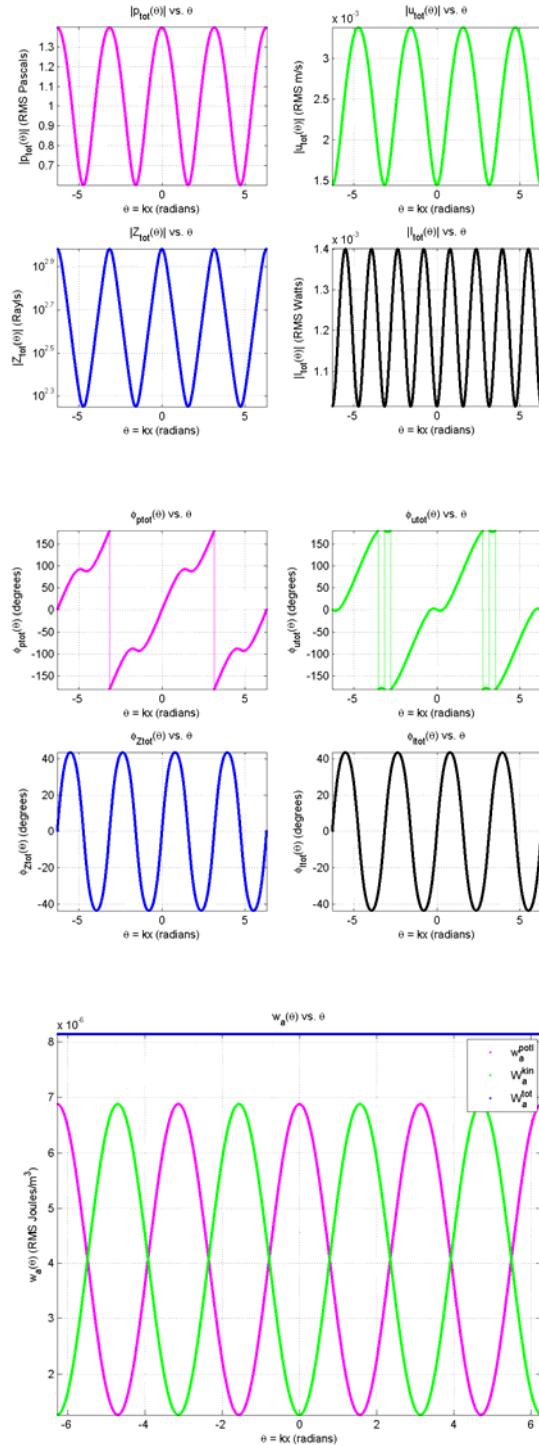


Figure 6.  $|\tilde{R}| = 0.5$

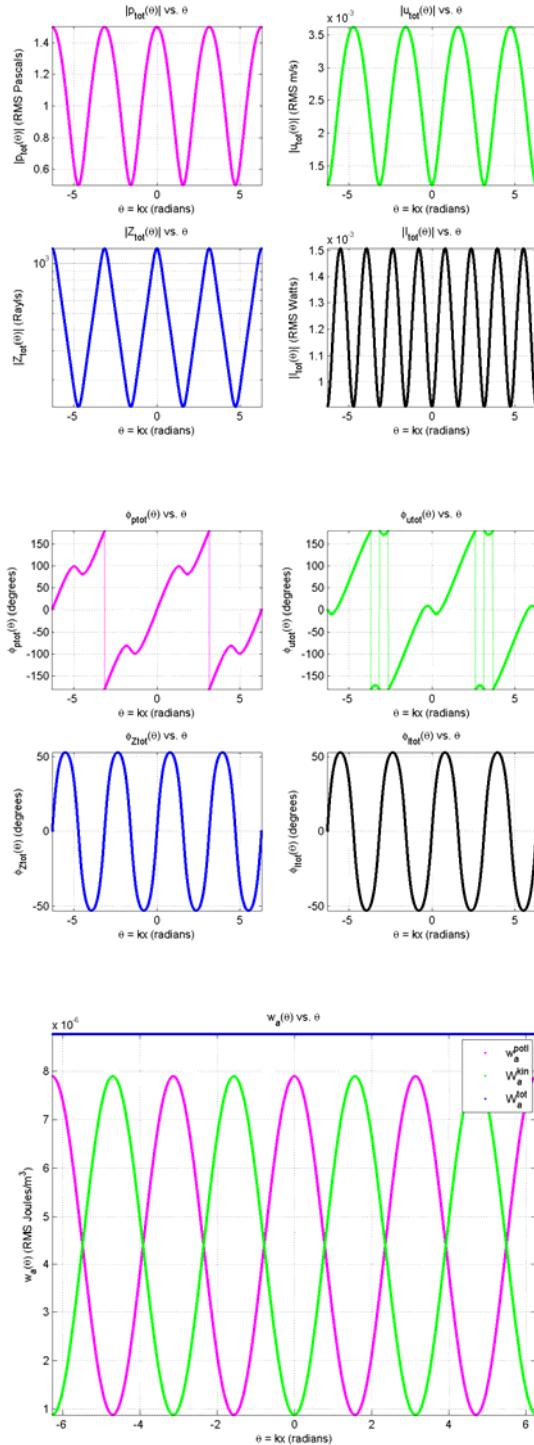


Figure 7.  $|\tilde{R}| = 0.6$

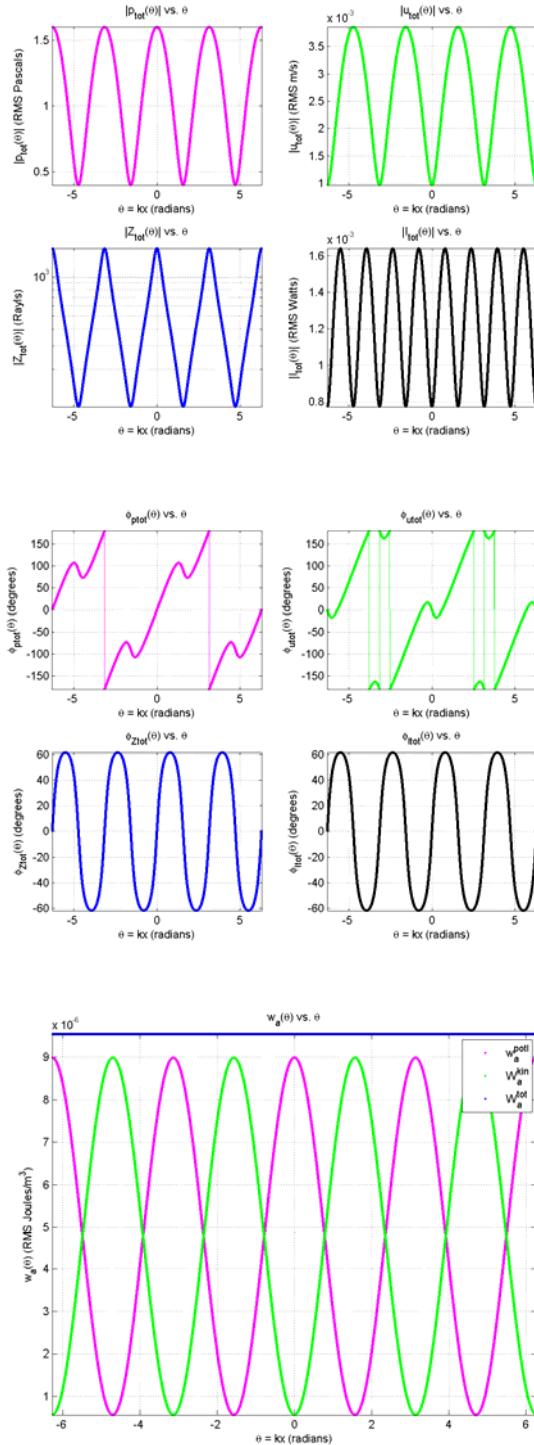


Figure 8.  $|\tilde{R}| = 0.7$

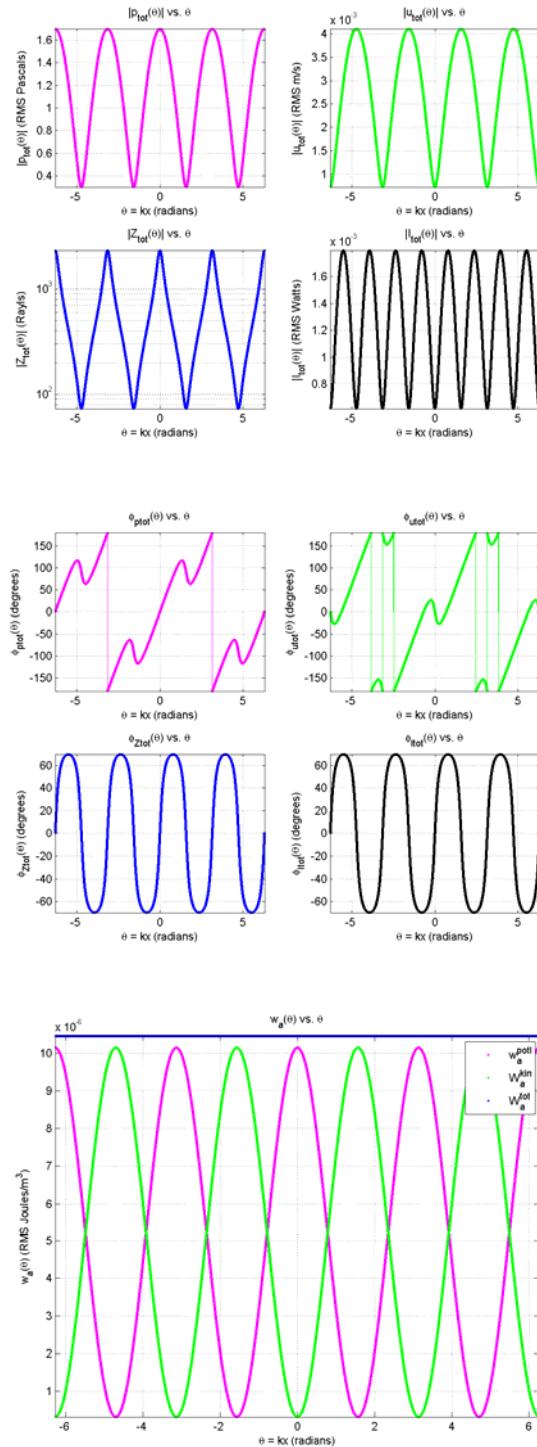


Figure 9.  $|\tilde{R}| = 0.8$

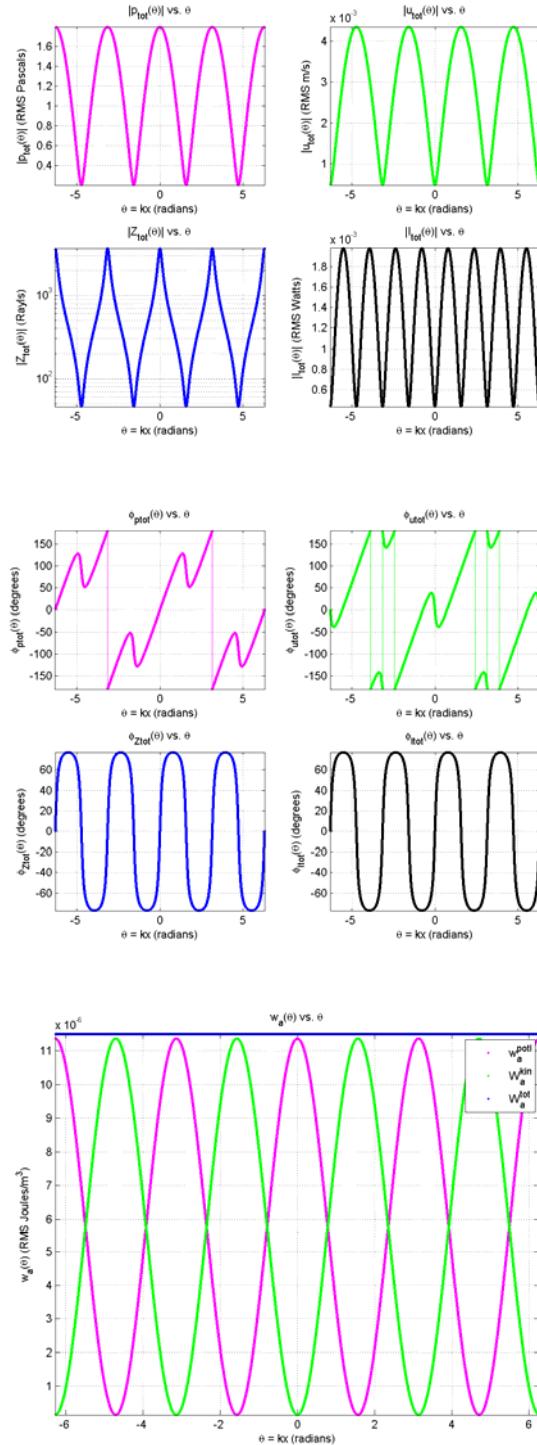


Figure 10.  $|\tilde{R}| = 0.9$

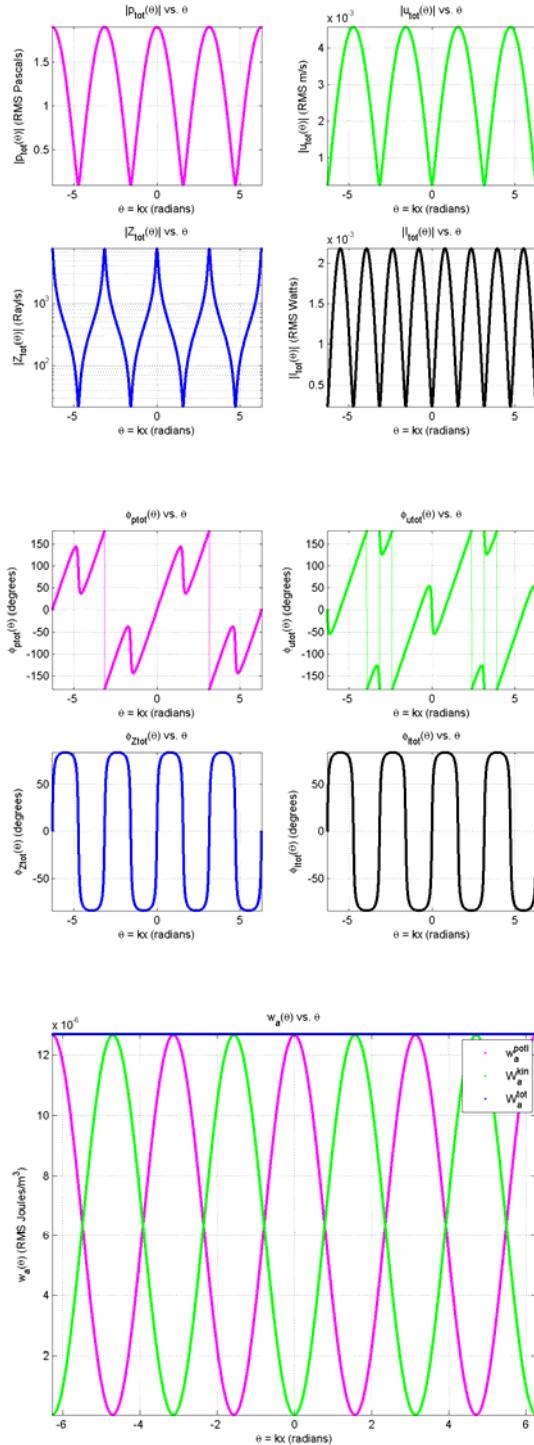
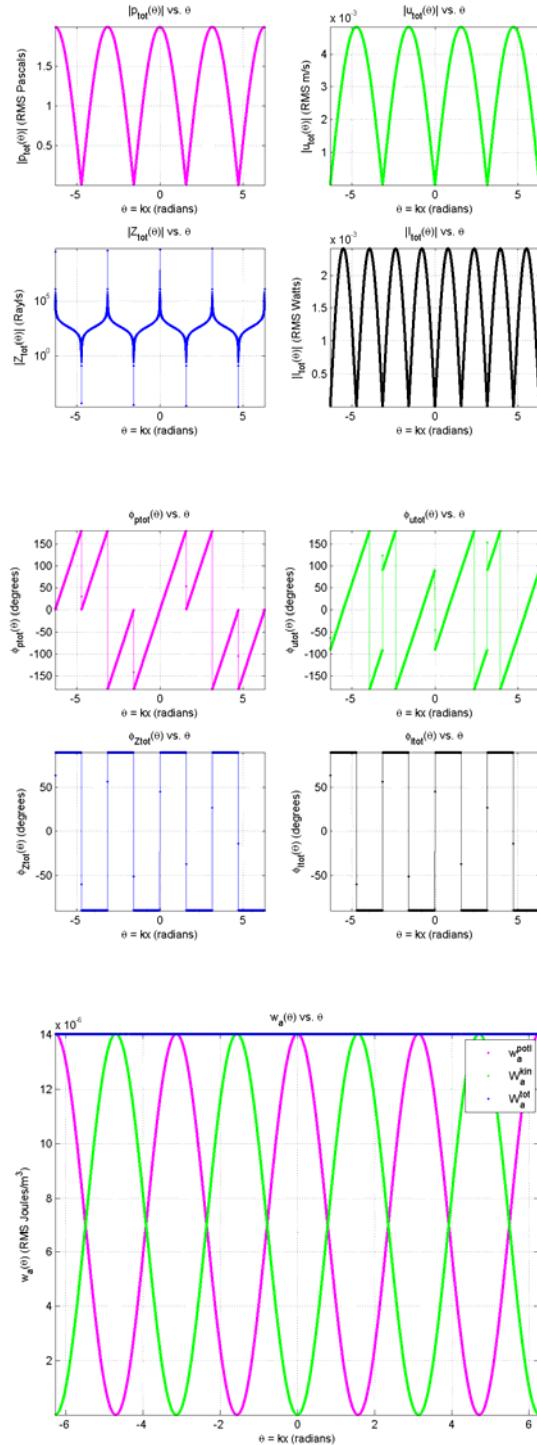


Figure 11.  $|\tilde{R}| = 1.0$



## Listing of the MATLAB code:

```
%=====
% Simple_Standing_Wave.m
%
% Study of the acoustical physics associated with the linear superposition
% of two counter-propagating 1-D monochromatic plane/traveling waves,
% = a standing wave.
%
%=====
% Author: Prof. Steven Errede 4/5/2014 07:50 hr
%
% Last Update: 4/7/2014 10:20 hr {SME}
%=====
clear all;
close all;

npts = 40000;

thetar = zeros(1,npts);

MgPtot = zeros(1,npts);
MgUtot = zeros(1,npts);
MgZtot = zeros(1,npts);
MgItot = zeros(1,npts);

PhPtot = zeros(1,npts);
PhUtot = zeros(1,npts);
PhZtot = zeros(1,npts);
PhItot = zeros(1,npts);

Wapot = zeros(1,npts);
Wakin = zeros(1,npts);
Watot = zeros(1,npts);

rho0 = 1.204; % Density of air @ NTP (kg/m^3)
c = 344.0; % Longitudinal speed of sound @ NTP (m/s)
z0 = rho0*c; % Longitudinal specific acoustic impedance of free air (Rayls)

p0 = 1.000; % Over-pressure amplitude (RMS Pascals)

% |~R| = |~B|/|~A|
%MgR = 1.000 - 1.0e-7;
MgR = 0.100;
%MgR = 0.000 + 1.0e-7;

delphio_ba = 0.0; % = phi0b - phi0a (degrees)
delphio_bar = (pi/180.0)*delphio_ba; % (radians)

thetar_lo = -2.0*pi + 1.0e-7; % = kx_lo
thetar_hi = 2.0*pi; % = kx_hi

dthetar = (thetar_hi - thetar_lo)/npts;

for j = 1:npts;
    thetar(j) = thetar_lo + (j-1)*dthetar; % = kx

    MgPtot(j) = p0 *sqrt(1.0 + (2.0*MgR*cos(2.0*thetar(j) + delphio_bar)) + (MgR*MgR));
    MgUtot(j) = (p0/(rho0*c))*sqrt(1.0 - (2.0*MgR*cos(2.0*thetar(j) + delphio_bar)) + (MgR*MgR));
    MgZtot(j) = MgPtot(j)/MgUtot(j);
    MgItot(j) = MgPtot(j)*MgUtot(j)/2.0;

    Pnum = sin(thetar(j))*(1.0 + (MgR*cos(2.0*thetar(j) + delphio_bar))) +
           cos(thetar(j))*(-MgR*sin(2.0*thetar(j) + delphio_bar));
    Pden = cos(thetar(j))*(1.0 + (MgR*cos(2.0*thetar(j) + delphio_bar))) -
           sin(thetar(j))*(-MgR*sin(2.0*thetar(j) + delphio_bar));
    PhPtot(j) = (180.0/pi)*atan2(Pnum,Pden);

    Unum = sin(thetar(j))*(1.0 - (MgR*cos(2.0*thetar(j) + delphio_bar))) -
           cos(thetar(j))*(-MgR*sin(2.0*thetar(j) + delphio_bar));

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Uden      = cos(theta(j))*(1.0 - (MgR*cos(2.0*theta(j) + delphi0_bar))) +
            sin(theta(j))*(MgR*sin(2.0*theta(j) + delphi0_bar));
PhUtot(j) = (180.0/pi)*atan2(Unum,Uden);

Znum      = (2.0*MgR*sin(2.0*theta(j) + delphi0_bar));
Zden      = (1.0 - (MgR*MgR));
PhZtot(j) = (180.0/pi)*atan2(Znum,Zden);

Inum      = (2.0*MgR*sin(2.0*theta(j) + delphi0_bar));
Iden      = (1.0 - (MgR*MgR));
PhItot(j) = (180.0/pi)*atan2(Inum,Iden);

Wapot(j) = 0.5 * (MgPtot(j)*MgPtot(j))/(z0*c);
Wakin(j) = 0.5*rho0*(MgUtot(j)*MgUtot(j));
Watot(j) = Wapot(j) + Wakin(j);
end

figure (01);
subplot(2,2,1);
plot (theta,MgPtot,'m',theta,MgPtot,'m.');
axis tight;
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('|p_{tot}({\theta})| (RMS Pascals)');
title ('|p_{tot}({\theta})| vs. {\theta}');

subplot(2,2,2);
plot (theta,MgUtot,'g.',theta,MgUtot,'g.');
axis tight;
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('|u_{tot}({\theta})| (RMS m/s)');
title ('|u_{tot}({\theta})| vs. {\theta}');

subplot(2,2,3);
%plot (theta,MgZtot,'b',theta,MgZtot,'b.');
semilogy(theta,MgZtot,'b',theta,MgZtot,'b.');
axis tight;
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('|Z_{tot}({\theta})| (Rayls)');
title ('|Z_{tot}({\theta})| vs. {\theta}');

subplot(2,2,4);
plot (theta,MgItot,'k',theta,MgItot,'k.');
axis tight;
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('|I_{tot}({\theta})| (RMS Watts)');
title ('|I_{tot}({\theta})| vs. {\theta}');

figure (02);
subplot(2,2,1);
plot (theta,PhPtot,'m',theta,PhPtot,'m.');
axis tight;
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('|\phi|_{ptot}({\theta}) (degrees)');
title ('|\phi|_{ptot}({\theta}) vs. {\theta}');

subplot(2,2,2);
plot (theta,PhUtot,'g',theta,PhUtot,'g.');
axis tight;
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('|\phi|_{utot}({\theta}) (degrees)');
title ('|\phi|_{utot}({\theta}) vs. {\theta}');
subplot(2,2,3);
plot (theta,PhZtot,'b',theta,PhZtot,'b.');
%semilogy(theta,PhZtot,'b',theta,PhZtot,'b.');
axis tight;

```

```
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('{\phi}_Ztot({\theta}) (degrees)');
title ('{\phi}_Ztot({\theta}) vs. {\theta}');

subplot(2,2,4);
plot (thetar,PhItot,'k',thetar,PhItot,'k.');
axis tight;
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('{\phi}_Itot({\theta}) (degrees)');
title ('{\phi}_Itot({\theta}) vs. {\theta}');

figure (03);
plot (thetar,Wapot,'m.',thetar,Wakin,'g.',thetar,Watot,'b.');
axis tight;
grid on;
xlabel ('{\theta} = kx (radians)');
ylabel ('w_a({\theta}) (RMS Joules/m^3)');
title ('w_a({\theta}) vs. {\theta}');
legend ('w_a^{potl}', 'w_a^{kin}', 'w_a^{tot}'');
```