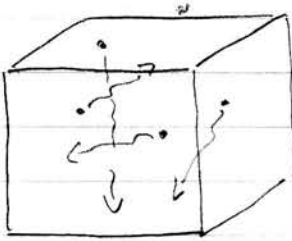


(1)

Lecture 9

Last time we showed that photon gas at temperature T inside a cavity of volume V has spectrum:



$$u(\omega) = \frac{\pi \omega^3}{\pi^2 c^3} \frac{1}{e^{\beta \hbar \omega} - 1}$$

$$\text{and } u_{\text{tot}} = \frac{U}{V} = \int_0^\infty d\omega u(\omega) = \frac{\pi^2}{15} \left(\frac{k_B T}{\hbar^3 c^3} \right)^4$$

$$U = \alpha V T^4 \quad (\text{notice } U \sim V \text{ since it's an extensive quantity})$$

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = 4\alpha V T^3 \quad \text{heat capacity} \sim T^3$$

$$\text{For fixed cavity } V \quad dU = T dS \quad \text{so}$$

$$dS = \frac{dU}{T} = 4\alpha V T^2 dT$$

$$S = \int \frac{dU}{T} = \frac{4}{3} \alpha V T^3 \quad (\text{const. of integration} = 0 \text{ so } S(0) = 0)$$

$$\therefore F = U - TS = V(\alpha T^4 - \frac{4}{3} \alpha T^4) = -\frac{1}{3} \alpha V T^4$$

$$\text{pressure } p = -\left(\frac{\partial F}{\partial V} \right)_T = \frac{1}{3} \alpha T^4 = \frac{1}{3} \frac{U}{V}$$

(Notice that for ideal gas of particles with mass, we found $p = \frac{2}{3} \frac{U}{V}$ PS #3)

(2)

So far these are properties of photon gas inside cavity
 we want to know properties of photons emitted by a hot object

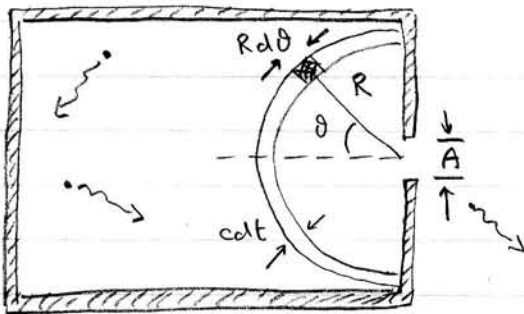
Imagine a cavity with a small hole of size A
 photons emitted have same spectrum as photons inside

Define the "energy flux density" J_u as rate energy is emitted
 per unit area (units $\frac{W}{m^2}$). Also called intensity

By dimensional analysis, expect

$$J_u = c u_{\text{hot}} \times (\text{geometrical factor})$$

Now derive this exactly:

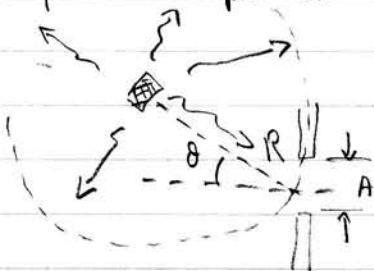


Photons that escape cavity now
 in time interval dt came
 from hemispherical shell some
 time before

Energy in a chunk of shell is:

$$\frac{U}{V} R^2 \sin\theta d\theta d\phi c dt$$

Not all energy from chunk of volume escaped, because not
 all photons point the right way



Photons emitted in every direction
 Equal chance to land anywhere on
 sphere of area $4\pi R^2$
 Area subtended by hole is $A \cos\theta$
 so prob. escape = $\frac{A \cos\theta}{4\pi R^2}$

(3)

So energy that escaped is:

$$\begin{aligned}
 dE_{\text{escape}} &= \int \frac{A \cos \theta}{4\pi R^2} \frac{U}{V} R^2 \sin \theta d\theta d\phi c dt \\
 &\quad \text{hemisphere} \\
 &= \frac{A c dt U}{4\pi V} \underbrace{\int_0^{2\pi} d\phi}_{2\pi} \underbrace{\int_0^{\pi/2} \sin \theta \cos \theta d\theta}_{\frac{1}{2}} \\
 &= \frac{1}{4} A c dt \frac{U}{V}
 \end{aligned}$$

$$\Rightarrow \text{Power} = \frac{dE_{\text{escape}}}{dt} = \frac{1}{4} A c \frac{U}{V}$$

$$J_u = \frac{\text{Power}}{A} = \frac{1}{4} c \frac{U}{V} = \frac{1}{4} c u_{\text{tot}}$$

$$J_u = \frac{\pi^2 (k_B T)^4}{60 \hbar^3 c^2} = \sigma_B T^4$$

$$\sigma_B = 5.67 \times 10^{-8} \frac{\text{W}}{\text{m}^2 \text{K}^4}$$

↑
Stefan-Boltzmann constant

This described flow of energy out of cavity, but what about an object in space radiating? (The sun is not a cavity with a hole in it!)

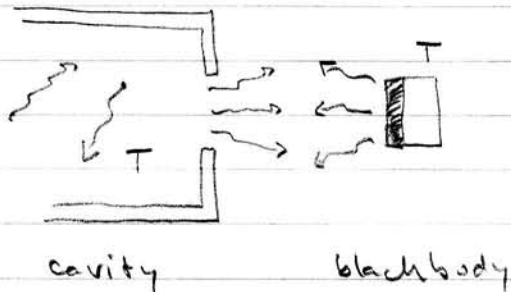
Equations turn out to be exactly the same.

Consider a black object - absorbs all EM radiation
- also emits EM radiation

object is a "blackbody"

Proof that blackbody emits radiation the same as a cavity with a hole

Thought experiment - hole in a cavity, blackbody face each other. Both at same T same size A



- Each emits photons and absorbs each other's photons
- In fact, each absorbs same fraction of other's radiation

What power is emitted by the blackbody?

If less than cavity, then energy flows to blackbody

If more " " then " " out of blackbody

Neither can happen because T is the same. (At thermal equilibrium, not net energy flow)

$$\therefore J_{\text{cavity}} = J_{\text{blackbody}} = \frac{\text{Power}}{\text{Area}}$$

Not only that, but spectra must be the same.

Use same argument, but now place a filter between cavity hole and blackbody that lets through frequencies $\omega, \omega + d\omega$

Blackbody radiation

- absorbs all light (no reflections hence "black")
- emits radiation according to its temperature spectrum given by Planck law
- total energy flux by Stefan Boltzmann law
- $P_{\text{emit}} = P_{\text{absorbed}}$ black body radiates as much energy as it absorbs

What about bodies that reflect light (i.e. objects that are not black)?

Go back to thought experiment

Let's say that for every 3 photons emitted by cavity that hit object, 1 is reflected 2 absorbed (at one λ)



To remain at equilibrium with cavity at T , the object must emit 2 photons — no net energy flow

$$P_{\text{abs}}(\omega) = a(\omega) P_{\text{inc}}(\omega)$$

↑
absorption coefficient at ω (here $\frac{2}{3}$)

$$P_{\text{emit}}(\omega) = P_{\text{abs}}(\omega)$$

$$e(\omega) P_{\text{inc}}(\omega) = a(\omega) P_{\text{inc}}(\omega)$$

↑
emissivity at ω (here $\frac{2}{3}$) $\therefore a(\omega) = e(\omega)$

Also $P_{\text{ref}}(\omega) = r(\omega) P_{\text{inc}}$

↑
reflectance at ω

By energy balance $P_{\text{emit}} + P_{\text{ref}} = P_{\text{inc}}$ so

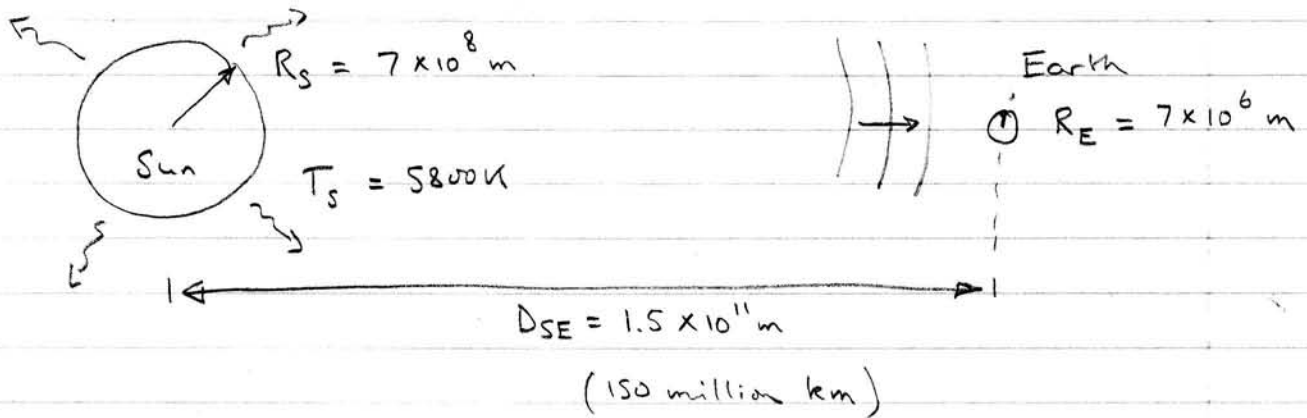
$$a(\omega) = e(\omega) = 1 - r(\omega)$$

- A good absorber is a good emitter (Ex: black body $a=e=1$)
- A good reflector is a poor absorber and emitter

(6)

Assuming e is constant with ω : $J_u = \sigma e T^4$
 For many things $e = e(\omega)$

Ex: The Sun and the Earth



$$\begin{aligned}
 P_s &= \sigma A_s T_s^4 = \sigma 4\pi R_s^2 T_s^4 \\
 &= (5.67 \times 10^{-8}) 4\pi (7 \times 10^8)^2 (5800)^4 \\
 &= 3.9 \times 10^{26} \text{ W} \quad \text{sun's luminosity}
 \end{aligned}$$

What is the energy flux at the earth?

$$J_s = \frac{P_s}{4\pi D_{SE}^2} = \frac{3.9 \times 10^{26} \text{ W}}{4\pi (1.5 \times 10^{11} \text{ m})^2} = 1370 \frac{\text{W}}{\text{m}^2} \quad \text{solar constant}$$

So the incident power at the earth is

$$J_s \times \text{Cross section of Earth} = J_s (\pi R_E^2) = P_{\text{inc}}$$

Assuming the earth is an ideal blackbody ($a = e = 1$)

$$P_{\text{abs}} = J_s \pi R_E^2 = P_{\text{emit}} = \sigma 4\pi R_E^2 T_E^4$$

so...

(7)

$$T_E = \left(\frac{J_S}{4\sigma} \right)^{1/4} = \left(\frac{1370}{4 \times 5.67 \times 10^{-8}} \right)^{1/4} = 279 \text{ K} \quad !! \quad (\text{pretty close to } 15^\circ\text{C})$$

$$T_E = \left(\frac{\cancel{\sigma} 4\pi R_S^2 T_S^4}{4\sigma 4\pi D_{SE}^2} \right)^{1/4} = \left(\frac{R_S}{2D_{SE}} \right)^{1/2} T_S$$

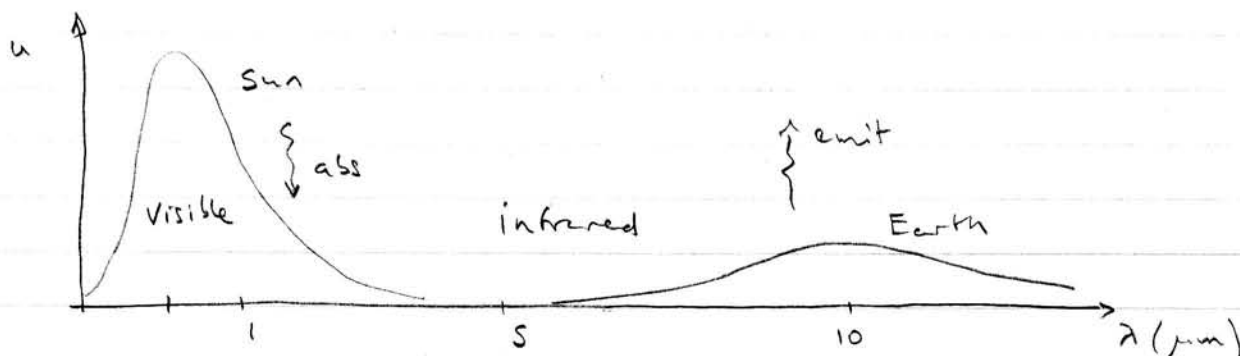
Unfortunately, earth is not an ideal black body, it reflects $\sim 30\%$ of sunlight, so we should replace 1370 with $1370 \times (1 - 0.3)$

$$\Rightarrow T_E = \left(\frac{1370 \times 0.7}{4 \times 5.67 \times 10^{-8}} \right)^{1/4} = 255 \text{ K} \quad !! \quad (\text{frigid!})$$

This is also incorrect: solution is that $e(\omega)$ is a strong function of ω

Solar radiation peaked in near infrared — most light in visible
($\hbar\omega_{\text{max}} = 2.82 k_B T_S = 1.41 \text{ eV} \rightarrow 880 \text{ nm}$)

Earth radiation peaked in far infrared $\rightarrow \sim 10 \mu\text{m}$
because $T_E \sim 300 \text{ K}$



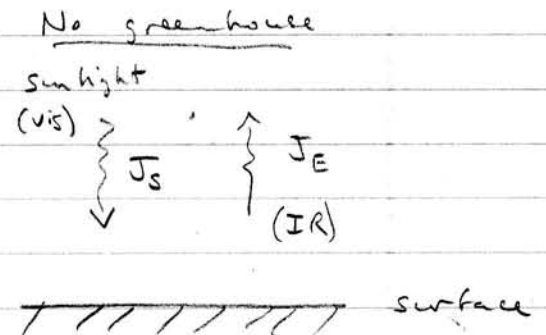
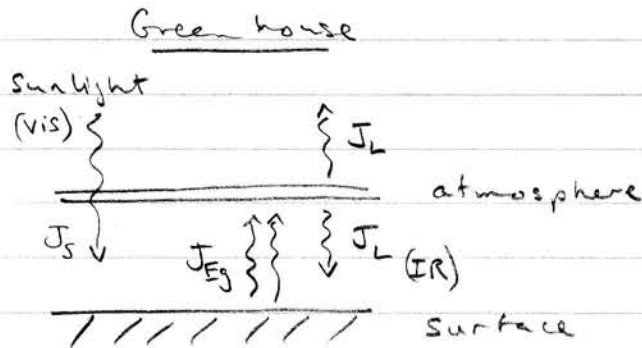
Earth absorbs visible light from sun (some reflected by clouds)
but has to re-emit it as IR light (because $T_E \ll T_S$)

Problem is that atmosphere absorbs IR light

$\Rightarrow$ this leads to greenhouse effect

Simple model for Greenhouse effect (see U+K)

Treat atmosphere as single layer above surface.
Transparent to visible, absorbs IR light



with greenhouse effect incident energy flux is J_s
emitted flux from Earth is J_{Eg} , from atmosphere layer
 J_L . Atmosphere emits to space and back to earth

By energy balance $J_s = J_L$ and $J_s + J_L = J_{Eg}$
 $\therefore J_{Eg} = 2J_s$

In contrast with no greenhouse $J_s = J_E$

$$\therefore \frac{J_{Eg}}{J_E} = 2 = \left(\frac{T_{Eg}}{T_E} \right)^4$$

$$T_{Eg} = 2^{1/4} T_E = 2^{1/4} \cdot 255 \text{ K} = 303 \text{ K} \quad (\text{a bit better})$$