

Announcements

Office hours:

- Xueda Wen **Wednesday** 10-11 AM 4111 ESB
- Peter Sahanggamu **Wednesday** 3-4 PM 390C L-S Interpass
- Jon Thaler **Thursday** 5-7 PM 158 Loomis
- Chun Kit Chan **Friday** 2-3 PM 279 Loomis

Lectures:

Starting with Friday's lecture, I will post the lecture slides on the web page the night before (early, I hope).

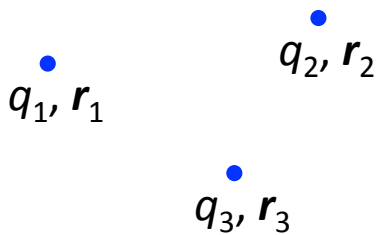
Caveat: I don't guarantee that the slides will be exactly what I show.
(I hope they'll be close.)

Work and Energy (2.4)

We have been talking about the electric potential, and we'll spend a lot of time calculating it.

However, there is another aspect of the problem that we also need to deal with, the amount of energy (Joules, not Volts) needed to assemble a collection of charges.

Consider a collection of point charges:



Note:

It is tacitly assumed that the charges are brought in from infinity, where there is no initial energy.

How much work does it take to assemble this configuration?

A tempting, but **incorrect**, answer: (inspired by $U = qV$)

We can easily calculate the potential at every charge (due to the others):

$$V(\vec{r}_i) = \frac{1}{4\pi\epsilon_0} \sum_{j \neq i} \frac{q_j}{r_j} \quad \text{where } \vec{r}_j = \vec{r}_i - \vec{r}_j.$$

So, is $W = \sum_i q_i V(\vec{r}_i)$??? No !!!

Note:

$j \neq i$, because a charge does not exert a force on itself.

The problem with this:

We are wrong by a factor of two. To see why, consider two point charges:

$$\begin{array}{cc} \bullet & \bullet \\ q_1, \mathbf{r}_1 & q_2, \mathbf{r}_2 \end{array}$$

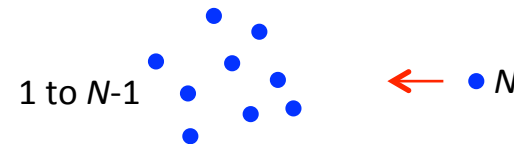
Put it together one charge at a time. The first one takes no energy. The second one takes this much work:

$$W = \left(\frac{1}{4\pi\epsilon_0} \frac{q_1}{r_{12}} \right) q_2 = q_1 V_1(\vec{r}_2) = q_2 V_2(\vec{r}_1)$$

So, if we sum over both charges, we'll be double counting.

By superposition, this argument can be extended to any collection of charges.

Another way to see this. Bring in the N^{th} charge:



The additional energy (work) is only due to the force that 1 to $N-1$ exert on N , not the force that N exerts on 1 to $N-1$, because they are not moving.

So, the correct answer is:

$$W = \frac{1}{2} \frac{1}{4\pi\epsilon_0} \sum_{i=1}^N \sum_{j \neq i}^N \frac{q_i q_j}{r_{ij}}$$

Continuous charge distributions:

For a continuous charge distribution, the sums become a single volume integral:

$$W = \frac{1}{2} \int \rho(\vec{r}) \underbrace{V(\vec{r})}_{\text{Was: } \sum_j \frac{q_j}{r_{ij}}} d\text{Vol}$$

Was: q_i

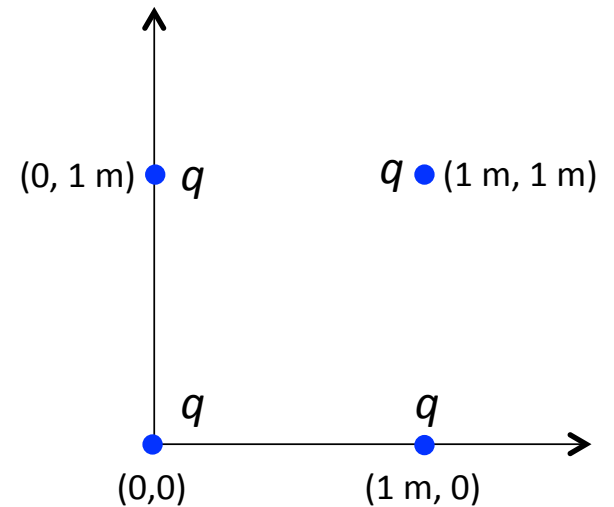
There is only one integral, because the second sum is implicit in $V(r)$.
Also, we no longer need to worry about $i \neq j$.

Example:

Four equal charges at the corners of a square:
 $q = 10^{-6} \text{ C}$.

The potential at each corner is that due to the other three charges:

$$V = \frac{10^{-6} \text{ C}}{4\pi\epsilon_0} \left(\frac{2}{1 \text{ m}} + \frac{1}{\sqrt{2} \text{ m}} \right) = 24.3 \text{ kV}$$



So, the energy required to assemble this configuration (what you could get back by taking it apart) is:

$$W = \frac{1}{2} (4 \times 24.3 \text{ kV} \times 10^{-6} \text{ C}) = 0.486 \text{ J}$$

Example:

Calculate the energy stored in a spherical shell of radius $R = 1$ m, that has a uniform surface charge density, $\sigma = 10^{-6}$ C/m².

This problem has spherical symmetry. At the surface: $V = \frac{Q_{\text{tot}}}{4\pi\epsilon_0 R} = \frac{\sigma R}{\epsilon_0} = 113$ kV

$$W = \frac{1}{2} \int \sigma V dA = \frac{1}{2} \left(\frac{\sigma^2 R}{\epsilon_0} \right) (4\pi R^2) = 0.71 \text{ J}$$

Comments:

- A microCoulomb is a lot of charge. It gives significant energies.
- Because this is surface charge, $\int \rho dV \rightarrow \int \sigma dA$

Configuration Energy in Terms of E

This involves some math, but it's instructive.

$$W = \frac{1}{2} \int \rho V d\text{Vol}$$

$$\text{Use: } \rho = \epsilon_0 \vec{\nabla} \cdot \vec{E}$$

$$= \frac{\epsilon_0}{2} \int (\vec{\nabla} \cdot \vec{E}) V d\text{Vol}$$

Integrate by parts: $u dv = d(uv) - v du$. Here, u is V and v is \vec{E} .

$$= \frac{\epsilon_0}{2} \left[\oint_{\text{surface}} (\vec{E} V) \cdot d\vec{a} - \int_{\text{Vol}} \vec{E} \cdot \vec{\nabla} V d\text{Vol} \right]$$

Now, note that: $\vec{E} = -\vec{\nabla} V$.

$$W = \frac{\epsilon_0}{2} \int_{\text{Vol}} E^2 d\text{Vol} \text{ plus the surface integral.}$$

To see how this works with vectors, write the x , y , and z components explicitly.

Commentary:

Integration by parts is a very common manipulation, because the surface term can often be ignored, for example by letting the volume of integration become very large (or infinite).

This works if the u and v quantities (here, E and V) go to zero sufficiently rapidly.

Example: A localized charge distribution

For large r , the surface area increases as r^2 , E falls as $1/r^2$, and V falls as $1/r$. EV falls as $1/r^3$, so it works.

BEWARE of situations where this is not the case, such as an infinite line of charge.

BEWARE also of using E to calculate the energy of point charges.

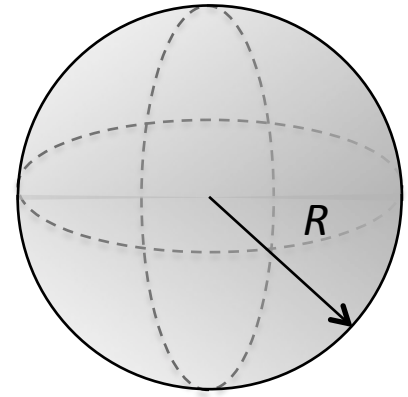
You'll get infinity (next slide).

Calculating the energy by integrating E^2 looks like a bookkeeping trick. However, it is possible to have E in the absence of any charges (EM radiation). In this situation, the E and B fields themselves carry energy – it's not merely a trick.

A puzzle:

In next week's homework, you'll calculate the energy of a uniform sphere of charge.

$$\text{Answer: } \text{Energy} = \frac{3}{5} \frac{1}{4\pi\epsilon_0} \frac{Q^2}{R}$$



An interesting question:

Suppose the electron is a small sphere of charge.

For what radius, R , does the electrostatic energy equal the rest energy, $m_e c^2$?

$$m_e c^2 = 511 \text{ eV} = 8.2 \times 10^{-14} \text{ J, so } R = 1.7 \times 10^{-15} \text{ m.}$$

However, the electron is known to be smaller than 10^{-22} m !!

TL;DR: Classical E&M fails at these distance scales.

This question is addressed in quantum field theory (Physics 582-3).

Note: The 3/5 factor depends on the charge distribution.

What distribution do you think gives the smallest factor
(minimizing R for a given E)?