

## Notes:

- The torque makes  $\mathbf{m}$  line up with  $\mathbf{B}$  (paramagnetism).
- Just like in the electrical case:  $U = -\vec{m} \cdot \vec{B}$
- $\mathbf{F} = 0$  in a uniform field.
- If the field is not uniform:  $\vec{F} = -\vec{\nabla} U = \nabla(\vec{m} \cdot \vec{B})$
- One can minimize  $U$  by: Rotating  $\mathbf{m}$ , or by moving to larger  $\mathbf{B}$ .

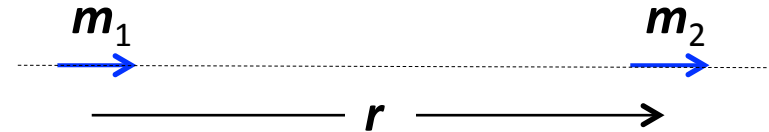
## Microscopic picture of magnetic materials:

I only mention this because Griffiths (p. 258) is incorrect.  
There are **two origins** of magnetic moment.

- Current loops (*e.g.*, electrons orbiting in atoms).
- Intrinsic magnetic moments (due to spin).  
The electron's spin and magnetic moment are not due to a spatial motion of charge. In fact (as far as we can tell), the electron is a point particle. The spin of elementary particles is a purely quantum mechanical effect and can't be understood classically.

## Interacting magnetic dipoles: (G, problem 6.3)

Calculate the force between two dipoles.



The field due to a dipole is:

$$\vec{B} = \frac{\mu_0 m}{4\pi r^3} (2\cos\theta \hat{r} + \sin\theta \hat{\theta}) \quad \text{just like an electric dipole.}$$

In this case,  $\theta = 0$ , so  $\mathbf{B}$  points along  $\hat{r}$  as does  $m_2$ .

$$\text{So, } \vec{F} = \vec{\nabla}(\vec{m} \cdot \vec{B}) = -\frac{3\mu_0 m_1 m_2}{\pi r^4} \hat{r}$$

← You can do this in spherical or Cartesian coordinates

# Diamagnetism (6.1.3)

## and the adiabatic approximation

Griffiths shows the microscopic origin of diamagnetism (induced  $\mathbf{m}$  opposes  $\mathbf{B}$ ). However, [his derivation has a hidden assumption](#), which I'll point out.

Consider an electron in a circular orbit about the nucleus.

It behaves like a current loop:

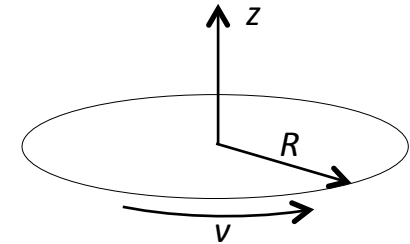
$$I = -e \frac{v}{2\pi R}$$

So, there is a dipole moment along  $z$ :

$$m_z = -\frac{1}{2}evR$$

The circular motion is due to the electric force:

$$\frac{1}{4\pi\epsilon_0} \frac{e^2}{R^2} = m_e \frac{v^2}{R}$$



[Now, turn on  \$B\_z\$ .](#)

This is where the approximation comes in. We assume that the change of the magnetic force per period is small:  $\frac{d(evB)}{dt} \left( \frac{2\pi R}{v} \right) \ll m_e \frac{v^2}{R}$ . In this case the orbit remains circular (not elliptical). [This is called the adiabatic approximation](#) (slow changes).

End 10/23/13