

Maxwell's Equations, Completed (7.3)

The equations we have so far are inconsistent.

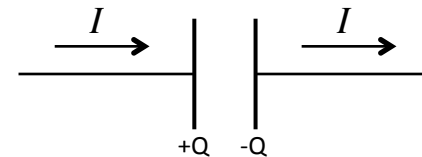
In particular, there is a problem with Ampere's Law (see Oct. 18 lecture):

Apply this vector identity: $\vec{\nabla} \cdot (\vec{\nabla} \times \vec{v}) = 0$ to Ampere's law: $\vec{\nabla} \times \vec{B} = \mu_0 \vec{J}$

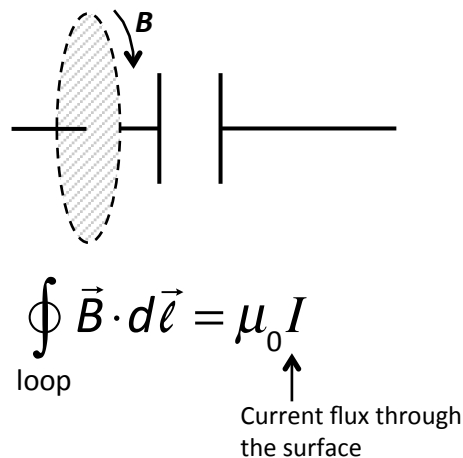
$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) = \mu_0 \vec{\nabla} \cdot \vec{J} = -\mu_0 \frac{\partial \rho}{\partial t}$$

Ampere's law must be modified when we get away from the static situation.

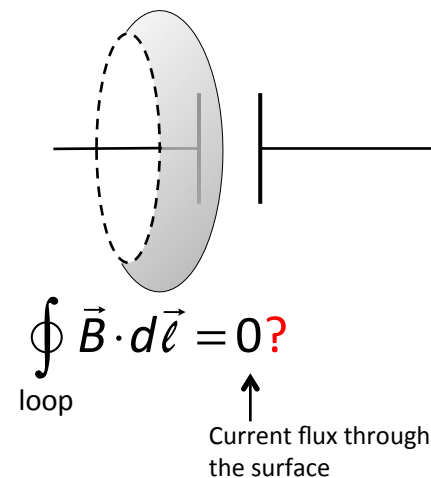
Recall from P212 a capacitor that is being charged up:



Draw an Amperean loop around the wire:



Same loop,
different surface \rightarrow



Can we fix the problem? Or, does Ampere's Law fail?

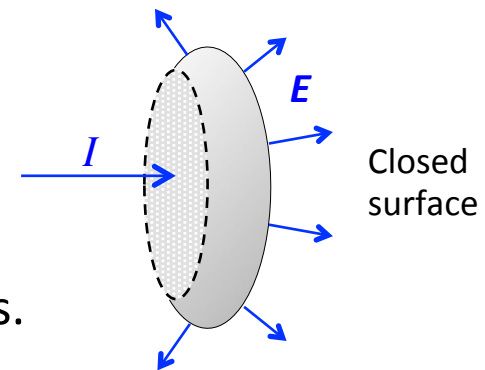
The continuity equation and Gauss's law save us:

$$\text{Continuity: } \vec{\nabla} \cdot \vec{J} = -\frac{\partial \rho}{\partial t} \Rightarrow \oint_{\text{surface}} \vec{J} \cdot d\vec{a} = - \int_{\text{volume}} \frac{\partial \rho}{\partial t} = -\frac{dQ_{\text{in}}}{dt}$$

$$\text{Gauss: } Q_{\text{in}} = \epsilon_0 \int_{\text{surface}} \vec{E} \cdot d\vec{a} \Rightarrow \frac{dQ_{\text{in}}}{dt} = \epsilon_0 \int_{\text{surface}} \frac{\partial \vec{E}}{\partial t} \cdot d\vec{a}$$

$$\text{Therefore: } \oint_{\text{surface}} \vec{J} \cdot d\vec{a} = -\epsilon_0 \oint_{\text{surface}} \frac{\partial \vec{E}}{\partial t} \cdot d\vec{a}$$

In our capacitor example: Current flows in, and outward electric flux increases.



There is no reason the current flow and the electric flux need to be spatially separated. In general: $\oint_{\text{surface}} (\vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t}) \cdot d\vec{a} = 0$

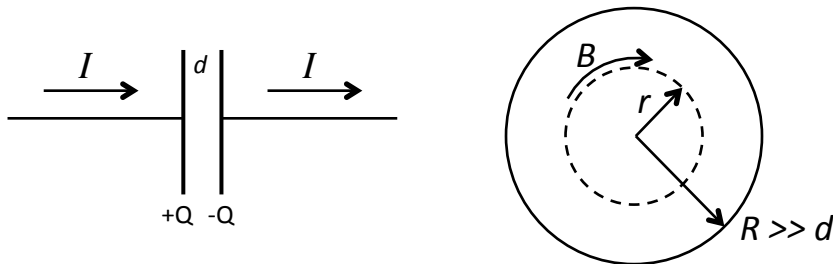
Therefore, this modification of Ampere's Law works:

$$\oint_{\text{loop}} \vec{B} \cdot d\vec{\ell} = \mu_0 \int_{\text{surface}} (\vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t}) \cdot d\vec{a} \Rightarrow \vec{\nabla} \times \vec{B} = \mu_0 \left(\vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right)$$

A changing electric field produces a magnetic field.

Displacement current

Example: Magnetic field in a capacitor.



\mathbf{E} is uniform in the gap: $E = \frac{\sigma}{\epsilon_0} = \frac{Q}{\pi R^2 \epsilon_0}$. $\frac{dE}{dt} = \frac{I}{\pi R^2 \epsilon_0}$

There is also a magnetic field in the ϕ direction:

$$2\pi r B = \epsilon_0 \mu_0 (\pi r^2) \frac{dE}{dt} \Rightarrow B = \frac{\mu_0 I}{2\pi R^2} r$$

The displacement current fills the gap, not just along a line. So, to enclose the entire displacement current, we need $r = R$.

At $r = R$: $B = \frac{\mu_0 I}{2\pi R}$. This is the same as around a wire, as expected.

Maxwell's equations:

$$\begin{aligned}\vec{\nabla} \cdot \vec{E} &= \frac{\rho}{\epsilon_0} & \vec{\nabla} \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t} \\ \vec{\nabla} \cdot \vec{B} &= 0 & \vec{\nabla} \times \vec{B} &= +\mu_0 \left(\vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right)\end{aligned}$$

Beware the \pm signs.

- A changing electric field produces a magnetic field
- $\vec{\nabla} \times \vec{E}$ has no “current” term, and $\vec{\nabla} \cdot \vec{B}$ has no “charge” term, because there are no magnetic monopoles.

End 11/22/13