

## Meissner effect:

Consider this situation:

In the superconductor:

$$\vec{\nabla} \times \vec{J} = -\frac{q^2 n}{m} \vec{\nabla} \times \vec{A} = -\frac{q^2 n}{m} \vec{B}$$

$$\vec{\nabla} \times (\vec{\nabla} \times \vec{B}) = \vec{\nabla} (\underbrace{\nabla \cdot \vec{B}}_{=0}) - \nabla^2 \vec{B}$$

$$= \mu_0 (\vec{\nabla} \times \vec{J})$$

$$= -\frac{\mu_0 q^2 n}{m} \vec{B}$$

So:  $\nabla^2 \vec{B} = +k \vec{B}$

Specifically:  $\nabla^2 B_y = +k B_y$

$$B_y = B_0 e^{\pm \sqrt{k} x}$$

$B_{\parallel}$  is continuous

London's condition

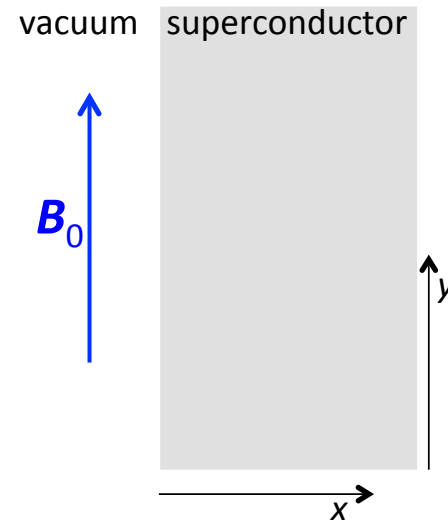
Math identity

Maxwell

London's condition

$$k = \frac{\mu_0 q^2 n}{m}$$

$\frac{1}{\sqrt{k}}$  is called the penetration depth,  $\lambda_L$ .



$\lambda_L$  is typically a few tens of nanometers.

## Magnetic levitation above a superconductor:

$\mathbf{B} = 0$  inside the superconductor.

Boundary condition:  $B_z = 0$  at  $z = 0$ .

This problem can be solved using images.

Remember the dipole field (Sep. 25 lecture).

$$\vec{B}_{\text{dipole}} = \frac{\mu_0 m}{4\pi r^3} (2\cos\theta \hat{r} + \sin\theta \hat{\theta})$$

An image dipole,  $-\mathbf{m}$ , at  $-z$  will cancel  $B_z$  at  $z = 0$ .

This works for any orientation of  $\mathbf{m}$ , not only the one I drew.

The force on  $\mathbf{m}$  is the force due to the image.

$$\vec{F} = \vec{\nabla}(\vec{m} \cdot \vec{B})$$

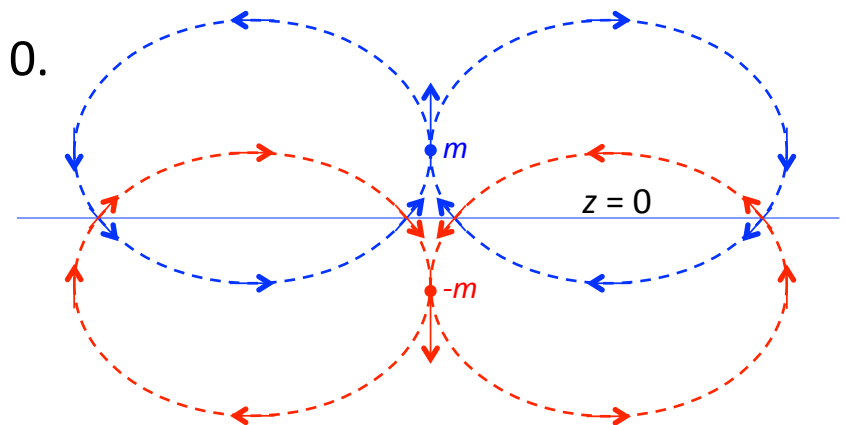
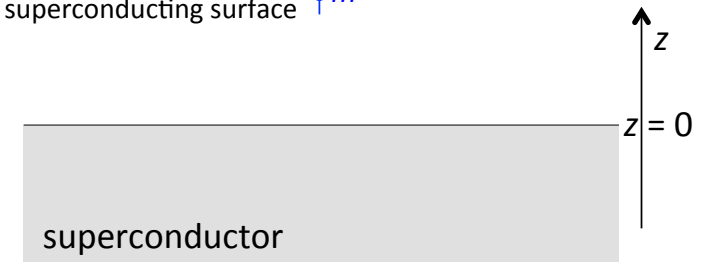
where  $\mathbf{B}$  is the field at  $\mathbf{m}$  due to the image.

In our situation:  $r = 2z$ , and  $\theta = \pi$  ( $180^\circ$ ), because the image points down:

$$B_z(\text{due to image}) = -\frac{\mu_0 m}{2\pi r^3} \Rightarrow F_z = \frac{d}{dr} \left( -\frac{\mu_0 m^2}{2\pi r^3} \right) = +\frac{3\mu_0 m^2}{2\pi r^4}$$

The magnet is levitated.

a small magnet above a  
superconducting surface  $\uparrow m$



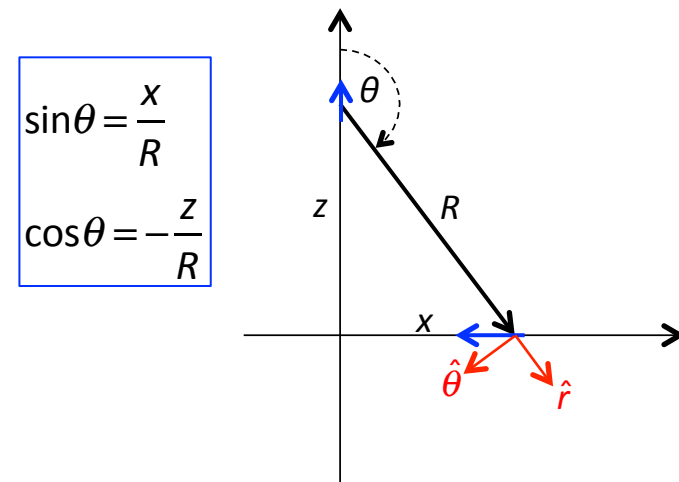
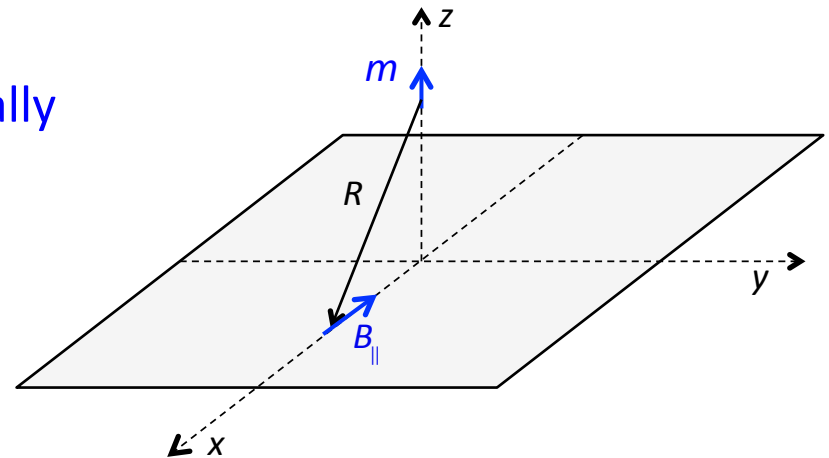
The equilibrium height is given by  $\frac{3\mu_0 m^2}{2\pi(2z)^4} = Mg$ .

As with image charge, the **image dipole is actually surface current**. We can calculate this current using the tangential component of  $\mathbf{B}$  and Ampere's law. There is cylindrical symmetry, so  $B_{\parallel}$  must be radial (*i.e.*, along  $x$  in the figure).

$$\begin{aligned} B_x &= 2(B_r \sin\theta + B_\theta \cos\theta) \\ &= \frac{\mu_0 m}{2\pi R^3} (2\cos\theta \sin\theta + \sin\theta \cos\theta) \\ &= -\frac{3\mu_0 mzx}{2\pi R^5} \hat{x} \end{aligned}$$

The surface current flows in circles (replace  $x$  with polar coordinate,  $r$ ):

$$\vec{K} = \frac{1}{\mu_0} \hat{z} \times \vec{B} = \frac{-3mz}{2\pi R^5} r \hat{\phi}$$



End 12/2/13