

Lecture 7 - Elasticity

Elasticity Stress and Strain in Crystals Kittel – Ch 3

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Elastic Behavior is the fundamental distinction between solids and liquids

- **Similarity:** both are “condensed matter”
- **A solid or liquid in equilibrium has a definite density** (mass per unit volume measured at a given temperature)
- **The energy increases if the density (volume) is changed from the equilibrium value - e.g. by applying pressure**

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Elastic Behavior is the fundamental distinction between solids and liquids

- **Difference:**
- **A solid maintains its shape**
 - The energy increases if the shape is changed – “shear”
- **A liquid has no preferred shape**
 - It has no resistance to forces that do not change the volume

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Strain and Stress

Strain is a change of relative positions of the parts of the material

Stress is a force /area applied to the material to cause the strain

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Pressure and Bulk Modulus

- **Consider first changes in the volume – applies to liquids and any crystal**
- **General approach:**
 $E(V)$ where V is volume

Can use either $E_{\text{crystal}}(V_{\text{crystal}})$ or $E_{\text{cell}}(V_{\text{cell}})$
since $E_{\text{crystal}} = N E_{\text{cell}}$ and $V_{\text{crystal}} = N V_{\text{cell}}$

- **Pressure = $P = -dE/dV$ (units of Force/Area)**
- **Bulk modulus $B = -V dP/dV = V d^2E/dV^2$ (same units as pressure)**
- **Compressibility $K = 1/B$**

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Total Energy of Crystal

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Elasticity

- Up to now in the course we considered only perfect crystals with no external forces
- **Elasticity describes:**
 - Change in the volume and shape of the crystal when external stresses (force / area) are applied
 - Sound waves
- Some aspects of the elastic properties are determined by the **symmetry** of the crystal
- Quantitative values are determined by **strength and type of binding** of the crystal?

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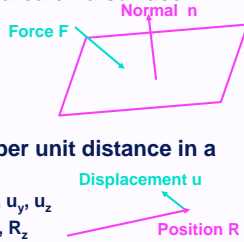
Elastic Equations

- The elastic equations describe the relation of stress and strain
- **Linear relations for small stress/strain**
Stress = (elastic constants) x Strain
- Large elastic constants \Rightarrow the material is stiff - a given strain requires a large applied stress
- We will give the general relations - **but we will consider only cubic crystals**
 - The same relations apply for isotropic materials like a glass
 - More discussion of general case in Kittel

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Elastic relations in general crystals

- **Strain and stress are tensors**
- Stress e_{ij} is force per unit area on a surface
 - Force is a vector F_x, F_y, F_z
 - A surface is defined by the normal vector n_x, n_y, n_z
 - $3 \times 3 = 9$ quantities
- Strain σ_{ij} is displacement per unit distance in a particular direction
 - Displacement is a vector u_x, u_y, u_z
 - A position is a vector R_x, R_y, R_z
 - $3 \times 3 = 9$ quantities



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Elastic Properties of Crystals

- **Definition of strain**
Six independent variables:
 $e_1 \equiv e_{xx}, e_2 \equiv e_{yy}, e_3 \equiv e_{zz},$
 $e_4 \equiv e_{yz}, e_5 \equiv e_{xz}, e_6 \equiv e_{xy}$
- Using the relation $e_{xy} = e_{yx}$ etc.
- **Stress**
 $\sigma_1 \equiv \sigma_{xx} = X_x, \sigma_2 \equiv Y_y, \sigma_3 \equiv Z_z,$
 $\sigma_4 \equiv Y_z, \sigma_5 \equiv X_z, \sigma_6 \equiv X_y$
- Here X_y denotes force in x direction applied to surface normal to y.
 $\sigma_{xy} = \sigma_{yx}$ etc.
- **Linear relation of stress and strain**
Elastic Constants C_{ij}
 $\sigma_i = \sum_j C_{ij} e_j, (i, j = 1, 6)$
 (Also compliances $S_{ij} = (C^{-1})_{ij}$)

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Strain energy

- **For linear elastic behavior, the energy is quadratic in the strain (or stress)**
Like Hooke's law for a spring
- Therefore, the energy is given by:
 $E = (1/2) \sum_i e_i \sigma_i = (1/2) \sum_{ij} e_i C_{ij} e_j, (i, j = 1, 6)$
- Valid for all crystals
- Note 21 independent values in general (since $C_{ij} = C_{ji}$)

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Symmetry Requirements Cubic Crystals

- **Simplification in cubic crystals due to symmetry** since x, y, and z are equivalent in cubic crystals
- For cubic crystals all the possible linear elastic information is in 3 quantities:
 $C_{11} = C_{11} = C_{22} = C_{33}$
 $C_{12} = C_{13} = C_{23}$
 $C_{44} = C_{55} = C_{66}$
- Note that by symmetry $C_{14} = 0$, etc
- Why is this true for cubic crystals?

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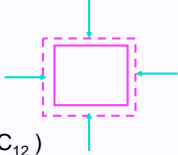
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Elasticity in Cubic Crystals

- Elastic Constants** C_{ij} are completely specified by 3 values C_{11} , C_{12} , C_{44}

$$\sigma_1 = C_{11} e_1 + C_{12} (e_2 + e_3), \text{ etc.}$$

$$\sigma_4 = C_{44} e_4, \text{ etc.}$$
- Pure change in volume – compress equally in x, y, z
- For pure dilation $\delta = \Delta V / V$
 $e_1 = e_2 = e_3 = \delta / 3$
- Define $\Delta E / V = 1/2 B \delta^2$
- Bulk modulus** $B = (1/3) (C_{11} + 2 C_{12})$



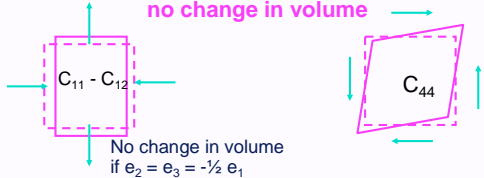
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Elasticity in Cubic Crystals

- Elastic Constants** C_{ij} are completely specified by 3 values C_{11} , C_{12} , C_{44}

$$\sigma_1 = C_{11} e_1 + C_{12} (e_2 + e_3), \text{ etc.}$$

$$\sigma_4 = C_{44} e_4, \text{ etc.}$$
- Two types of shear – no change in volume

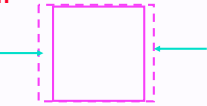


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Elasticity in Cubic Crystals

- Pure uniaxial stress and strain**
- $\sigma_1 = C_{11} e_1$ with $e_2 = e_3 = 0$
- $\Delta E = (1/2) C_{11} (\delta x/x)^2$
- Occurs for waves where there is no motion in the y or z directions

Also for a crystal under $\sigma_1 \equiv X_x$ stress if there are also stresses $\sigma_2 \equiv Y_y$, $\sigma_3 \equiv Z_z$ of just the right magnitude so that $e_2 = e_3 = 0$



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Elastic Waves

- The general form of a displacement pattern is $\Delta \mathbf{r}(\mathbf{r}) = u(\mathbf{r}) \mathbf{x} + v(\mathbf{r}) \mathbf{y} + w(\mathbf{r}) \mathbf{z}$
- A traveling wave is described by $\Delta \mathbf{r}(\mathbf{r}, t) = \Delta \mathbf{r} \exp(i\mathbf{k} \cdot \mathbf{r} - i\omega t)$
- For simplicity consider waves along the x direction in a **cubic crystal**

Longitudinal waves (motion in x direction) are given by $u(x) = u \exp(ikx - i\omega t)$

Transverse waves (motion in y direction) are given by $v(x) = v \exp(ikx - i\omega t)$

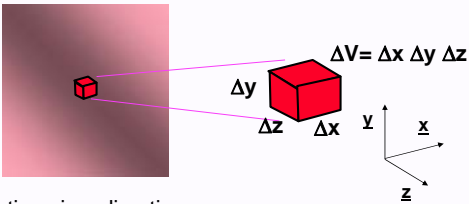
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Waves in Cubic Crystals

- Propagation follows from Newton's Eq. on each volume element
- Longitudinal waves:**
 $\rho \Delta V d^2 u / dt^2 = \Delta x dX_x / dx = \Delta x C_{11} d^2 u / dx^2$
 (note that strain is $e_1 = d u / dx$)
- Since $\Delta V / \Delta x =$ area and ρ area = mass/length = ρ_L , this leads to
 $\rho_L u / dt^2 = C_{11} du / dx$
 or
 $\omega^2 = (C_{11} / \rho_L) k^2$
- Transverse waves** (motion in the y direction) are given by
 $\omega^2 = (C_{44} / \rho_L) k^2$

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Elastic Waves



- Variations in x direction
- Newton's Eq: $ma = F$
- Longitudinal: displacement u along x,
 $\rho \Delta V d^2 u / dt^2 = \Delta x dX_x / dx = \Delta x C_{11} d^2 u / dx^2$
- Transverse: displacement v along y,
 $\rho \Delta V d^2 v / dt^2 = \Delta x dY_x / dx = \Delta x C_{44} d^2 v / dx^2$

Net force in x direction

Net force in y direction

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Sound velocities

- The relations before give (valid for any elastic wave):

$$\omega^2 = (C / \rho_L) k^2 \quad \text{or} \quad \omega = s k$$

- where s = sound velocity
- Different for longitudinal and transverse waves
- Longitudinal sound waves can happen in a liquid, gas, or solid
- Transverse sound waves exist only in solids
- More in next chapter on waves

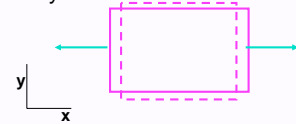
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Young's Modulus & Poisson Ratio

- Consider crystal under tension (or compression) in x direction
- If there are no stresses $\sigma_2 \equiv Y_y$, $\sigma_3 \equiv Z_z$ then the crystal will also strain in the y and z directions

- Poisson ratio defined by $(dy/y) / (dx/x)$

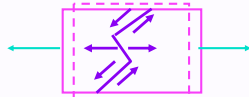
- Young's modulus defined by $Y = \text{tension} / (dx/x)$
- Homework problem to work this out for a cubic crystal



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When does a crystal break?

- Consider crystal under tension (or compression) in x direction
- For large strains, when does it break?
- Crystal planes break apart – or slip relative to one another
- Governed by “dislocations”
- See Kittel – Chapter 20



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Next Time

- Vibrations of atoms in crystals
- Normal modes of harmonic crystal
- Role of Brillouin Zone
- Quantization and Phonons
- Read Kittel Ch 4

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