

Phys 487 Discussion 11 – Fermi’s Golden Rule

Given • $H(t) = H^{(0)} + H'(t)$, • $\{E_n^{(0)}, |n^{(0)}\rangle\}$ = the eigen-* of $H^{(0)}$ • initial state $|\psi(t=0)\rangle = |i^{(0)}\rangle$

then
$$\boxed{|\psi(t)\rangle = \sum_n c_n(t) e^{-i\omega_n t} |n^{(0)}\rangle}$$
 with
$$\boxed{i\hbar \dot{c}_f(t) = \sum_n H'_{fn} e^{i\omega_{fn} t} c_n(t)}$$

• $\omega_{fn} \equiv (E_f^{(0)} - E_n^{(0)}) / \hbar$
 • $H'_{fn} \equiv \langle f^{(0)} | H' | n^{(0)} \rangle$

& to 1st order in $H' \ll H^{(0)}$,
$$\boxed{c_f(t) \approx \delta_{fi} + \frac{1}{i\hbar} \int_0^t H'_{fi}(t') e^{i\omega_{fi} t'} dt'}$$
 \rightarrow
$$\boxed{P_{i \rightarrow f} = |c_f(t)|^2}$$

Problem 1 : Fermi’s Golden Rule for a constant perturbation

In class, we derived Fermi’s Golden Rule for a perturbation that **oscillates sinusoidally with time** (typically in the form of EM radiation) and that is ON for a very long time. As it happens, Fermi’s Golden Rule *also* applies for another extremely common type of perturbation: a potential that is **constant with time**, it merely turns on and off. An example: a simple Stark and/or Zeeman effect experiment when a field $\vec{E}(\vec{r})$ and/or $\vec{B}(\vec{r})$ is turned on at some time $t = 0$ and turned off later. So off we go!

The simplest time-dependent perturbation is a constant potential V that just “turns on” at some time $t = 0$:

$$V(t) = 0 \text{ for } t < 0 \quad \& \quad V(t) = V = \text{constant for } t \geq 0.$$

Important: we are NOT saying that V is constant versus POSITION, only with respect to time. In all of our time-dependent PT work, it is implied that the perturbation labelled “ V ” or “ H ” DOES in general have some \vec{r} - dependence, it will just end up in a transition matrix element $V_{fi} = \langle \psi_f | V | \psi_i \rangle$ that we will have to calculate (as an integral over position) if we want a specific answer for a specific V . If we ever need to specify a potential that is independent of position, we will call it something like “ V_0 ” to denote one single scalar value.

Now suppose we have a system with a solvable unperturbed Hamiltonian H_0 plus the on/off perturbation $V(t)$ given above. What is the transition probability $P_{i \rightarrow f} = |c_f(t)|^2$ to first order?

(a) Derive the following result :
$$P_{i \rightarrow f} = \frac{|V_{fi}|^2}{\hbar^2} \left[\frac{\sin(\omega_{fi} t / 2)}{\omega_{fi} t / 2} \right]^2 t^2 \text{ for } i \neq f.$$

You will need the “half-angle formula” $1 - \cos \theta = 2 \sin^2(\theta / 2)$.

► Is your first thought that the result is a typo? It is always my first thought when seeing that expression for a simple *time-independent* perturbation that just turns ON once and OFF once! “We saw that $\sin^2(\omega t / 2)$ stuff when we worked with sinusoidal perturbations in class, surely it is just a copy/paste error?” Indeed one would think that such a term only appears for sinusoidal perturbations, but no! Start your calculation from time-dependent PT basics (back to the formula sheet!), and observe how that same time-dependent term arises even for our much simpler ON/OFF perturbation. (Actually, look closely: is $P_{i \rightarrow f}$ the same or just similar to the sinusoidal case?)

(b) Prove the following weird but important Dirac delta-function relation :
$$\delta(ax) = \frac{\delta(x)}{a}.$$

► Remember that the defining properties of the Dirac delta are on your 486 formula sheet, consult those to derive/prove the above relation, and the one in the next part.

(c) Prove that the following is a delta function :
$$\lim_{a \rightarrow \infty} \frac{1}{\pi} \frac{\sin^2(ax)}{ax^2} = \delta(x).$$
 (You will need $\int_{-\infty}^{\infty} \frac{\sin^2 x}{x^2} dx = \pi$.)

(d) Combining the above, show that the **transition rate energy is conserved**, i.e. where $(E_f - E_i) \rightarrow 0$.

$$R_{i \rightarrow f} \equiv \frac{P_{i \rightarrow f}}{t} = \frac{2\pi}{\hbar} |V_{fi}|^2 \delta(E_f - E_i)$$

in the limit where

This relation is one form of **Fermi's Golden Rule** for energy-conserving transitions.

(e) Is it reasonable to insist that energy is conserved despite the change in potential energy? Return to expression (a) and consider its dependence on $\omega_{fi} = \hbar(E_f - E_i)$. As you can quickly check with some sort of machine, the function $\sin^2 x / x^2$ is peaked at $x=0$ and has a width of about π . (It reaches $0.4 \approx 1/2$ at $x = \pm\pi/2$). Given this info, what range of ω_{fi} values keeps the transition probability $P_{i \rightarrow f}$ within a factor of about 2 of its maximum value? Your answer will involve time, t . Does the range of probable transition frequencies increase or decrease with t ?

(f) Hopefully what you found was that, as $t \rightarrow \infty$, the width in reasonably-probable transition frequencies goes to zero. This is a very important result! What does t represent, exactly? Once you know, you can say something like this

“In the limit that the perturbation $V(t)$ _____(words)_____, the only transition frequency with any finite probability is $\omega_{fi} = \text{_____}$, which means that energy is _____ in this limit.”

We have thus clarified the conditions under which part (d) is a valid result.

(g) You just found that **energy is conserved more exactly** in the transition from state i to state f **as the time t that the perturbation has been ON increases**. What relation involving a German name is this related to?

(h) The perturbation can never be on *forever*, i.e. we can never reach the limit $t \rightarrow \infty$, so there is always some non-zero range of final-state energies $E_f \approx E_i$ that can be reached from an initial-state energy E_i ... but a transition $E_i \rightarrow E_f$ can only occur if there a state with energy E_f actually exists. It is customary to inject information about the availability of final states into Fermi's Golden Rule using the quantity $\rho(E_f)$ = the density of final states = #states per unit energy. This quantity has units of 1/energy. The energy-conserving delta function $\delta(E_f - E_i)$ in our earlier version of Fermi's Golden Rule *also* has units of 1/energy. To get the most familiar form of F.G.R., we simply replace the one-final-state-only δ -function with the density of states:

$$R_{i \rightarrow f} \equiv \frac{P_{i \rightarrow f}}{t} = \frac{2\pi}{\hbar} |V_{fi}|^2 \rho(E_f) \Big|_{E_f=E_i}$$

Fermi's Golden Rule