

**1) Linear differential operators:**

- a) Let  $w(x) > 0$ . Consider the differential operator  $\hat{L} = id/dx$ . Find the formal adjoint of  $L$  with respect to the inner product  $\langle u|v \rangle_w = \int wu^*v dx$ , and find the corresponding surface term  $Q[u, v]$ .
- b) Now do the same for the operator  $M = d^4/dx^4$ , for the case  $w = 1$ . Find the adjoint boundary conditions defining the domain of  $M^\dagger$  for the case

$$\mathcal{D}(M) = \{y, y^{(4)} \in L^2[0, 1] : y(0) = y'''(0) = y(1) = y'''(1) = 0\}.$$

(Hint: you may find the identity

$$f^{(4)}g - fg^{(4)} = \frac{d}{dx} \{f'''g - f''g' + f'g'' - fg'''\}$$

to be of use.)

**2) Sturm-Liouville forms:** By constructing appropriate weight functions convert the following common operators into Sturm-Liouville form:

- a)  $\hat{L} = (1 - x^2) d^2/dx^2 + [(\mu - \nu) - (\mu + \nu + 2)x] d/dx$ .
- b)  $\hat{L} = (1 - x^2) d^2/dx^2 - 3x d/dx$ .
- c)  $\hat{L} = d^2/dx^2 - 2x(1 - x^2)^{-1} d/dx - m^2(1 - x^2)^{-1}$ .

**3) Discrete approximations and self-adjointness:** Consider the second order inhomogeneous equation  $Lu \equiv u'' = g(x)$  on the interval  $0 \leq x \leq 1$ . Here  $g(x)$  is known and  $u(x)$  is to be found. We wish to solve the problem on a computer, and so set up a discrete approximation to the ODE in the following way:

- replace the continuum of independent variables  $0 \leq x \leq 1$  by the discrete lattice of points  $0 \leq x_n \equiv n/N \leq 1$  Here  $N$  is a positive integer and  $n = 0, 1, 2, \dots, N$ ;
- replace the functions  $u(x)$  and  $g(x)$  by the arrays of real variables  $u_n \equiv u(x_n)$  and  $g_n \equiv g(x_n)$ ;
- approximate the continuum differential operator  $d^2/dx^2$  by the finite difference operator  $\mathcal{D}^2$ , defined by  $\mathcal{D}^2 u_n \equiv (u_{n+1} - 2u_n + u_{n-1})/a^2$  where  $a = N^{-1}$  is the lattice spacing.

Now do the following problems:

- a) Impose continuum Dirichlet boundary conditions  $u(0) = u(1) = 0$ . Decide what these correspond to in the discrete approximation, and write the resulting set of algebraic equations in matrix form. Show that the corresponding matrix is real and symmetric.
- b) Impose the periodic boundary conditions  $u(0) = u(1)$  and  $u'(0) = u'(1)$ , and show that these require us to set  $u_0 \equiv u_N$  and  $u_{N+1} \equiv u_1$ . Again write the system of algebraic equations in matrix form and show that the resulting matrix is real and symmetric.

c) Consider the non-symmetric  $N \times N$  matrix operator

$$D^2 u = \frac{1}{a^2} \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & \dots & 0 \\ 1 & -2 & 1 & 0 & 0 & \dots & 0 \\ 0 & 1 & -2 & 1 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & \dots & 0 & 1 & -2 & 1 & 0 \\ 0 & \dots & 0 & 0 & 1 & -2 & 1 \\ 0 & \dots & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} u_N \\ u_{N-1} \\ u_{N-2} \\ \vdots \\ u_3 \\ u_2 \\ u_1 \end{pmatrix}.$$

- i) What vectors span the null space of  $D^2$ ?
- ii) To what continuum boundary conditions for  $d^2/dx^2$  does this matrix correspond?
- iii) Consider the matrix  $(D^2)^\dagger$ , To what continuum boundary conditions does this matrix correspond? Are they the adjoint boundary conditions for the operator in part ii)?

**4) Factorization:** Schrödinger equations of the form

$$-\frac{d^2\psi}{dx^2} - l(l+1)\text{sech}^2 x \psi = E\psi$$

are known as *Pöschel-Teller equations*. By setting  $u = l \tanh x$  and following the strategy of this problem one may relate solutions for  $l$  to those for  $l - 1$  and so find all bound states and scattering eigenfunctions for any integer  $l$ .

a) Suppose that we know that  $\psi = \exp \left\{ - \int^x u(x') dx' \right\}$  is a solution of

$$L\psi \equiv \left( -\frac{d^2}{dx^2} + W(x) \right) \psi = 0.$$

Show that  $L$  can be written as  $L = M^\dagger M$  where

$$M = \left( \frac{d}{dx} + u(x) \right), \quad M^\dagger = \left( -\frac{d}{dx} + u(x) \right),$$

the adjoint being taken with respect to the product  $\langle u|v \rangle = \int u^* v dx$ .

b) Now assume  $L$  is acting on functions on  $[-\infty, \infty]$  and that we not have to worry about boundary conditions. Show that given an eigenfunction  $\psi_-$  obeying  $M^\dagger M\psi_- = \lambda\psi_-$  we can multiply this equation on the left by  $M$  and so find a eigenfunction  $\psi_+$  with the same eigenvalue for the differential operator

$$L' = MM^\dagger = \left( \frac{d}{dx} + u(x) \right) \left( -\frac{d}{dx} + u(x) \right)$$

and *vice-versa*. Show that this correspondence  $\psi_- \leftrightarrow \psi_+$  will fail if, *and only if*,  $\lambda = 0$ .

- c) Apply the strategy from part b) in the case  $u(x) = \tanh x$  and one of the two differential operators  $M^\dagger M$ ,  $MM^\dagger$  is (up to an additive constant)

$$H = -\frac{d^2}{dx^2} - 2 \operatorname{sech}^2 x.$$

Show that  $H$  has eigenfunctions of the form  $\psi_k = e^{ikx}P(\tanh x)$  and eigenvalue  $E = k^2$  for any  $k$  in the range  $-\infty < k < \infty$ . The function  $P(\tanh x)$  is a polynomial in  $\tanh x$  which you should be able to find explicitly. By thinking about the exceptional case  $\lambda = 0$ , show that  $H$  has an eigenfunction  $\psi_0(x)$ , with eigenvalue  $E = -1$ , that tends rapidly to zero as  $x \rightarrow \pm\infty$ . Observe that there is no corresponding eigenfunction for the other operator of the pair.