

1) Conducting Strip: A thin insulated conducting strip of width $2a$ extends infinitely far in the z direction. It lies in the plane $y = 0$ between $x = -a$ and $x = +a$ and is held at potential V_0 . The rest of the $y = 0$ plane consists of two semi-infinite conducting sheets, one occupying all $x > a$, and one all $x < -a$. Both these sheets are grounded *i.e.* held at potential $V = 0$.

- a) Represent the potential $V(x)$ in the $y = 0$ plane as a Fourier integral

$$V(x) = \int_{-\infty}^{\infty} \frac{dk}{2\pi} A(k) e^{ikx},$$

and explicitly compute $A(k)$.

- b) There is no charge or conductors elsewhere, so the electrostatic potential in the rest of three dimensional space satisfies Laplace's equation with the boundary condition $V \rightarrow 0$ as $|y| \rightarrow \infty$. Write down the Fourier integral giving the function $V(x, y, z)$ at all points in space.
- c) Compute the electric field $\mathbf{E} = -\nabla V$ near $y = 0$ on either side of the conducting planes, and hence find the Fourier integral which gives the charge distribution $\sigma(x)$ on the three conducting sheets.
- d) Perform the integration over k , and hence find $\sigma(x)$ as an explicit (and elementary) function of x .

2) A Recent Qual Problem: A sphere of radius a is made by joining two conducting hemispheres along their equators. The hemispheres are electrically insulated from one another and maintained at two different potentials V_1 and V_2 .

- a) Starting from the general expression

$$V(r, \theta) = \sum_{l=0}^{\infty} \left(a_l r^l + \frac{b_l}{r^{l+1}} \right) P_l(\cos \theta)$$

find an integral expression for the coefficients a_l, b_l that are relevant to the electric field *outside* the sphere. Evaluate the integrals giving b_1, b_2 and b_3 .

- b) Use your results from part a) to compute the electric dipole moment of the sphere as function of the potential difference $V_1 - V_2$.
- c) Now the two hemispheres are electrically connected and the entire surface is at one potential. The sphere is immersed in a uniform electric field \mathbf{E} . What is its dipole moment now?

3) Tides and Gravity : The Earth is not exactly spherical. Two major causes of the deviation from sphericity are the Earth's rotation and the tidal forces it feels from the Sun and the Moon. In this problem we will study the effects of rotation and tides on a self-gravitating sphere of fluid of uniform density ρ_0 .

- a) Consider the equilibrium of a roughly spherical body of fluid rotating homogeneously with angular velocity ω_0 . Show that the effect of rotation can be accounted for by introducing an “effective gravitational potential”

$$\varphi_{\text{eff}} = \varphi_{\text{grav}} + \frac{1}{3}\omega_0^2 R^2 (P_2(\cos \theta) - 1),$$

where R, θ are spherical coordinates defined with their origin in the centre of the body and \hat{z} along the axis of rotation.

- b) A small planet is in a circular orbit about a distant massive star. It rotates about an axis perpendicular to the plane of the orbit so that it always keeps the same face directed towards the star. Show that the tidal forces can be taken into account by introducing an effective external potential

$$\varphi_{\text{tidal}} = -\Omega^2 R^2 P_2(\cos \theta),$$

together with a potential of the same sort as in part a) that accounts for the once-per-orbit rotation. Here Ω is the orbital angular velocity, and R, θ are spherical coordinates defined with their origin at the centre of the planet and \hat{z} pointing at the star.

- c) The external potentials slightly deform the initially spherical planet and the surface is given by

$$R(\theta, \phi) = R_0 + \eta P_2(\cos \theta).$$

Show that, to first order in η , this deformation does not alter the volume of the body. Observe that positive η corresponds to a prolate spheroid and negative η to an oblate one.

- d) The gravitational field of the deformed spheroid can be found by approximating it as an undeformed homogeneous sphere of radius R_0 , together with a thin spherical shell of radius R_0 and surface mass density $\sigma = \rho_0 \eta P_2(\cos \theta)$. Use the general axisymmetric solution

$$\varphi(R, \theta, \phi) = \sum_{l=0}^{\infty} \left(A_l R^l + \frac{B_l}{R^{l+1}} \right) P_l(\cos \theta)$$

of Laplace's equation, together with Poisson's equation

$$\nabla^2 \varphi = 4\pi G \rho(\mathbf{r})$$

for the gravitational potential, to obtain expressions for φ_{shell} in the regions $R > R_0$ and $R \leq R_0$.

e) The surface of the fluid will be an equipotential of the combined potentials of the homogeneous sphere, the thin shell, and the effective external potential of the tidal or centrifugal forces. Use this fact to find η (to lowest order in the angular velocities) for the two cases. Do not include the centrifugal potential from part b) when computing the purely tidal distortion. I only made the planet rotate synchronously in order to keep the the fluid stationary with respect to the tidal bulges.

(See: <http://co-ops.nos.noaa.gov/restles1.html> for a nice explanation of why we ignore the spatial variation of the centrifugal force when computing tidal effects.)

(I get $\eta_{\text{rot}} = -\frac{5}{2} \frac{\omega_0^2 R_0}{4\pi G \rho_0}$, and $\eta_{\text{tide}} = \frac{15}{2} \frac{\Omega^2 R_0}{4\pi G \rho_0}$, but I may have made a mistake!)