

1) Fermat's principle: According to Fermat's principle, the path taken by a ray of light between any two points makes stationary the travel time between those points. A medium is characterized optically by its *refractive index* n , such that the speed of light in the medium is c/n . In this problem we will assume that all light rays lie in the $x - y$ plane.

- a) Use Fermat's principle to show that light propagates along straight lines in homogeneous media (*i.e.*, media in which n is independent of position).
- b) Consider the propagation of light from a flat slab of glass of refractive index n_1 into to another with refractive index n_2 . By examining paths that need not be differentiable at the flat interface, establish Snell's law: $n_1 \sin \theta_1 = n_2 \sin \theta_2$.
- c) A planar light ray propagates in an layered medium with refractive index $n(x)$. Use Fermat's principle to establish a generalized Snell's law in the form $n \sin \psi = \text{constant}$ by finding the equation for stationary paths for

$$F_1 = \int n(x) \sqrt{1 + y'^2} dx.$$

(Here the prime denotes differentiation with respect to x .) Repeat this exercise with $x \leftrightarrow y$ by finding a similar equation for the stationary paths of

$$F_2 = \int n(y) \sqrt{1 + y'^2} dx.$$

By using suitable definitions of the angle of incidence ψ , in each case show that the two formulations of the problem of a layered medium give physically equivalent answers. (In the second formulation you will find it easiest to use the first integral.)

2) Hyperbolic Geometry: This problem involves a version of the Poincaré model for the non-Euclidean geometry of Lobachevski. You may want to review the history of Euclid's fifth postulate on Wikipedia.

- a) Show that the stationary paths for the functional

$$F_3 = \int \frac{1}{y} \sqrt{1 + y'^2} dx,$$

with y restricted to lying in the upper half plane are circles of arbitrary radius **with centers on the x axis**. These paths are the *geodesics*, or minimum length paths, in a space with Riemann metric

$$ds^2 = \frac{1}{y^2} (dx^2 + dy^2), \quad y > 0$$

- b) **OPTIONAL** Show (by giving a compass-and-ruler geometric construction) that if we call these geodesics “lines”, then one and only one line can be drawn though two given points.
- c) **OPTIONAL** Two straight lines are said to be *parallel* if, and only if, they meet at “infinity”, *i.e.* on the x axis. (Verify that the x axis is indeed infinitely far from any point with $y > 0$.) Given a line q and a point A not lying on that line, show that we can use your method from part (b) to construct *two* lines passing through A that are parallel to q , and that between these two lines lies a pencil of lines passing through A that never meet q . (Google “pencil of lines” if you have not heard the phrase before!)

3) Drums and Membranes: The shape of a distorted drumskin is described by the function $h(x, y)$, which gives the height to which the point (x, y) of the flat undistorted drumskin is displaced.

- a) Show that the area of the distorted drumskin is equal to

$$\text{Area} = \int dx dy \sqrt{1 + \left(\frac{\partial h}{\partial x}\right)^2 + \left(\frac{\partial h}{\partial y}\right)^2},$$

where the integral is taken over the area of the flat drumskin.

- b) Show that for small distortions, the area reduces to

$$\mathcal{A}[h] = \text{const.} + \frac{1}{2} \int dx dy |\nabla h|^2,$$

where $\nabla \equiv \mathbf{e}_x \partial_x + \mathbf{e}_y \partial_y$.

- c) Show that if h satisfies the two-dimensional Laplace equation then \mathcal{A} is stationary with respect to variations that vanish at the boundary.
- d) Suppose the drumskin has mass ρ_0 per unit area, and surface tension T . Write down the Lagrangian controlling the motion of the drumskin and derive the equation of motion that follows from it.

4) Magnetostatics: We wish to find the magnetic field $\mathbf{B}(x)$, $x \in \mathbb{R}^3$, produced by a (compactly supported) current distribution $\mathbf{J}(x)$ in a material with position-dependent permeability $\mu(x)$. Consider the functional

$$F[\mathbf{A}] = \int_{\mathbb{R}^3} \left\{ \frac{1}{2\mu(x)} |\nabla \times \mathbf{A}|^2 - \mathbf{J} \cdot \mathbf{A} \right\} d^3x,$$

where $\mathbf{A}(x)$ is a candidate vector potential for the field $\mathbf{B} \equiv \nabla \times \mathbf{A}$.

- i) Show that $F[\mathbf{A}]$ is *gauge invariant* (*i.e.* it is unchanged in value when \mathbf{A} is replaced by $\mathbf{A} + \nabla \chi$) provided that $\nabla \cdot \mathbf{J} = 0$.
- ii) Show that the stationarity condition $\delta F / \delta A_i(x) = 0$ leads to \mathbf{A} obeying the appropriate Maxwell equation. (You may assume that $\nabla \times \mathbf{A}$ is zero at infinity.)

- iii) Replace the current distribution by the physical magnetic field it produces, $\nabla \times \mathbf{H}_{physical} = \mathbf{J}$. By “completing the square,” show that $F[\mathbf{A}]$ takes its *minimum* value when \mathbf{A} is a vector potential for the physical magnetic field.