

Physics 489 S 04 Lecture 3

Drude theory of transport in metals (Ashcroft and Mermin, chapter 1,3) Non-Equilibrium Properties

1. Drude theory of transport in metals (Ashcroft and Mermin, chapter 1)
Model: Z electrons detach from each atom and flow freely through solid.
Assumptions for classical theory of transport:
 - (1) Between collisions, electrons are free and independent.
 - (2) Collision randomizes direction of velocity.
 - (3) Relaxation time assumption. Mean free time = τ .
 - (4) Collision thermalizes magnitude of velocity to local temperature

2. DC electrical conductivity.
 $V = IR$ or $\mathbf{j} = \sigma \mathbf{E}$ with $\sigma = \frac{ne^2\tau}{m}$
 $\tau = \frac{\ell}{v}$; T-dependence through properties of τ or ℓ .
Classical mean free paths $\ell = \tau v_{gas} \approx 10A$ seemed reasonable to Drude.

3. Semiclassical Theory
Quantum (Fermi) mean free path $\ell = \tau v_{Fermi} \approx 1000A$ at room T - remarkably long.
Observed experimentally. ℓ can be very long at low T.
What about uncertainty principle in using a classical description?
Valid *only* for $\ell p_f > \hbar/2$.

4. Thermal conductivity.
 $\mathbf{j}_Q = -\kappa \nabla T$ with $\kappa = \frac{1}{3} v^2 \tau c_V$
Experimental Wiedemann-Franz law $\frac{\kappa}{\sigma T} = \text{constant} = \text{Lorentz number}$.
Drude gives $\frac{\kappa}{\sigma T} = \frac{3}{2} \left(\frac{k_B}{e^2}\right)^2$ with fortuitous quantitative agreement.
Classical $c_V = \frac{3}{2} k_B$ becomes Quantum $c_V \propto T$ (100 times smaller).
Classical $\frac{1}{2} m v^2 = \frac{3}{2} k_B T$ becomes Quantum $\frac{1}{2} m v^2 = E_F$ (100 times larger).
Semiclassical theory gives $\frac{\kappa}{\sigma T} = \frac{\pi^2}{3} \left(\frac{k_B}{e^2}\right)^2$ with even better agreement.

5. Thermopower and Seebeck Effect – Discuss briefly
Must have electric field to prevent current flow in presence of temperature gradient.
 $E = Q \nabla T$, $Q = -\frac{1}{3ne} c_V$

6. Dynamical response: $\frac{d\mathbf{p}}{dt} = \mathbf{f} - \frac{\mathbf{p}}{\tau}$
 AC conductivity: $\sigma(\omega) = \sigma_0 \frac{1}{1-i\omega\tau}$ with $\sigma_0 = \frac{ne^2\tau}{m}$.
7. Optical Properties - response in electromagnetic field.
 Dielectric function $\epsilon(\omega) = 1 + 4\pi i\sigma(\omega) = n(\omega)^2$, where n is the refractive index.
 Wave velocity $= \frac{\omega}{k} = \frac{c}{n(\omega)}$
 Examples:
- When $\omega\tau \ll 1$, $\sigma(\omega)$ is mainly real [$\frac{1}{\tau} \approx 10^{13} Hz$].
 Imaginary part of ϵ gives absorption; large at low ω .
 Real part of n gives wave velocity.
 - When $\omega\tau \gg 1$, $\sigma(\omega)$ is mainly imaginary
 $\epsilon(\omega) = 1 - (\frac{\omega_P}{\omega})^2$ is real, with ω_P = plasma frequency
 When $\omega < \omega_P$, n is imaginary, reflectivity = 1, metal is opaque and shiny.
 When $\omega > \omega_P$, n is real and the metal becomes transparent (in ultraviolet).
 When $\omega = \omega_P$, plasma oscillation occurs.
8. Hall effect.
 Must have electric field to prevent transverse current flow in presence of magnetic field.
 $E_y = R_H H_z j_x$ with $R_H = \frac{-1}{nec}$
 Depends on density and sign of charges only.
9. Major failures of free, independent electron theory (Ashcroft and Mermin, chapt. 3)
- Solids have an enormous range of behaviors:
 - Some solids are insulators; Conductivity = 0 at T=0); No linear T term in specific heat
 - Semiconductors
 - Sometimes hall constant of real metals sometimes have opposite sign.
 - Magnetism, superconductivity,