489 S 04 Lecture 10

## $\begin{array}{c} Physics~489~S~04~Lecture~10\\ Quantized~Lattice~Vibrations,~Thermal~Properties~(A~\&~M~chapter~23) \end{array}$

1. Review of classical theory:

Within harmonic approximation, the exact solution of lattice vibrations is a set of independent harmonic oscillators, labelled by wavevector  $\mathbf{k}$ 

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For each **k** there are s = D \* S modes (D = dimension; S = # atoms per cell) with frequency  $\omega_s(\mathbf{k})$ 

There are D acoustic modes with  $\omega \to 0$  as  $\mathbf{k} \to 0$ 

- 2. Quantum Theory (Using results stated in Review of Statistics, Lecture 1) Each oscillator of frequency  $\omega$  has quantized energies  $E_n = (n+1/2)\hbar\omega$  Quantum particles (phonons) obeying Bose statistics Average occupation: Planck distribution,  $\langle n \rangle = (e^{\beta\hbar\omega} 1)^{-1}$ ,  $\beta = 1/k_BT$  Total energy density:  $(1/V)E_{tot} = (1/V)\sum_{s,\mathbf{k}}\hbar\omega_s(\mathbf{k})\left(\langle n_s(\mathbf{k}) \rangle + \frac{1}{2}\right)$
- 3. Specific Heat  $C_V = (1/V)dE_{tot}/dT = (1/V)\sum_{s,\mathbf{k}}\hbar\omega_s(\mathbf{k})\frac{d< n_s(\mathbf{k})>}{dT}$ For  $k_BT >> \hbar\omega$ ,  $< n>= k_BT/\hbar\omega$ , and  $C_V$  approaches the classical value  $Dk_B$  per nucleus, D= dimension For low T only vibrations with  $\omega_s(\mathbf{k})$  of order  $k_BT$  contribute to  $C_V$ For sufficiently low T in all solids, only acoustic modes with  $\omega=v_{sound}k$  contribute and  $C_V \propto T^D$ , where D= dimension
- 4. Debye model (correct for acoustic modes, qualitatively correct for other modes) Approximate all frequencies by  $\omega = v_{sound}k$ ,  $v_{sound} = \text{spherical average}$   $k_{max} = k_D$  defined to sum to the right number of modes In 3 dimensions, k is in a sphere with volume  $(4\pi/3)k_D^3 = (2\pi)^3/V_{atom}$  Debye Temperature  $\Theta_D$  defined by  $k_B\Theta_D = \omega_{max} = v_{sound}k_D$  At low T,  $C_V = 3(N/V)k_B(4\pi^4/5)(T/\Theta_D)^3$  (in 3D) Can consider  $C_V$  at low T to define  $\Theta_D$  for a given material
- 5. Einstein Model:  $\omega_s(\mathbf{k})$  modelled as constant  $\omega_s$  independent of k (Appropriate for optic modes only) Gives  $C_V \propto exp(-\beta\hbar\omega_s)$  for low T
- 6. Typical magnitudes
  Phonon energy  $\hbar\omega = 0$  to a maximum value of  $\approx 10^{-2}$  to  $\approx 10^{-1}$  eV
  Debye energy  $-\approx 10^{-2}$  to  $\approx 10^{-1}$  eV, Debye Temperature  $-\approx 100$  to  $\approx 1000$  K
- 7. Density of states = number of states per unit energy per unit volume Determines all quantities that depend only on energy, e.g.,  $C_V$   $g(\omega) = (1/V) \sum_{s,\mathbf{k}} \delta(\omega \omega_{s,\mathbf{k}})$  For vibrations at low  $\omega$ ,  $g(\omega) \propto \omega^{D-1}$ , where D is the dimension.

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Figure 1: Phonon dispersion and density of states in 1 and 3 dimensions

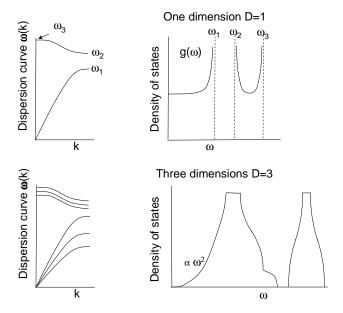


Figure 2: Debye approximation to the dispersion and density of states in 3 dimensions. Typical specific heat which is reasonably approximated by Debye model.

