

**Physics 489 S 04 Lecture 22**

**Semiconductors I: Properties of homogeneous materials, A & M Ch. 28,  
Kittel Ch. 8.**

1. Semiconductors are materials whose conductivity can be modified greatly in controllable ways by temperature, adding foreign atoms, and electric fields  
Usually they are insulators at  $T=0$  with narrow gaps ( $\approx 0$  to  $\approx 2$  eV) in their perfect crystalline form
2. Key properties: Magnitude of gap, effective masses of bands at top of filled bands and bottom of empty bands

Examples: (Most values are from Kittel, Chapt. 8)

	Gap		m*e		m*h
	OK	300K			
Si	1.17 eV	1.12 eV	1.0 , 0.2me		0.52, 0.16
Ge	0.75	0.67	1.6, 0.082		0.34, 0.043
GaAs	1.5	1.4	0.066		0.5, 0.082
InSb	0.23	0.16	0.015		0.39, 0.021

3. Electrons and Holes

Occupation of band-like states given by Fermi function

For  $E_c - \mu \gg k_B T$ , and  $\mu - E_v \gg k_B T$ ,

Probability for state in conduction band to be full =

$$f(E) \approx \exp(-(E - \mu)/k_B T), \text{ for } E > E_c$$

Probability for state in valence band to be empty =

$$1 - f(E) \approx \exp(-(\mu - E)/k_B T), \text{ for } E < E_v$$

Integrating in effective mass approx. gives the density of electrons and holes:

$$n = N \exp(-(E_c - \mu)/k_B T),$$

$$p = P \exp(-(\mu - E_v)/k_B T),$$

where  $N = 2(m_e^* k_B T / 2\pi \hbar^2)^{3/2}$ , etc., which varies slowly with T.

4. Law of mass action:  $np = NP \exp(-E_g/k_B T)$  independent of Fermi Energy  $\mu$ .
5. Intrinsic semiconductors (perfect crystals) at  $T > 0$ ,  
Charge neutrality:  $n = p = n_i(T) = (NP)^{1/2} \exp(-E_g/k_B T)$   
 $\mu$  is at center of gap at  $T=0$ ,  $\mu(T) = \mu(0) + 0.5 k_B T \ln(N/P)$
6. Extrinsic semiconductors: doped crystals  
Charge of carriers:  $n - p = N_{donors} - N_{acceptors}$ ,  $np = n_i(T)^2$   
If  $N_{donors} > N_{acceptors}$ ,  $\mu$  is increased nearer conduction band, n increased, p decreased  
If  $N_{donors} < N_{acceptors}$ ,  $\mu$  is decreased nearer val. band, p increased, n decreased

## 7. Conduction properties

Mobilities:  $\mu_e = ne^2\tau_e/m_e^*$ ,  $\mu_h = ne^2\tau_h/m_h^*$

Conductivity:  $\sigma = n_e\mu_e + n_h\mu_h$  - electron and hole always add!

Hall Effect:  $R_H = (1/ec)(p - n(\mu_e/\mu_h)^2)/(p + n(\mu_e/\mu_h)^2)$

Thermoelectric Effect:  $\Pi = \text{heat} / \text{charge per carrier}$  - can determine sign of carriers

## 8. Doping – Donors and Acceptors: Examples: As or Ga in Si

## 9. Bound States of electrons or holes - Hydrogenic Binding:

Binding =  $E_d$  or  $E_a = Ryd(m^*/\epsilon^2) \approx 13.6eV(0.1/10^2) \approx 0.01 eV$

Radius =  $R \approx \text{bohr radius}(\epsilon/m^*) \approx 0.52A(10/0.1) \approx 50A$

Because of interactions among electrons, there is an energy cost to put two electrons (or two holes) on one impurity site. This favors having one electron (or hole) which has a spin that can be detected by magnetic resonance. Example of effects discussed in Lecture 1 and given in problem(s).

## 10. Thermal excitation of free carriers from donor or acceptor sites

In n type crystal,  $N_{donors} > N_{acceptors}$ , at low T the density of free carriers is:

$$n/N_d = 2N \exp(-(E_d/2)/k_B T), \text{ where } E_d \text{ plays the role of a gap}$$

Similarly for p type crystals

At high T, n and p approach the intrinsic regime  $n \approx p \approx n_i(T)$

## 11. Degenerate metal formed at T=0 by overlap of donor or acceptor sites

If the impurity states overlap strongly, the exclusion principle forces the electrons to tend to pair with opposite spins - for large overlap they form a metal:

Mott Transition - insulator to metal as density increases

Very roughly this happens when the hydrogenic states overlap strongly, i.e. , for densities where  $N_{donors}$  or  $N_{acceptors} > (1/4\pi R^3/3)$

Heavily doped semiconductors are metallic with finite conductivity  $\sigma$  at T= 0