

**561 Fall 2005 Lecture 16****Brief Survey of Strongly-Interacting Electron Systems**

Most complete review to date: M. Imada, A. Fujimori, and Y. Tokura, "Metal Insulator Transitions", "Rev. Mod. Phys. 70, 1039 (1998), which reviews theory and experiment on correlated-electron systems, especially transition metal oxides.

See also E. Dagotto, "Correlated Electrons in High Temperature Superconductors", Rev. Mod. Phys. 66, 763 (1994).

A. Damascelli and Z. Hussain and Z-X. Shen, "Angle-resolved photoemission studies of the cuprate superconductors", Rev. Mod. Phys. 75, 463 (2003).

Colossal Magnetoresistance: M. B. Salamon and M. Jaime, "The physics of manganites: Structure and transport", Rev. Mod. Phys. 73, 583 (2001).

**1. Zaanen-Allen-Sawatzky classification scheme**

Discussed in Imada, et al.

A. Example that was believed to be the realization of a Mott-Hubbard insulator:  $V_2O_3$

- The only states near  $E_F$  are 3d states
- Hubbard model with 10 3d states
- Metal-insulator transition believed to be mainly a property of the interacting electrons in the Hubbard model for the 3d states
- Other states (Oxygen 2p, 4s, etc.) are either strongly bound below the Fermi energy or well above - not involved crucially in the transition
- See figures from Imada review

B. Charge Transfer Insulators

- NiO - nominally  $d^8$  antiferromagnetic insulator with Gap  $\approx 4eV$
- Photoemission shows  $d^8 \rightarrow d^7$  is far below  $E_F$  by about 10 eV.
- $d^8 \rightarrow d^9$  lowest empty states
- Oxygen 2p states (band-like "Ligand" states highest occupied states
- Very different from band calculations, which find states at Fermi energy to be mainly 3d like a Mott-Hubbard system
- Comparison of  $Cu_2O$  ("Ordinary" Band Insulator with filled d shell) and  $CuO$  (one hole per Cu, which leads to correlated antiferromagnetic insulator )

**2. Cu Oxide planar systems: Antiferromagnetic Insulators to High-Tc superconductors**

- Gross Features of Bands: Simple nondegenerate Hubbard model with one band (mainly Oxygen 2p in character) crossing  $E_F$

- Undoped systems (e.g.  $CaCuO_2$ ;  $La_2CuO_4$  (La214);  $YBa_2CuO_3$  (YBa123); (Bi,Tl)Ca 221; (Bi,Tl)BaCa2212; etc. all have similar  $CuO_2$  square layers - counting gives a half-filled nondegenerate Hubbard model
- Would be a metal except that interactions lead to correlated antiferromagnetic insulator
- Doping with holes (e.g.,  $(La_{1-x}Sr_x)_2CuO_4$ ) leads to metal and superconductor for "optimal range of doping"
- States for added holes are mainly O p states - leads to idea of "O metal" - "t-J" model for holes on O that interact with spins on the Cu (spin 1/2 on each Cu)

$$H = t \sum_{\langle ij \rangle \sigma} c_{i\sigma}^\dagger c_{j\sigma} + J_1 \sum_i \sigma_i \cdot \mathbf{S}_i + J_2 \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j, \quad (1)$$

where  $\langle ij \rangle$  denotes nearest neighbor pairs.

The last term involving  $J_2$  leads to antiferromagnetic order, and the holes tend to promote ferromagnetism - i.e., a competition or frustration.

- Gross Features of Photoemission from metallic systems
- Very weak features near  $E_F$  in the doped materials are the metallic bands that lead to superconductivity
- Angle Resolved Photoemission (ARPES) - See review by Domacelli, et al. Many examples that show the Fermi surface - anomalous behavior of spectra that are non-Fermi-liquid-like as a function of  $E - E_F$ 
  - Examples from the High-Tc materials - recent very high precision experiments - Figures from Lect. 1
  - d-wave symmetry gap
  - "Pseudogap" in underdoped materials - depression in spectra near  $E_F$  - similar to the superconducting gap, but above  $T_c$ .

### 3. "Colossal Magneto-Resistance" (CMR) materials

- Cubic perovskite structure  $La_{1-x}Sr_xMnO_3$  (with some Jahn-Teller distortion)
- $LaMnO_3$  is a ferromagnetic insulator
- Can be understood in terms of a type of band picture: Mn d states form large local 3/2 spins - modelled as classical spins - plus additional itinerant electrons that move in a field of the spins - double exchange mechanism (Zener, 1951) - the quantum version is a variant of the "t-J" model where  $J_2$  is small so that one finds ferromagnetism
- Doped system (each Sr subtracts one electron) has metal-insulator (M-I) transition
- Leads to large coupling of conductivity and magnetism "Colossal Magneto-Resistance" near the phase transition

- What happens at M-I transition - matter of controversy:
  - Just disorder in spins (C. M. Varma, PRB 54, 7328 (1996))
  - Disorder in spins plus Jahn-Teller distortions (Millis, et al. PRL 74, 5144 (1995))
    - in this model the dynamic J-T effect makes resistivity very large in metallic phase

**4. Stripes** Formation of inhomogeneities or "stripes" or "pasta phases" - seen in some materials - theories by Zaanen and by Emery and Kivelson