Let's now return to the last terms in the Standard Model Lagrangian we haven't studied yet:

\[ L \sim -\frac{1}{4} W^{\mu\nu} W_{\mu\nu} - \frac{1}{4} B^{\mu\nu} B_{\mu\nu} + (D_{\mu} H)^*(D_{\mu} H) + m^2 H^*H - \lambda (H^*H)^2 \]

As with the Abelian case, the wrong-sign mass term will lead to spontaneous symmetry breaking. First let's minimize the potential:

\[ V(H) = -m^2 H^*H + \lambda (H^*H)^2 \]

\[ \frac{\partial V}{\partial H} = -m^2 H + 2\lambda H (H^*H) = 0 \quad \Rightarrow \quad H^*H = \frac{m^2}{2\lambda}. \]

Note that this condition only determines the norm of \( H \), \( |H|^2 \equiv H^*H + H^*H \). Since \( SU(2) \) gauge transformations rotate \( H \to e^i \theta H \), we can choose a gauge where \( H = 0 \).

Write \( H = \exp \left( 2i \frac{\pi(x)\tau^i}{\nu} \right) \left( \frac{\nu}{\nu + \frac{\nu(x)}{2}} \right) \nu / \nu = \frac{m}{\sqrt{\lambda}}, \tau^a = \frac{1}{2} \sigma^a \ (SU(2) \text{ generators}) \)

(The \( \frac{1}{\sqrt{2}} \) is there so \( D_\mu H^* D^\mu H \) contains \( \frac{1}{2} \partial_\mu \partial^\mu h \), as appropriate for a real scalar \( h \)). Use unitary gauge to set \( \pi(x) = 0 \) everywhere.

Covariant derivative is \( D_\mu H = \partial_\mu H - ig W^a_\mu \tau^a H - \frac{i}{2} g' B^a_\mu H \)

First, let's look only at the terms without \( h \) (i.e., set \( h = 0 \) for now)

\( H \to \frac{\nu}{\sqrt{2}} (0) \). Since \( B \) is Abelian, rewrite non-derivative term as

\[ -i g \left( W^a_\mu \tau^a + \frac{i}{2} \frac{g'}{g} B^a_\mu \right) = -i g \left( W^a_\mu \sigma^a + \frac{g'}{g} B^a_\mu \right) \]

\[ \Rightarrow \quad \text{Hermitean} \]

\[ \Rightarrow \quad |D_\mu H|^2 = g^2 \left( \frac{g'}{g} B^a_\mu + W^a_\mu \right)^2 \left( \frac{g'}{g} B^a_\mu - W^a_\mu \right)^2 \left( W^a_\mu - i W^a_\mu \right) \left( W^a_\mu + i W^a_\mu \right) \]

\[ \Rightarrow \quad |D_\mu H|^2 = \frac{g^2}{8} \left( \frac{g'}{g} B^a_\mu + W^a_\mu \right)^2 \left( \frac{g'}{g} B^a_\mu - W^a_\mu \right)^2 \left( W^a_\mu - i W^a_\mu \right) \left( W^a_\mu + i W^a_\mu \right) \]
the three gauge bosons which become massive are

\[ W^\pm, W^3 \text{ (mass } m_w = \frac{gV}{2}) \text{, and } \frac{g}{2} \beta - W^3. \]

However, QFT tells us we need to preserve the normalization of the gauge kinetic terms, so we should perform a rotation of the fields \( \beta_n \) and \( W^3 \) to define the mass eigenstate. Specifically:

\[
\begin{pmatrix}
Z_n \\
A_n
\end{pmatrix} = \begin{pmatrix}
\cos \theta_n & -\sin \theta_n \\
\sin \theta_n & \cos \theta_n
\end{pmatrix}
\begin{pmatrix}
W^3 \\
\beta_n
\end{pmatrix}, \text{ with } \tan \theta_n = \frac{g'}{g} \text{ (Weinberg angle). Then}
\]

\[-\frac{1}{4} W^3_{\mu \nu} W^{3 \mu \nu} - \frac{1}{4} \beta_{\mu \nu} \beta^{\mu \nu} \rightarrow -\frac{1}{4} Z_{\mu \nu} Z^{\mu \nu} - \frac{1}{4} F_{\mu \nu} F^{\mu \nu} \text{ (rotations preserve norm)}\]

\[ Z_{\nu \mu} \delta_{\nu \mu} \quad \delta Z_{\nu \mu} \delta A_{\nu \mu} \]

Also, \( \frac{g'}{g} \beta - W^3 = \tan \theta_n \beta_n - W^3 = -\frac{1}{\cos \theta_n}(W^3 \cos \theta_n - \beta_n \sin \theta_n) = -\frac{Z_n}{\cos \theta_n} \]

\[ \Rightarrow \text{ we identify } Z_n \text{ with the Z boson and } A_n \text{ with the photon, and their Lagrangian is } \]

\[ L = -\frac{1}{4} F_{\mu \nu} F^{\mu \nu} - \frac{1}{4} Z_{\mu \nu} Z^{\mu \nu} + \frac{1}{2} m^2 Z_n Z_n, \]

with \( m^2 = \frac{1}{2\cos \theta_w} gV \). But photon remains massless! We can express this as

\[ SU(2) \times U(1)_Y \rightarrow U(1)_{EM} \text{; the electroweak symmetry is spontaneously broken to electromagnetism.} \]

What about electric charge? We want to find the part of the gauge kinetic term that couples to the photon, which is a linear combination of \( W^3 \) and \( \beta_n \). We have previously identified \( Z^3 + Y \) as the electric charge, so let's find its coefficient in the covariant derivative:

\[ D_m = \partial_m - ig W^a \nabla^a - ig' Y \beta_n \]

\[ = \partial_m - i g \frac{g'}{g} (W^+ n \tau^i + W^- n \tau^{-i}) - i \frac{1}{\sqrt{g^2 + g'^2}} Z_n (g^2 \tau^3 - g'^2 Y) - i \frac{gg'}{\sqrt{g^2 + g'^2}} A_n (\tau^3 + Y) \]

where \( W^\pm = \frac{1}{\sqrt{2}} (W^l \pm i W^-) \) and \( \tau^\pm = \frac{1}{\sqrt{2}} (\tau^1 \pm i \tau^2) \).
From the coefficient of $T^3 + Y$, we can extract the electric charge:

$$e = \frac{g'g}{\sqrt{g^2 g'}} = g \sin \theta_w = g' \cos \theta_w$$

Finally, we can treat $W^\pm_m$ as a complex vector field, with mass term $m_w^2 W^+_m W^-_m$, where

$$m_w = \frac{g v}{2}$$

The notation $W^\pm$ is appropriate, since $W^\pm$ have electric charges, $\pm 1$.

Let $(T^3 + Y)$ act on $W^+_m = W^+_m T^3$. $Y$ acts as $0$ since $SU(2)_L$ and $U(1)_Y$ commute. $W$ is in the adjoint of $SU(2)$, so $T$ acts as a commutator:

$$[T^3, T^\pm] = \pm T^\mp$$

(recall raising and lowering operators from QM!)

so $W^\pm$ have electric charge $\pm 1$. By similar reasoning, $Z$ is neutral.

Predictions of the Higgs mechanism:

- The standard model contains a massless photon, a neutral massive gauge boson $Z$, and a charged massive gauge boson $W$. Their masses are related as $m_Z = \frac{m_W}{\cos \theta_w}$, so $W$ is lighter than the $Z$.

- Electric charge is related to the gauge couplings $g$ and $g'$ as $e = g \sin \theta_w$.

- Four parameters in the Lagrangian $g, g', m$, and $\lambda$.
  - Four physical parameters $e, \theta_w, m_w$, and $m_h = \frac{\lambda v^2}{2}$.
  - Unfortunately, $m_h$ independent from other three! Can't predict the Higgs mass.

- Standard Model fields couple to $W^\pm$ and $Z$ through covariant derivative $D_m = \partial_m - i \frac{g}{2} (W^+_m T^3 + W^-_m T^-) - i \frac{g}{\cos \theta_w} Z^+_m Q^+_2 - i e A^+_m Q^+_2$ where

$$Q^+_2 \equiv T^3 - \sin^2 \theta_w Q$$ is the "charge" under the $Z$-boson. Different for RadL field!
Why do we need the Higgs boson in the first place? Even if we knew nothing about the Yukawa terms and the underlying gauge invariance, the existence of a massive vector boson with self-interactions is pathological without the Higgs.

To see this, consider the process $W_L^+Z_L \rightarrow W_L^+Z_L$, where the subscript $L$ means longitudinally polarized. This process only exists for massive (since massless vectors are transverse), nonabelian (since abelian vectors have no self-interactions) vectors.

The component of the $Z$ which interacts with the $W$ is $W^3$, so this is very much like gluon-gluon interactions with some $SU(2)$ group theory factors instead of $SU(3)$. (See Schwartz Sec. 29.1 for the full set of Feynman rules.) First let's carefully define polarization vectors:

For general $p^\mu$:

$$\epsilon_1 = \frac{i}{m_w} (p_1, 0, 0, E) \quad \epsilon_2 = \frac{1}{m_w} p_2^\mu + \frac{2m_w}{t-2m_w^2} p_+^\mu$$

(slightly for $\epsilon_3, \epsilon_4$, with $t = (p_1 - p_3)^\mu = (p_2 - p_4)^\mu$)

These satisfy $\epsilon_i \cdot p_i = 0$, not normalized but we'll make it for argument to follow.

First matrix element:

$$iM = (i e \cot \theta_W)^2 \epsilon_1^\mu \epsilon_2^\nu \epsilon_3^{\alpha \beta} \epsilon_4^{\gamma \delta} \times \frac{i}{s-m_w^2} (-\gamma^\lambda k^\mu + \frac{1}{m_w^2} k^\lambda k^\mu) \times$$

$$\left[ -\gamma^\nu (p_2 - p_1)^\lambda + \gamma^\lambda (p_2 k)^\nu - \gamma^\lambda (k p_1)^\nu \right] \left[ -\gamma^{\alpha \delta} (p_4 - p_3)^\mu + \gamma^{\mu \delta} (p_4 k)^\alpha - \gamma^{\mu \delta} (k p_3)^\alpha \right]$$

$w/k = p_1 + p_2 = p_3 + p_4$. 

Putting the Higgs back in
Plugging in polarization vectors (since we have fixed our initial spin states):

\[ M_{s} = \frac{e^{2} \cot \theta_{w}}{4m_{w}^{2} m_{z}^{2}} \left[ 2s_{u} + s^{2} - 2m_{w}^{2} \frac{3s_{u} + s^{2}}{s_{u} + u} + 2m_{z}^{2} \frac{s^{2} - 3s_{u} - 2u^{2}}{s_{u} + u} - \frac{m_{w}^{2}}{m_{w}^{2}} s + \mathcal{O}(1) \right] \]

This looks like a problem: at large enough \( s \), amplitude grows without bound, eventually we will violate unitarity.

To understand this behavior, look at \( E > m_{w} \), where \( E_{w} = \frac{1}{m_{w}} \)

\[ M \sim (\text{propagator}) \times (\text{polarization})^{q} \sim E^{q} \sim s^{q} \]

\[ \frac{E_{w}}{m_{w}^{2}} \sim \left( \frac{E}{m_{w}} \right)^{q} \]

\[ \sim \frac{1}{m_{w}^{q}} \]

Things are actually not as bad as they seem:

\[ M_{u} = \frac{e^{2} \cot \theta_{w}}{4m_{w}^{2} m_{z}^{2}} \left[ 2s_{u} + u^{2} - 2m_{w}^{2} \frac{3s_{u} + u^{2}}{s_{u} + u} + 2m_{z}^{2} \frac{u^{2} - 3s_{u} - 2u^{2}}{s_{u} + u} - \frac{m_{w}^{2}}{m_{w}^{2}} u + \mathcal{O}(1) \right] \]

\[ M_{+} = \frac{e^{2} \cot \theta_{w}}{4m_{w}^{2} m_{z}^{2}} \left[ -s^{2} - 4s_{u} + 2m_{w}^{2} + m_{z}^{2} \frac{s^{2} + 6s_{u} + u^{2}}{s_{u} + u} + \mathcal{O}(1) \right] \]

So there is a partial cancellation (much like the Abelian case, where 3- and 4-point couplings are related):

\[ M_{\text{tot}} = -\frac{m_{w}^{2}}{4m_{w}^{2}} e^{2} \cot \theta_{w} (s + u) + \mathcal{O}(1) = \frac{t}{v^{2}} + \mathcal{O}(1) \]

But this still grows with energy! Specifically, using partial-wave unitarity (Schwartz 24.1.5), we must have \( \frac{E^{2}}{v^{2}} \times \frac{1}{32 \pi} < 1 \)

\( \Rightarrow E < \sqrt{32 \pi} v \approx 2.9 \text{ TeV} \). Therefore, some new physics must appear at this energy scale to restore unitarity.
In the Standard Model, the Higgs rescues unitarity. (Before 2012 we did not know that Nature picked this solution!)

Higgs interactions are simple to determine: just take \( v \to v + h \)

\[
\begin{align*}
\Rightarrow m_w^2 w^+ w^- & = \frac{v^2}{4} w^+ w^- \quad \Rightarrow (v + h)^2 \frac{v^2}{4} w^+ w^- \\
& = \frac{2h}{v} \frac{v^2}{4} w^+ w^- + \cdots \\
& = 2 \frac{h}{v} m_w^2 w^+ w^- + \cdots \\
\text{(same for } Z)\end{align*}
\]

Importantly, this implies Higgs couples proportional to mass! (will see this more next week). Here, we have an additional diagram

\[
\begin{align*}
\text{Higgs interactions determined by } v \to v + h \text{ in Lagrangian, we will do this for Yukawa terms next time.}
\end{align*}
\]