$$H = H_{hfs} + H_B = hA_{hfs} \frac{\overline{I} \cdot \overline{J}}{\hbar^2} + \frac{\mu_B}{\hbar} (g_J J + g_I I)B_z$$

(a) [3 pts] In general, finding the eigenstates and energies in this system requires a numerical solution (diagonalization of this Hamiltonian). However, we stated in class that for states with J = 1/2, as in the S_{1/2} ground states of hydrogen and the alkalis, the energies can be described by the Breit-Rabi formula

$$E = -\frac{\Delta E_{hfs}}{4F^+} - g_I m \mu_B B_z \pm \frac{\Delta E_{hfs}}{2} \sqrt{1 + \frac{4mx}{2F^+} + x^2}$$

Here, $F^+ = (I + 1/2)$, $\frac{\Delta E_{hfs}}{h} = A_{hfs}F^+$, $m = m_I \pm m_J$ is the *z* component of the total angular momentum, and the term *x* is given by $x = (g + g_I) \mu_B B_z / \Delta E_{hfs}$. Show that the state energies follow this Breit-Rabi formula for a general field value B_z .

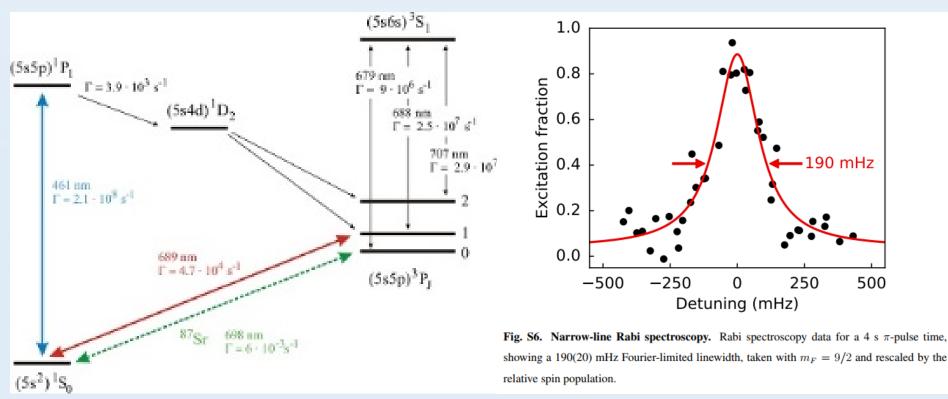
L7: 2-level systems & classical resonance

Foot 1.6, 1.8

- spin-1/2 particle (electron/proton spin)
- (pseudo)spin-1/2 particle (hyperfine, etc.)

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 - optical excitations



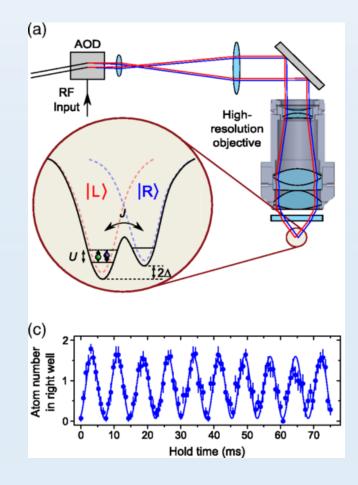
L7: 2-level systems & classical resonance

Foot 1.6, 1.8

- spin-1/2 particle (electron/proton spin)
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 - energy state of anharmonic oscillator
 - rot/vib states of molecules
 - etc.

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L7: 2-level systems & classical resonance

Foot 1.6, 1.8

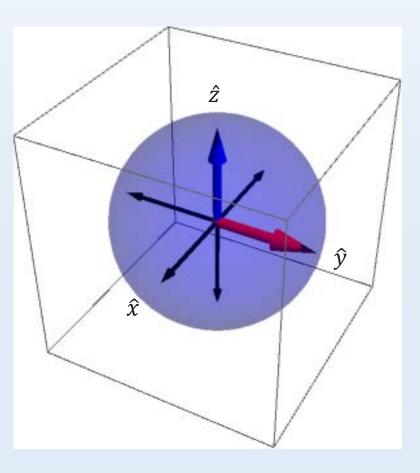
Examples of "2-level" systems

- spin-1/2 particle (electron/proton spin)
- (pseudo)spin-1/2 particle (hyperfine, etc.)
 - optical excitations
 - energy state of anharmonic oscillator
 - rot/vib states of molecules
 - etc.
- double-well site index
- interferometer path
- photon polarization / other d.o.f.
- states of a macroscopic object
- field occupation number
- etc.

Ignoring intricacies, all these 2-level systems are described by a similar formalism

Let's review some basic terminology

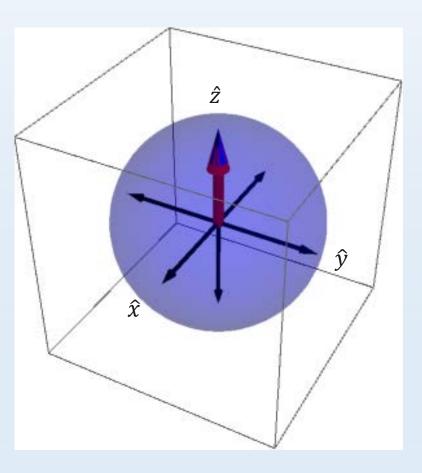
Field purely along z



$$\bar{B} = B_0 \hat{z}$$

spin in x-y plane

Field purely along z

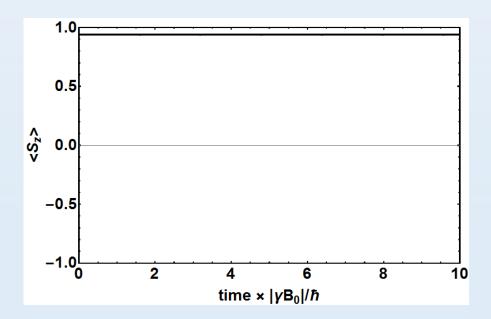


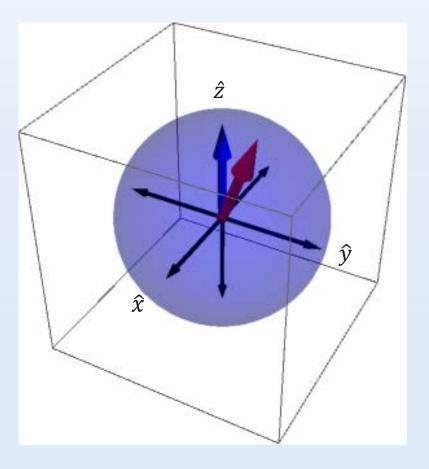
$$\bar{B} = B_0 \hat{z}$$

spin along z-axis

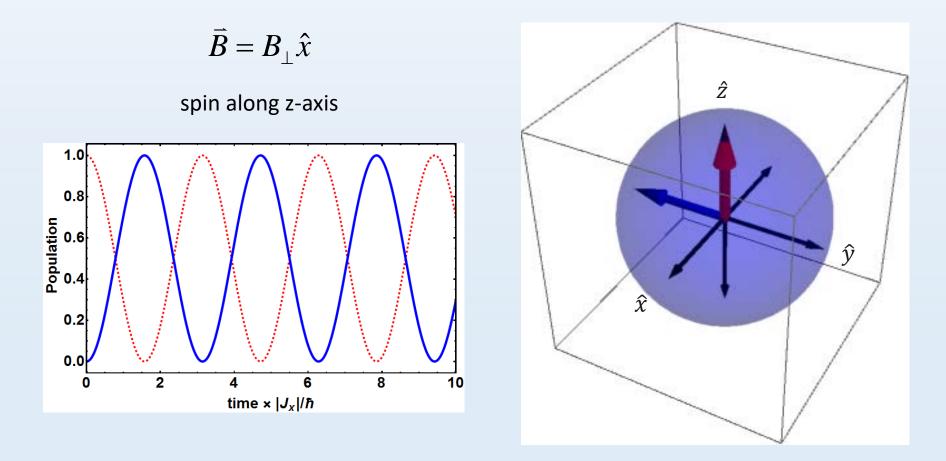
General Larmor precession

 $\bar{B} = B_0 \hat{z}$





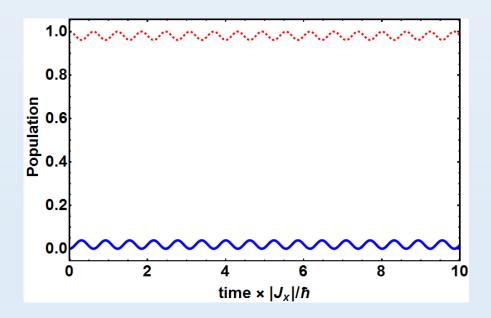
Field purely along x

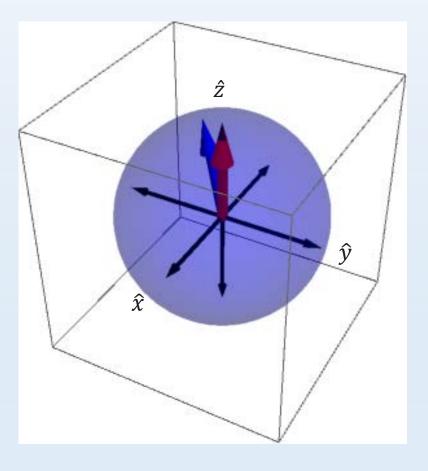


What would the dynamics look like if our spin was initially pointed along x?

Weak transverse field

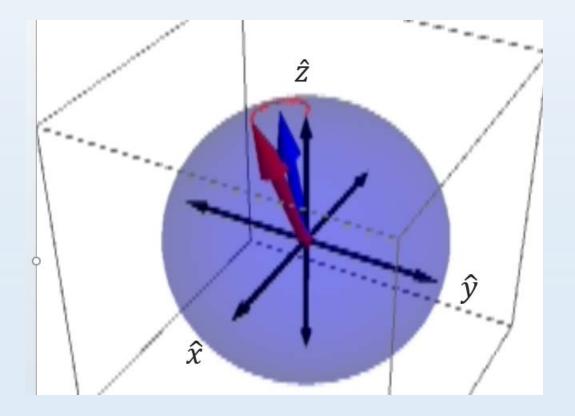
 $\vec{B} = B_0 \hat{z} + B_\perp \hat{x}$





 $B_0 = 10 B_{\perp}$

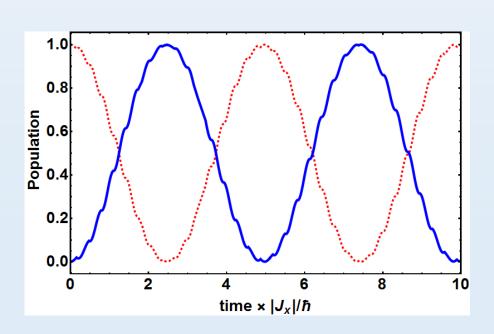
How can we get the spin to keep going towards the "south pole," i.e. -z?

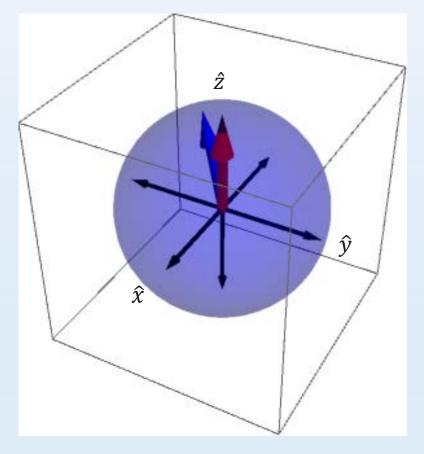


Open response: 5 points towards HW#2

Inverting the field in coordination with the precession

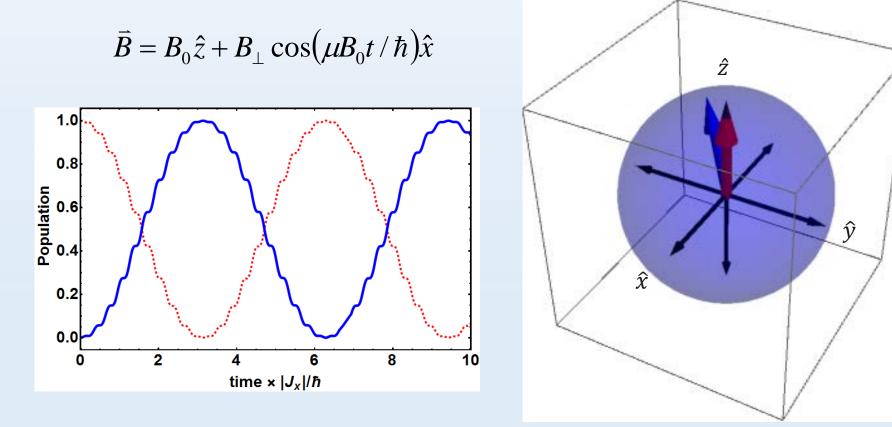
 $\vec{B} = B_0 \hat{z} + B_\perp (t) \hat{x}$





 $B_0 = 10B_\perp$

Field oscillating at the Larmor frequency



 $B_0 = 10 B_{\perp}$

many other ways as well...